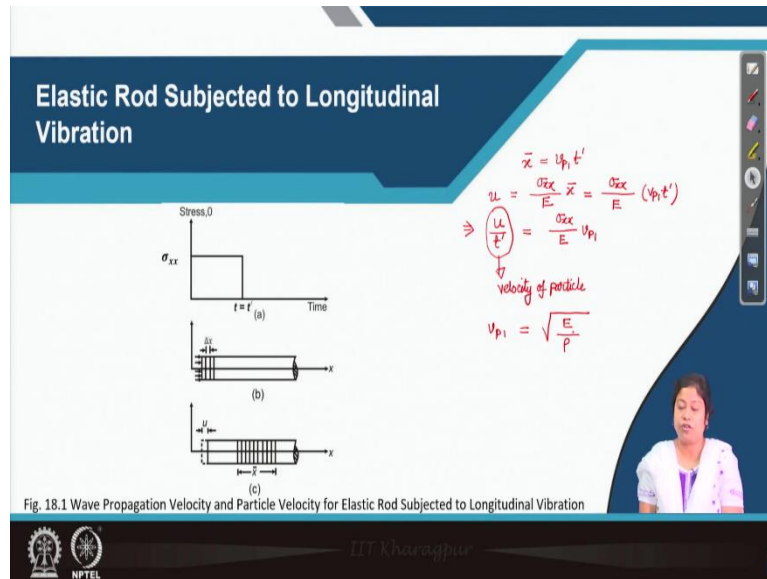


Soil Dynamics
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Lecture 18

Wave Propagation in an Elastic Rod of Finite Length

Hello everyone, today we will continue our last class on wave propagation.

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First, we will see what we have discussed last class. So, last class we have studied wave propagation in an infinite length elastic rod. So, they are first we have studied what will be the equation of motion when the longitudinal wave is propagating in the rod, thereafter, we have studied what will be the equation of motion when torsional wave is propagating in the rod.

Today, first I had like to say something about the longitudinal wave propagating through the in rod of infinite length in this diagram, what we can see in this diagram, we can see that u is the here you can see u is the this displacement or I can see completion of the rod because of the propagation of longitudinal wave in this rod and \bar{x} is the length of that element which is under consideration.

Then what will be the expression for \bar{x} , I am writing here itself we can write \bar{x} is equal to $v_{p1} t'$, just give me one minute here, yes, $v_{p1} t'$ likewise, we can write that displacement u also in this case. So, how much is the displacement is in this case it is $\sigma_{xx} \bar{x}$ divided by E which is elastic modulus times \bar{x} .

Now, in place of \bar{x} what I can write I can write the v_p 1 times t dashed, so, let us do that σ_x divided by E times v_p 1 times t dashed. Now, from this we can calculate the velocity of the particles by calculating u divided by t dashed which is equal to σ_x divided by E times v_p 1. So, here these term u divided by t dashed is nothing but the velocity of a particle.

So, from this what we can see, velocity of the particle depends upon σ_x whereas, the velocity of the longitudinal wave depends upon only the material properties, which are elastic modulus and the density of the material. That means, v_p 1 we have already seen is equal to E divided by ρ , where E is elastic modulus and ρ is the density of the material.

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Elastic Rod of Finite Length Subjected to Longitudinal Vibration

- For longitudinal wave: $\frac{\partial^2 u}{\partial t^2} = v_{p1}^2 \frac{\partial^2 u}{\partial x^2}$ where, $v_{p1} = \sqrt{\frac{E}{\rho}}$
- General solution: $u = F(x)G(t)$
- Therefore, $\dot{u} = F(x)\dot{G}(t)$ and $\ddot{u} = F(x)\ddot{G}(t) = \frac{\partial^2 u}{\partial t^2}$
- Similarly, $u' = F'(x)G(t)$ and $u'' = F''(x)G(t) = \frac{\partial^2 u}{\partial x^2}$
- Since, $\frac{\partial^2 u}{\partial t^2} = v_{p1}^2 \frac{\partial^2 u}{\partial x^2}$
- Therefore, $F(x)\ddot{G}(t) = v_{p1}^2 F''(x)G(t)$
- Or, $v_{p1}^2 \frac{F''(x)}{F(x)} = \frac{\ddot{G}(t)}{G(t)} = -\omega^2$

Now, let us see when elastic rod of finite length is subjected to longitudinal vibration. So, we already have the equation for wave propagation in one direction one dimensional or 1d direction. So, that equation you can see here $\frac{\partial^2 u}{\partial t^2}$ is equal to v_p 1 square times $\frac{\partial^2 u}{\partial x^2}$. So, here v_p 1 is the velocity of the body wave or we can call a longitudinal wave, we can also call it as P wave.

So, now, here if what is u , u is the displacement. So, now, we can consider that u is equal to the product of function x and function G , F is a function of space that is x and G is a function of time. Then, what is u dot, u dot is the first derivative of u with respect to the time t therefore, we can write for u dot is equal to a fixed times G dot t because, as I already said u dot is the first derivative of u with respect to time not with respect to space then u 2 dot will be a fixed times G 2 dot t which is $\frac{\partial^2 u}{\partial t^2}$.

Similarly, we can calculate the first derivative and second derivative of u with respect to space that means x . So, $u \text{ dash } t$ is equal to $F \text{ dash } x$ times Gt and $u \text{ double dash}$ is equal to $F2 \text{ dash } x$ times Gt . Now, we have these equations that $\text{del } 2u \text{ by del } t^2$ is equal to v_p^2 times $\text{del } 2u \text{ by del } x^2$. So, here if we will write what we get from previous two lines, then what we will write, we will get Fx on the left-hand side we have a Fx times $G^2 \text{ dot } t$ which is equal to v_p^2 times Fx , sorry, $F \text{ double dash } x$ times Gt .

Then next step we can write v_p^2 times $F \text{ 2 dash } x$ divided by Fx is equal to $G^2 \text{ dot } p$ divided by Gt . Now, we have already learned when we have a function, we can express one function by harmonic motion, then its acceleration is proportional to the displacement and also acceleration is proportional to the frequency of motion and direction of acceleration is opposite to the direction of displacement.

So, in this case, if we express Gt by harmonic motion, then we can write it same way. So, that is a reason what we have done here, we have taken Gt as...

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$$G(t) = C_1 \sin \omega t + C_2 \cos \omega t \quad G'(t)$$

$$\dot{G}(t) = \frac{d}{dt} [G(t)] = -\omega^2 [C_1 \sin \omega t + C_2 \cos \omega t]$$

$$\frac{\ddot{G}(t)}{G(t)} = -\omega^2$$

Elastic Rod of Finite Length Subjected to Longitudinal Vibration

- For longitudinal wave: $\frac{\partial^2 u}{\partial t^2} = v_{p1}^2 \frac{\partial^2 u}{\partial x^2}$ where, $v_{p1} = \sqrt{\frac{E}{\rho}}$
- General solution: $u = F(x)G(t)$
- Therefore, $\dot{u} = F(x)\dot{G}(t)$ and $\ddot{u} = F(x)\ddot{G}(t) = \frac{\partial^2 u}{\partial t^2}$
- Similarly, $u' = F'(x)G(t)$ and $u'' = F''(x)G(t) = \frac{\partial^2 u}{\partial x^2}$
- Since $\frac{\partial^2 u}{\partial t^2} = v_{p1}^2 \frac{\partial^2 u}{\partial x^2}$
- Therefore, $F(x)\ddot{G}(t) = v_{p1}^2 F''(x)G(t)$
- Or, $v_{p1}^2 \frac{F''(x)}{F(x)} = \frac{\ddot{G}(t)}{G(t)} = -\omega^2$




I can write here, we can take Gt as... Let us take C_1 sine omega t plus C_2 cosine omega t , then what we can write $G_2 t$ that means I am taking the second derivatives of Gt which will be equal to, here you can see I have stepped one, I have jumped one step, that is $G \dot{t}$, you can calculate first $G \dot{t}$ and then also you can calculate $G_2 \dot{t}$, or you can jump one step and directly write this also.


So now from here, if we will write $G_2 \dot{t}$ divided by Gt , what we are getting that is equal to minus omega square. So, the same thing I have written here that $G_2 \dot{t}$ divided by Gt is equal to minus omega square. Now, eventually, $G_2 \dot{t}$ by Gt is equal to v_{p1}^2 times F_2 dashed x by F_x , that is the reason this entire thing is equal to minus omega square.

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Elastic Rod of Finite Length Subjected to Longitudinal Vibration



- Assume: $G(t) = A_1 \cos \omega t + A_2 \sin \omega t$ ✓
- Assume: $F(x) = C_1 \cos \frac{\omega x}{v_{p1}} + C_2 \sin \frac{\omega x}{v_{p1}}$
- Boundary conditions:
 - at $x = 0, u = 0 \Rightarrow G(t)F(x) = 0$
 - Hence, at $x = 0 F(x) = 0$
 - at $x = L, \frac{\partial u}{\partial x} = 0$ (since no strain at the free end)



$$G(t) = C_1 \sin \omega t + C_2 \cos \omega t \quad \dot{G}(t)$$

$$\ddot{G}(t) = \frac{\partial^2}{\partial t^2} [G(t)] = -\omega^2 [C_1 \sin \omega t + C_2 \cos \omega t]$$

$$\frac{\ddot{G}(t)}{G(t)} = -\omega^2$$

Now, I am just making one change here, here I have taken C_1 , C_2 . Now, in this slide, you can say that two constant terms associated with the cosine ωt and sine ωt are taken as A_1 and A_2 respectively. So, with these only, we will continue our discussion. So, let us assume $G(t)$ here is equal to $A_1 \cos \omega t + A_2 \sin \omega t$.

We need to assume also something for F_x , here we have taken F_x is equal to C_1 times cosine ωx divided by it will be $v_p 1$ and plus C_2 times sine ωx divided by $v_p 1$, please make it $v_p 1$. Then we need to set boundary condition. Let us take a simple case. When a soil sample is resting on a pedestal, that means its bottom is fixed.

We are not applying anything at the top of this sample. And the length of the sample is A and its diameter is D . So here, what boundary condition we can set, when x , let us take x is measured in this direction. So, when x is equal to 0 , that means at the bottom end, what we can say displacement is 0 . So, first boundary condition, I can write that x is equal to 0 , u is equal to 0 , then u means $G(t)$ times F_x that is also equal to 0 .

So, at x is equal to 0 , u is equal to 0 . So, what we can write we can write $G(t)$ times F_x is equal to 0 . Now under any circumstances, $G(t)$ cannot be 0 , then what will be 0 , F_x . So let us write at x is equal to 0 , F_x is 0 . When F_x is 0 , what we can write that we will see.

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$$F(x) = C_1 \cos \frac{\omega x}{v_{p1}} + C_2 \sin \frac{\omega x}{v_{p1}}$$

at $x=0$ $F(x) = F(0) = C_1 = 0$

$$\Rightarrow C_1 = 0$$

$$F(x) = C_2 \sin \frac{\omega x}{v_{p1}}$$

at $x=L$ $\frac{\partial u}{\partial x} = 0$

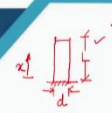
$$\frac{\partial u}{\partial x} = \left(\frac{\omega}{v_{p1}}\right) C_2 \cos \frac{\omega x}{v_{p1}}$$

$$\frac{\partial u}{\partial x} \Big|_{x=L} = \left(\frac{\omega}{v_{p1}}\right) C_2 \cos \frac{\omega L}{v_{p1}} = 0$$

$$\frac{\omega L}{v_{p1}} = (2n-1) \frac{\pi}{2} \quad \text{where } n = 1, 2, 3, \dots \text{ etc.}$$

For $n=1$ $v_{p1} = \frac{\omega L}{\pi/2} = \frac{2\omega L}{\pi}$
(Phase mode)

Elastic Rod of Finite Length Subjected to Longitudinal Vibration



- Assume: $G(t) = A_1 \cos \omega t + A_2 \sin \omega t$ ✓
- Assume: $F(x) = C_1 \cos \frac{\omega x}{v_{p1}} + C_2 \sin \frac{\omega x}{v_{p1}}$
- Boundary conditions:
 - at $x=0, u=0 \Rightarrow G(t)F(x) = 0$
 - Hence, at $x=0$ $F(x) = 0$
 - at $x=L, \frac{\partial u}{\partial x} = 0$ (since no strain at the free end)

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So first, this, I am going to a new page. So here we have taken $F(x)$ is equal to C_1 times cosine of ωx by v_{p1} plus C_2 times sine of ωx by v_{p1} , 1 stands for one dimension that it. Then when x is equal to 0 at x is equal to 0 $F(x)$ that is $F(0)$ is equal to C_1 and that we already defined that u is 0, so, $F(x)$ is 0 here. So, from this we can say C_1 is 0, then we can write $F(x)$ is equal to C_2 times sine of ωx divided by v_{p1} .

Now, see the second condition, second condition says that at x is equal to L that means at the top of these sample $\frac{\partial u}{\partial x}$ is equal to 0, since no strain at the free end. So, let us write this condition also at x is equal to L $\frac{\partial u}{\partial x}$ is equal to 0. So, first we need to find out what $\frac{\partial u}{\partial x}$ that is nothing but ω divided by v_{p1} times C_2 times cosine of ωx divided by v_{p1} .

Now here we will write in place of x , we will write L . So, $\frac{\partial u}{\partial x}$ at x is equal to L is ω divided by v_p times C_2 times cosine of ωL by v_p and that is equal to 0. If this is equal to 0, what does it mean? It means that the angle that is ωL divided by v_p is either $\frac{\pi}{2}$ or $2n$ minus 1 times $\frac{\pi}{2}$ where n is varying from one to any integer. So, let us write that then ωL divided by v_p I can write it v_p is equal to $2n$ minus 1 times $\frac{\pi}{2}$, where n is equal to 1, 2, 3 etcetera, it can be any integer.

So, with this now, what we can calculate, we from this we can calculate what will be v_p . So, here v_p will be how much? Let us take first mode of vibration that means n is equal to 1. So, v_p is equal to 1 that is first mod of vibration, v_p is equal to ωL divided by $\frac{\pi}{2}$. So, in this way, we can calculate the velocity of the longitudinal wave or P wave for first mode of vibration and also from these. So, this is our final. So, in this way we can calculate v_p . I can write one more step here just like this. Now come to the next.

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Elastic Rod of Finite Length Subjected to Longitudinal Vibration

- > at $x = 0 \Rightarrow C_1 = 0$
- > at $x = L, \frac{\partial u}{\partial x} = 0$
- > $\frac{\partial u}{\partial x} = G(t) \left[C_2 \frac{\omega}{v_p} \cos \frac{\omega L}{v_p} \right] = 0$
- > Since $G(t) \neq 0, \left[C_2 \frac{\omega}{v_p} \cos \frac{\omega L}{v_p} \right] = 0$
- > Thus, $\cos \frac{\omega L}{v_p} = 0$
- > $\frac{\omega L}{v_p} = (2n - 1) \frac{\pi}{2}$ where $n = 1, 2, 3, \dots$ etc
- > Hence for first mode of vibration i.e. $n = 1, \frac{\omega L}{v_p} = \frac{\pi}{2}$
- > So, $F(x) = C_2 \sin \frac{\pi x}{2L}$

Handwritten notes on the right side of the slide:

$$F(x) = C_2 \sin \frac{\omega x}{v_p}$$

$$v_p = \frac{\omega L}{\pi/2} \Rightarrow \frac{\omega}{v_p} = \frac{\pi}{2L}$$

$$F(x) = C_2 \sin \frac{\pi x}{2L}$$

What I said that is explained here you can go through this. So, finally, we can write F_x using the expression C_2 times in place of ωL , what we have I am just writing here we have considered actually F_x is equal to C_2 times sine of ωx divided by v_p and already we have established relationship between ω and v_p .

What was that relationship we have seen that v_p is equal to our ωL divided by $\frac{\pi}{2}$. So, from these what we can write in place of ω divided by v_p we can write $\frac{\pi}{2L}$. So, F_x is equal to C_2 times sine of $\frac{\pi x}{2L}$ when that is written here.

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Elastic Rod of Finite Length Subjected to Longitudinal Vibration

- Assume: $G(t) = A_1 \cos \omega t + A_2 \sin \omega t$
- Assume: $F(x) = C_1 \cos \frac{\omega x}{v_{p1}} + C_2 \sin \frac{\omega x}{v_{p1}}$
- Boundary conditions:
 - at $x = 0, u = 0 \Rightarrow C_1 = 0$
 - at $x = L, M \frac{\partial^2 u}{\partial t^2} = -AE \frac{\partial u}{\partial x}$

Handwritten notes on the right side of the slide:

- Diagram of a mass M on a rod of length L .
- Free-body diagram of mass M showing forces $F = Mg$ and F .
- Text: "at top i.e. $x=L$ "
- Equation: $\epsilon_{xx} = \frac{\partial u}{\partial x} = -\frac{F}{AE}$
- Equation: "For mass 'M', $M \frac{\partial^2 u}{\partial t^2} = F$ "
- Equation: $\epsilon_{xx} = \frac{\partial u}{\partial x} = -M \frac{\partial^2 u}{\partial t^2} \left(\frac{L}{AE} \right)$

Now, another boundary condition is possible, let us take the same soil sample resting on a pedestal fixed at the pedestal. Now, at the top what we can do let us take we have placed a mass in so, sorry. So, here a mass is placed at the top of this soil sample and mass is N capital N and the length of the soil sample and its diameter remains same L and d respectively. I can write I can show here L is the length and d is the diameter.

Now, what boundary condition we can write for these keys when x is equal to 0 that time u is 0 because $x=0$ means at the bottom it is fixed so, displace there is no displacement. So, from this we can write C_1 is 0. So, from the first boundary condition we can write C_1 is equal to 0. Now, second boundary condition what is happening at the top end.

So, at the top end a mass of magnitude capital M is placed. Now, if I will draw the Freebody diagram for this mass what it will be let us take these the weight of the mass we can write F is equal to capital Mg also here just, then what is happened to the soil sample, the soil sample, I am just showing I am just interested for the free top end.

So, at the top end it is also subjected to the same force capital F , the only difference is the direction of the force, if the for the mass M the direction of the forces downward for the soil the direction of the same force F under which it is subjected is in upward direction. So, for such case what we can write about the strain at top that is x is equal to L , how much is the strain, strain is $\frac{\partial u}{\partial x}$ or I can write ϵ_{xx} that is equal to F divided by AE .

In this case, what we have from the mass M if I will write the equation of motion what I can write for mass M for mass M what I can write, I can write mass times acceleration. So, mass times acceleration which is inertia force is equal to the unbalanced force we consider F positive in downward direction.

So, here F is positive if we consider F is positive in downward direction then for the other case, we get we need to consider minus F here and eventually when we deal with soil compressive forces positive and tensile forces negative. So, in that way the sign convention matches perfectly. So, now, what we will do here in place of F we can write mass times acceleration.

So, I can write ϵ_{xx} or $\frac{\partial u}{\partial x}$ is equal to minus M times $\frac{\partial^2 u}{\partial t^2}$ times 1 upon AE. So, from this we can write this expression that here I have multiplied both sides by minus AE, if I will multiply both sides by minus AE, what we can write mass times acceleration is equal to minus AE times $\frac{\partial u}{\partial x}$ that is our actual strain for this case. So, now using this boundary condition what we will see.

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Elastic Rod of Finite Length Subjected to Longitudinal Vibration

- at $x=0 \Rightarrow C_1 = 0$
- at $x=L, M \frac{\partial^2 u}{\partial t^2} = -AE \frac{\partial u}{\partial x}$

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So, alright. So, now, after setting the two boundary conditions what we can write?

(Refer Slide Time: 29:25)

The image shows a whiteboard with handwritten mathematical derivations. At the top, the wave equation is given as $\frac{\partial^2 u}{\partial t^2} = v_{p1}^2 \frac{\partial^2 u}{\partial x^2}$ and the displacement is assumed to be $u = F(x)G(t)$. Below this, the boundary condition at $x=L$ is stated as $M \frac{\partial^2 u}{\partial t^2} = -AE \frac{\partial u}{\partial x}$. The derivation then proceeds to find the acceleration $\ddot{u} = \frac{\partial^2 u}{\partial t^2} = F(x) \ddot{G}(t)$ and the strain $u' = \frac{\partial u}{\partial x} = F'(x) G(t)$. The function $F(x)$ is expressed as $F(x) = C_1 \cos \frac{\omega x}{v_{p1}} + C_2 \sin \frac{\omega x}{v_{p1}} = C_2 \sin \frac{\omega x}{v_{p1}}$. The derivative $F'(x) = \left(\frac{\omega}{v_{p1}}\right) (C_2 \cos \frac{\omega x}{v_{p1}})$ is then evaluated at $x=L$ to give $F'(x=L) = \left(\frac{\omega}{v_{p1}}\right) (C_2 \cos \frac{\omega L}{v_{p1}})$.

So, now we have this equation $\frac{\partial^2 u}{\partial t^2}$ that is equal to v_{p1}^2 times $\frac{\partial^2 u}{\partial x^2}$. Now, in this equation what we have assumed u is equal to function of space which is $F(x)$ times function of time, which is $G(t)$. Capital G represents the function of time and capital F represents the function of space.

What is the second boundary condition here at x is equal to L that means, at the top end or other end we have these boundary condition that mass times acceleration, which is inertia force is equal to minus AE times $\frac{\partial u}{\partial x}$. $\frac{\partial u}{\partial x}$ means longitudinal strain here. Now, since we have assumed that u is equal to $F(x)$ times $G(t)$. So, from this we can find out $\frac{\partial^2 u}{\partial t^2}$ which is the acceleration. So, I am once again writing it here.

So, what is acceleration then \ddot{u} which is $\frac{\partial^2 u}{\partial t^2}$ is equal to $F(x)$ times $\ddot{G}(t)$ and what is the expression for u' or we can write it as $\frac{\partial u}{\partial x}$ it is $F'(x)$ times $G(t)$. Now, already we have assumed the expression for $F(x)$ what is that $F(x)$ is equal to $C_1 \cos \frac{\omega x}{v_{p1}} + C_2 \sin \frac{\omega x}{v_{p1}}$.

So, here what is $F'(x)$, $F'(x)$ is. Now, in the expression of $F(x)$ we have already seen C_1 is equal to 0 from the first boundary condition. So, I can just write it as $C_2 \sin \frac{\omega x}{v_{p1}}$ then $F'(x)$ is equal to C_2 times better I write it first the ω divided by v_{p1} . So, $F'(x)$ is equal to ω divided by v_{p1} times C_2 times cosine of $\frac{\omega x}{v_{p1}}$.

So, now, I need to write F dashed at x is equal to L which is omega divided by vp 1 times C2 times cosine of omega L divided by vp 1. Now, we can use the expression of F dash L and that expression of this one. So, I can write it here itself. So, for these what I can do just go to the next page a new page.

(Refer Slide Time: 34:03)

$$\begin{aligned}
 M \frac{\partial^2 u}{\partial t^2} &= -AE \frac{\partial u}{\partial x} \quad \text{at } x=L \\
 \Rightarrow M [F(L) \ddot{G}(t)] &= -AE [F'(L) G(t)] \\
 \Rightarrow M \frac{\ddot{G}(t)}{G(t)} &= -AE \frac{F'(L)}{F(L)} \\
 \Rightarrow M (-\omega^2) &= -AE \frac{F'(L)}{F(L)} \\
 \Rightarrow M (-\omega^2) F(L) &= -AE F'(L) \\
 \Rightarrow M (-\omega^2) \left(C_2 \sin \frac{\omega L}{v_{p1}} \right) &= -AE \left(\frac{\omega}{v_{p1}} \right) \left(C_2 \cos \frac{\omega L}{v_{p1}} \right) \\
 \Rightarrow \tan \frac{\omega L}{v_{p1}} &= - \left(\frac{AE}{M} \right) \left(\frac{\omega}{v_{p1}} \right) \left(-\frac{1}{\omega^2} \right) \\
 &= \frac{AE}{M v_{p1} \omega}
 \end{aligned}$$

$$\begin{aligned}
 \frac{\partial^2 u}{\partial t^2} &= v_{p1}^2 \frac{\partial^2 u}{\partial x^2} \quad u = F(x)G(t) \\
 \text{at } x=L \quad M \frac{\partial^2 u}{\partial t^2} &= -AE \frac{\partial u}{\partial x} \\
 u &= F(x)G(t) \\
 \Rightarrow \ddot{u} &= \frac{\partial^2 u}{\partial t^2} = F(x) \ddot{G}(t) \\
 u' &= \frac{\partial u}{\partial x} = F'(x)G(t) \\
 F(x) &= C_1 \cos \frac{\omega x}{v_{p1}} + C_2 \sin \frac{\omega x}{v_{p1}} = C_2 \sin \frac{\omega x}{v_{p1}} \\
 F'(x) &= \left(\frac{\omega}{v_{p1}} \right) \left(C_2 \cos \frac{\omega x}{v_{p1}} \right) \\
 \underline{F'(x=L)} &= \left(\frac{\omega}{v_{p1}} \right) \left(C_2 \cos \frac{\omega L}{v_{p1}} \right)
 \end{aligned}$$

So, now, I can first write M times acceleration is equal to minus AE times del u of del x at x is equal to L. So, in the next line what we can write mass times acceleration, acceleration means del 2u of del t2 at x is equal to L, what I can write if L times G2 dot which is equal to minus AE times F dashed L times Gt or what I can write here mass times G2 dot divided by G is equal to minus AE times F dashed L divided by FL.

Now here, what is the value of G_2 dot divided by G , we know g is a harmonic function. So, already we have seen for harmonic function or harmonic motion, what we can write acceleration is proportional to or acceleration is equal to minus omega square times the equation of the motion, which is Gt here. So, I can actually write here M times minus omega square which is equal to minus AE times F dashed L divided by FL or I can write it also as M times minus omega square times function of L is equal to minus AE times F dashed L .

Now, here I can write the value of FL . So, M times minus omega square, what is the value of a FL . So, here, FL means I can write here itself, that is C_2 times sine omega L divided by v_{p1} . So, I can write it in the next page, then C_2 times sine omega L divided by v_{p1} in place of FL and that is equal to minus AE times F dash del, what is f dash del omega divided by v_{p1} times C_2 times cosine of omega L divided by v_{p1} . Now, from these you can see C_2 which is a nonzero quantity present in both sides.

So, I can cancel it out like this and I can rewrite the in this expression as what I can write I can write then omega divided by v_{p1} omega no I cannot write I can write omega divided by v_{p1} which is better I should write first the left hand side the right hand side. So, now, what I can write tangent of omega L divided by v_{p1} that is equal to minus AE divided by M this M is used here times, omega divided by v_{p1} times minus 1 divided by omega square.

So, simplifying it we can write AE divided by M times v_{p1} times omega. So, this is the value of tangent of omega L divided by v_{p1} .

(Refer Slide Time: 39:43)

The image shows a whiteboard with handwritten mathematical derivations. The equations are as follows:

$$\frac{\omega L}{v_{p1}} \tan \frac{\omega L}{v_{p1}} = \left(\frac{\omega L}{v_{p1}} \right) \left(\frac{AE}{M v_{p1} \omega} \right)$$

$$= \frac{AE}{M v_{p1}^2} = \frac{ALP}{M} = \frac{Ms}{M}$$

$$v_{p1}^2 = \frac{E}{\rho}$$

$$\Rightarrow \frac{E}{v_{p1}^2} = \rho$$

$$\frac{\omega L}{v_{p1}} \tan \frac{\omega L}{v_{p1}} = \frac{Ms}{M}$$

The whiteboard also features a vertical toolbar on the left side with various drawing tools like a pen, eraser, and highlighter. At the bottom, a Windows taskbar is visible with several application icons.

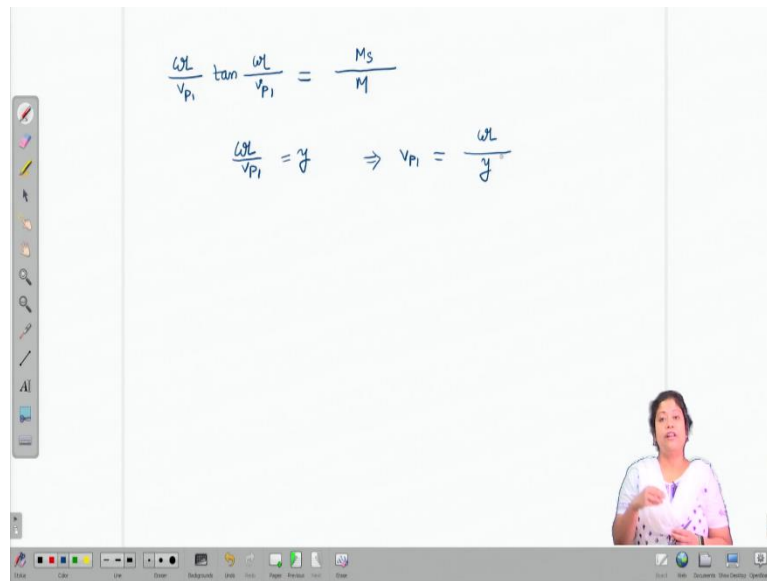
$$\begin{aligned}
M \frac{\partial^2 u}{\partial t^2} &= -AE \frac{\partial u}{\partial x} \quad \text{at } x=L \\
\Rightarrow M [F(L) \ddot{G}(t)] &= -AE [F'(L) G(t)] \\
\Rightarrow M \frac{\ddot{G}(t)}{G(t)} &= -AE \frac{F'(L)}{F(L)} \\
\Rightarrow M (-\omega^2) &= -AE \frac{F'(L)}{F(L)} \\
\Rightarrow M (-\omega^2) F(L) &= -AE F'(L) \\
\Rightarrow M (-\omega^2) \left(\cancel{C} \sin \frac{\omega L}{v_{p1}} \right) &= -AE \left(\frac{\omega}{v_{p1}} \right) \left(\cancel{C} \cos \frac{\omega L}{v_{p1}} \right) \\
\Rightarrow \tan \frac{\omega L}{v_{p1}} &= - \left(\frac{AE}{M} \right) \left(\frac{\omega}{v_{p1}} \right) \left(-\frac{1}{\omega^2} \right) \\
&= \frac{AE}{M v_{p1} \omega}
\end{aligned}$$

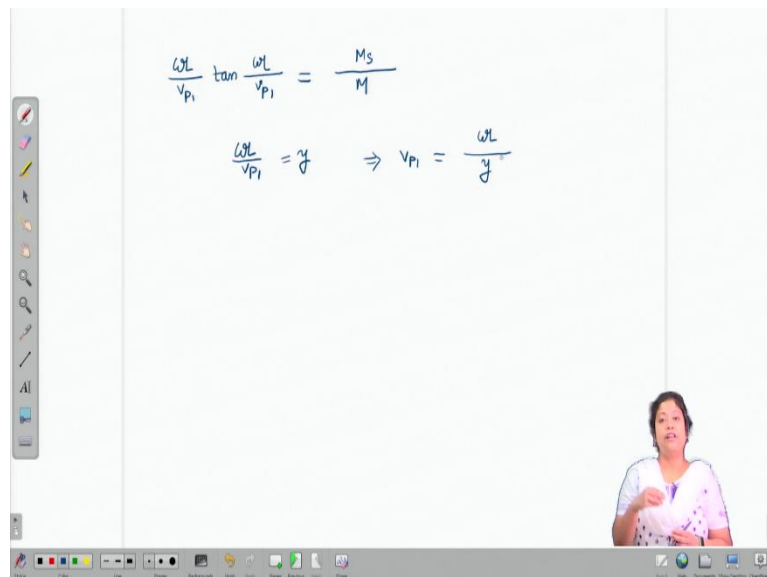
Now, what we can write here. Here if I will multiply omega L by vp 1 on both sides what we will get? Let us see, omega L divided by vp 1 times tangent of omega L divided by vp 1 which is equal to omega L divided by vp 1, times what we have on the right-hand side earlier AE divided by M times vp 1 times omega. So, I can write here AE divided by M times vp 1 times omega. So, finally, what we are getting here finally, I am getting AE or I can AE L I can write L before writing AE.

So, A L E divided by M times vp 1 square. So, what is AL here, AL is the volume of the rod or the soil sample when we are taking soil. So, AL is volume and what else we can what is vp 1 square, we have already seen vp 1 square means E divided by rho where E is modulus of elasticity and rho is the density of the material. So, from these what I can write, I can then write E divided by vp 1 square is equal to rho. So, you can see here that term E divided by vp squared is present.

So, I can then write this expression as aAL divided by M times rho. Now, AL which is the volume of the rod or sample times rho gives the mass of that sample or mass of the rod. So, I can for soil sample it is mass of the soil divided by M which is mass of the which is the mass kept at the top of this soil sample. So, in this way finally, what I can write omega L divided by vp 1 times tangent of omega L divided by vp 1 is equal to Ms divided by M.

(Refer Slide Time: 43:03)


$$\frac{\omega L}{v_{p1}} \tan \frac{\omega L}{v_{p1}} = \frac{M_s}{M}$$
$$\frac{\omega L}{v_{p1}} = y \Rightarrow v_{p1} = \frac{\omega L}{y}$$


$$\frac{\omega L}{v_{p1}} \tan \frac{\omega L}{v_{p1}} = \frac{M_s}{M}$$
$$\frac{\omega L}{v_{p1}} = y \Rightarrow v_{p1} = \frac{\omega L}{y}$$

We get I am writing here $\frac{\omega L}{v_{p1}} \tan \frac{\omega L}{v_{p1}}$ is equal to $\frac{M_s}{M}$, where M_s is the mass of the soil sample and M is the mass which is placed at the top of the soil sample. Now, we know how to solve this type of equation. So, we can solve this type of equation and find out something for $\frac{\omega L}{v_{p1}}$ let us take that something is y then from that we can express v_{p1} as $\frac{\omega L}{y}$.

So, in this way, we can find out the velocity of the P wave propagating through one propagating through soil element.

(Refer Slide Time: 44:13)

Elastic Rod of Finite Length Subjected to Longitudinal Vibration

- at $x = 0 \Rightarrow C_1 = 0$
- at $x = L, M \frac{\partial^2 u}{\partial t^2} = -AE \frac{\partial u}{\partial x}$

➤ $\tan \frac{\omega L}{v_{p1}} = \frac{AE}{M v_{p1} \omega}$

➤ Finally, $\frac{\omega L}{v_{p1}} \tan \frac{\omega L}{v_{p1}} = \frac{M_1}{M}$ where $n = 1, 2, 3, \dots$ etc.

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In this case, what we have considered is that we have considered one directional wave propagation.