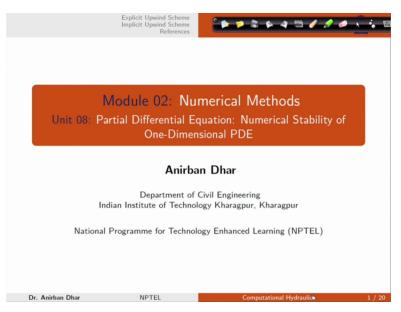
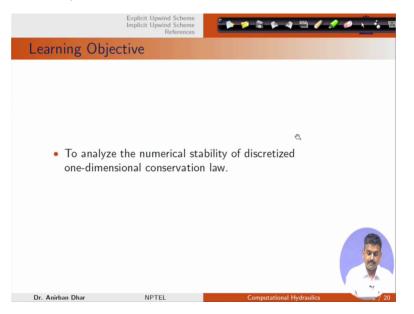
Computational Hydraulics ProfessorAnirban Dhar Department of Civil Engineering Indian Institute of Technology Kharagpur Lecture 12 Partial Differential Equation - Numeric Stability of One Dimensional PDE

Welcome to this lecture number 12 of the course computational hydraulics. We are in module 2, numerical methods. And in today's lecture we will becovering unit 8, partial differential equation, numerical stability of one dimensional PDE.

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What is the learning objective for this particular unit? At the end of this unit students will be able to analyze the numerical stability of discretized one-dimensional conservation law in terms of partial differential equation. (Refer Slide Time 01:02)



Let us consider our general equation with general variable phi. We have seen that initially with this first term in the right hand side we have solved boundary value problem. With our temporal term and boundary value problem or right hand term, this one and this one we have solved initial boundary value problem. And in these boundary value problem or initial boundary value problems, we have considered Sphi or source sink term.

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General Equation A form of differential equation with a general variable ϕ : $\frac{\partial(\Lambda_{\phi}\phi)}{\partial t} + \nabla.(\Upsilon_{\phi}\phi\mathbf{u}) = \underbrace{\nabla.(\Gamma_{\phi},\nabla\phi)}_{F_{\phi_o}} + S_{\phi}$ (1) where $\phi = \text{general variable}$ $\Lambda_{\phi}, \Upsilon_{\phi} = \text{problem dependent parameters}$ $\Gamma_{\phi} = \text{tensor}$ $F_{\phi_o} = \text{other forces}$ $S_{\phi} = \text{source/sink term}$		Explicit Upwind Scheme Implicit Upwind Scheme References	*****	/ / /
$\frac{\partial(\Lambda_{\phi}\phi)}{\partial t} + \nabla.(\Upsilon_{\phi}\phi\mathbf{u}) = \nabla.(\Gamma_{\phi}.\nabla\phi) + F_{\phi_{o}} + S_{\phi} $ (1) where $\phi = \text{general variable}$ $\Lambda_{\phi}, \ \Upsilon_{\phi} = \text{problem dependent parameters} \qquad IBVP$ $\Gamma_{\phi} = \text{tensor}$ $F_{\phi_{o}} = \text{other forces}$	General Equ	ation		
	where $\phi = \text{general var}$ $\Lambda_{\phi}, \ \Upsilon_{\phi} = \text{prob}$ $\Gamma_{\phi} = \text{tensor}$ $F_{\phi_o} = \text{other for}$	$\frac{\partial (\Lambda_{\phi}\phi)}{\partial t} + \nabla .(\Upsilon_{\phi}\phi \mathbf{u}) =$ iable lem dependent parameters rccs	$\nabla (\Gamma_{\phi}, \nabla \phi) + F_{\phi_o} + S_{\phi}$	(1)
Dr. Anirban Dhar NPTEL Computational Hydraulics	Dr. Asidore Diver	NOTE	Consudered Hadronick	

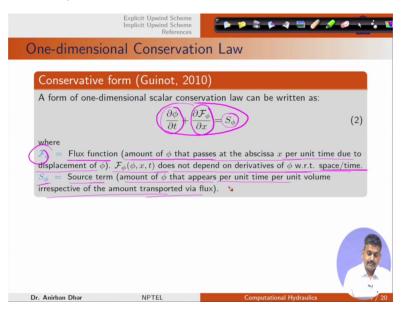
Now we need to consider the second term in the left hand side. This is somewhat related to velocity term. And we should analyze this term before starting our finite volume approach.

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	Explicit Upwind Scheme Implicit Upwind Scheme References		🖉 🥔 🔥 🐱
General Equat	ion		
where $\phi = { m general variable}$	e dependent parameters	$\nabla \cdot (\Gamma_{\phi} \cdot \nabla \phi) + F_{\phi_o} + S_{\phi}$	(1)
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So in this case, one-dimensional conservation law in terms of phi can be written as del phi by del t and del F phidel x Sphi. S phi is source sink term. This is temporal term and this is spatial term only in one dimension that is x. So what is this F phi? Fphi is the flux function. Amount of phi that passes at the abscissa x per unit time due to displacement of phi. And Fphi does not depend on derivative of phi with respect to space or time. And S phi is source sinkterm. This is amount of phi that appears per unit volume irrespective of the amount transported viaflux.

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So with this information we can proceed. We can see that we have phi which is function of x and t only. So one-dimensional in space and we have one time dimension.

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	Explicit Upwind Scheme Implicit Upwind Scheme References	Ê 🕨 🌮 🖉	\$ 4 = 1 8 6	• • • •
One-dim	ensional Conservati	on Law		
Conserv	vative form (Guinot, 201	10)		
where $\mathcal{F}_{\phi} = FI$ displacement $S_{\phi} = Sc$	Fone-dimensional scalar conset $\begin{array}{c} \partial \phi \\ \partial t \end{pmatrix} + \frac{\partial}{\partial t} \\ \\ \psi $	$rac{\partial \mathcal{F}_{\phi}}{\partial x}=S_{\phi}$ asses at the absorber of derivat pears per unit times the second seco	$\phi(x,t)$, cissa x per unit time du ives of ϕ w.r.t. space/t	
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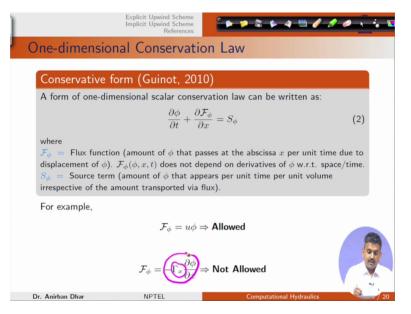
So for example, in this case phi can be u phi. If phi is equal to u phi and lambda phi equals to 1 and our upsilonphi is equal to 1, then we can get this equation without the right hand first term and other force term on the right hand side.

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	Explicit Upwind Scheme Implicit Upwind Scheme References	`► ♥≈►4 ≡/ /	🥔 👗 🍾 🛚
One-dimensiona	al Conservation	on Law	
Conservative for	rm (Guinot, 201	.0)	
A form of one-dime	nsional scalar conser	rvation law can be written as:	
	$\sqrt{\frac{\partial \phi}{\partial t}} + \frac{\partial \phi}{\partial t}$	$rac{\mathcal{F}_{\phi}}{\partial x} = S_{\phi}$	(2)
displacement of ϕ). \mathcal{F}	$F_{\phi}(\phi,x,t)$ does not deamount of ϕ that app	asses at the abscissa x per unit time epend on derivatives of ϕ w.r.t. space bears per unit time per unit volume flux).	
For example,	$\mathcal{F}_{\phi} = u \phi$	Allowed Al	=
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However any derivative in this Fphi is not allowed. So this is related to diffusion. So this gamma x is like diffusion coefficient. So this is not allowed for this kind ofequation because we have derivative of phi with respect to space or in one dimensional case.

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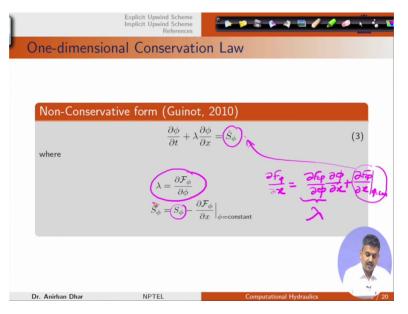
Sowe can write the non-conservative form where del phi by del t. And this is actually lambda into delphi by del x. This is written in terms of ourgeneral variable phi, not in terms of flux Fphi.

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]	One-dimensiona	Explicit Upwind Scheme mplicit Upwind Scheme References	n Law	<i>3 0</i> x i. v
	Non-Conservative	e form (Guinot, $\left(\frac{\partial \phi}{\partial \phi}\right)_{+} \lambda \frac{\partial}{\partial \phi}$		(3)
		∂t	x x	(-)
			×	
	Dr. Anirban Dhar	NPTEL	Computational Hydraulics	20

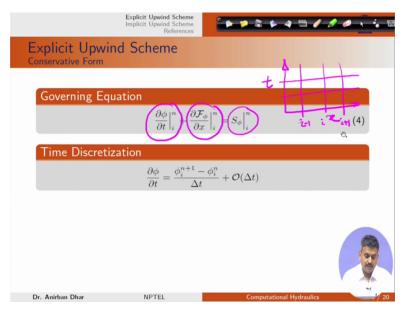
So what is this lambda? Lambda is equal to del Fby del phi. And we have this modified source sink term. This is Sphiminusdel phi. This is constant. Soactually we can write our Fphi del y del x as delphi. This phi into delphi by del x plus del F phi by del x where phi is constant value. So now this term is lambda and we can change this side for this term. And if we transfer this term on the right hand side and subtract from the original source sink term then we can get this modified one.

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So let us utilize our explicit upwind scheme. What is this upwind? Depending on this value of lambda or sign of lambda, we generally change the scheme. If it is positive or negative we need to change our discretization in space. So we have our original conservative form equation. We are discretizing it at I n level. This is present time level. So if we consider in this case we have one space access and one time access other level for this one. In this case if we have I,Iminus1, I, Iplus 1.

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Then for timediscretization we have nth level n and n plus 1. For space discretization we have Iminus 1,I, Iplus 1. These values are actually unknown values.

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	Explicit Upwind Scheme Implicit Upwind Scheme References	* * * * * * *	=////
Explicit Upwin Conservative Form	d Scheme		J. Je-je
		nH	111
Governing Equ	ation		
	$\left. \frac{\partial \phi}{\partial t} \right _i^n + \frac{\partial \mathcal{F}}{\partial x}$	$\left. \frac{\phi}{c} \right _{i}^{n} = S_{\phi} \right _{i}^{n}$	(4)
Time Discretiz	ation		
	$\frac{\partial \phi}{\partial t} = \frac{\phi_i^{n+1} - \phi_i^{n+1}}{\Delta t}$	$\frac{-\phi_i^n}{2} + \mathcal{O}(\Delta t)$	
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And these are our known values.

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	Explicit Upwind Scheme Implicit Upwind Scheme References	* * * * * * *	= <i>/ / o</i> 🛴 🖥 🖬
Explicit Upwind	Scheme		1 de ale
		nt	
Governing Equat	ion		
	$\left.\frac{\partial\phi}{\partial t}\right _{i}^{n} + \frac{\partial\mathcal{F}}{\partial x}$	$\left. \frac{\phi}{c} \right _{i}^{n} = S_{\phi} \right _{i}^{n}$	(4)
Time Discretizat	ion		
	$\frac{\partial \phi}{\partial t} = \frac{\phi_i^{n+1} - \phi_i^{n+1}}{\Delta t}$	$\frac{-\phi_i^n}{2} + \mathcal{O}(\Delta t)$	
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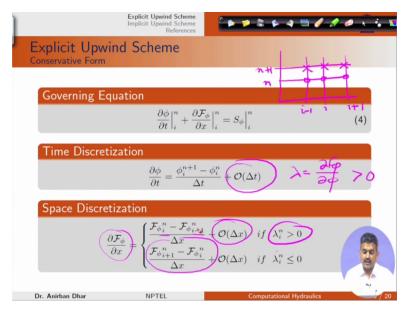
So with this information we can discretize our space derivative for flux. So interestingly in this case if lambda I n is positive, that means lambda equals to flux. So change in flux with respect tovariable phithat is positive, then we will be utilizing this derivative I, I minus 1 for delphi by del x.

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		Explicit Upwind Scheme Implicit Upwind Scheme References	► ♥ \$\$ \$ 4 \$\$ \$ \$ \$	e 🔪 is 🛛
1	Explicit Upwine Conservative Form	d Scheme		
			nt	
	Governing Equa	ation		
		$\left. \frac{\partial \phi}{\partial t} \right _i^n + \frac{\partial \mathcal{F}}{\partial t}$	$\left[\frac{\overline{\phi}}{v}\right]_{i}^{n} = S_{\phi}\Big _{i}^{n}$	(4)
	Time Discretiza	ation		
		$\frac{\partial \phi}{\partial t} = \frac{\phi_i^{n+1}}{\Delta t}$	$\frac{-\phi_i^n}{2} + \mathcal{O}(\Delta t)$ $\lambda = \frac{\partial f \phi}{\partial \phi}$	70
	Space Discretiz	ation		
	$\frac{\partial F}{\partial x}$	$ \begin{cases} \frac{\mathcal{F}_{\phi_i}^n - \mathcal{F}_{\phi_{i-1}}^n}{\Delta x} \\ \frac{\mathcal{F}_{\phi_{i+1}}^n - \mathcal{F}_{\phi_i}^n}{\Delta x} \end{cases} $	$\begin{aligned} &+ \mathcal{O}(\Delta x) if \lambda_i^n > 0 \\ &+ \mathcal{O}(\Delta x) if \lambda_i^n \le 0 \end{aligned}$	
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And if lambda is negative then we will be utilizing F Iplus 1 minus I value for this one. And in this case time discretization is of order delta t and space discretization is of order delta x. So overall accuracy for this one is delta x delta t.

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And for source sink term this is simple, we will be evaluating at I n level. Now we can write our final solution like this. If we change the sides for known variables on the right hand side. This is unknown level. (Refer Slide Time 10:50)

			Explicit Upwind Schem Implicit Upwind Schem Reference	e 🛁	Þ 🌮 🐍	\$ 4	- 11 /	e .	* 6 🛛
Exp	licit	Upwin	d Scheme						
S	ource	Term							
50	Jurce	Term		- P					
			S_{c}	$\phi = S_{\phi_i}^n$					
Fir	nal solu	ition can b	oe written as,						
	($\phi_i^{n+1} = \cdot$	$\begin{cases} \phi_i^n + \frac{\Delta t}{\Delta x} (\mathcal{F}_{\phi_{i-}}^n) \\ \phi_i^n + \frac{\Delta t}{\Delta x} (\mathcal{F}_{\phi_i}^n) \end{cases}$	$(-\mathcal{F}_{\phi_{i}}^{n})$	$+ \Delta t S_{\phi_i}^n$	if	$\lambda_i^n > 0$		
		- Chi	$\left(\phi_i^n + \frac{\Delta t}{\Delta x} (\mathcal{F}_{\phi_i})^n + \Delta $	$-\mathcal{F}_{\phi_{i+1}^n})$	$+\Delta t S_{\phi_i}^n$	if	$\lambda_i^n \le 0$		
									27
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This is known level. This flux is also known. These values are known. This is also known. So we can write like this.

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	Explicit Upwind Scheme Implicit Upwind Scheme References	ê 🗭 📁 🗿 🖗	4 🛎 🖉 🖋	🥔 🔪 🐍 🔽
Explicit Upwin	d Scheme			
Source Term				
Source Ferm	a	<i>a n</i>		
	$S_{\phi} =$	S_{ϕ_i}		
Final solution can l	pe written as,			
$\phi_i^{n+1} =$	$\begin{cases} \overbrace{\phi_i^n}^n + \frac{\Delta t}{\Delta x} (F_{\phi_{i-1}}^n) - \\ \phi_i^n + \frac{\Delta t}{\Delta x} (F_{\phi_i}^n - J) \end{cases}$	$\mathcal{F}_{\phi_i}^{\ \ n} + \Delta t \widehat{\mathcal{S}_{\phi_i}^{\ \ n}}$	$if \ \lambda_i^n > 0$	
	$\left(\phi_i^n + \frac{\Delta t}{\Delta x} (\mathcal{F}_{\phi_i^n} - \mathcal{F}_{\phi_i^n})\right)$	$F_{\phi_{i+1}}^n) + \Delta t S_{\phi_i}^n$	$if \ \lambda_i^n \le 0$	
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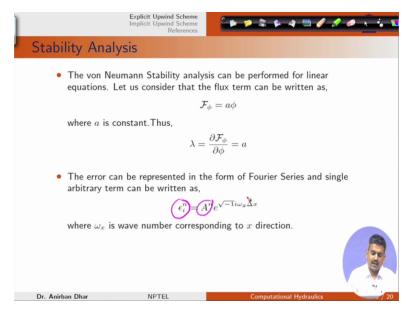
Now we need to analyze the stability of this particular scheme. But Von Neumann stability analysis can be performed for linear equations. Let us consider that the flux term can be written in terms of F phi where F phi equals to Aphi and A is some constant value. So lambda is delF phi by del phiwhich is A. Depending on the sign of the A, whether it is positive or negative, we can change the space discretization.

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	Explicit Upwind Scheme Implicit Upwind Scheme References	(*************************************	
Stability Anal	ysis		
equations. L		ysis can be performed for linear the flux term can be written as, $F_{\phi} = a\phi$	
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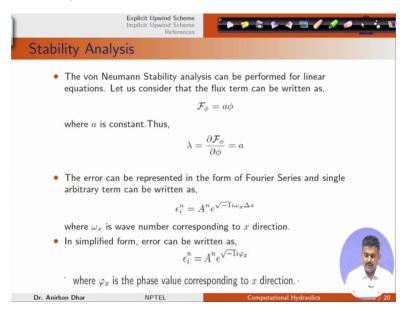
Now let us consider that error equation again it can be error term. Epsilon I n can be written in the terms of amplitude and this face thing forx direction.

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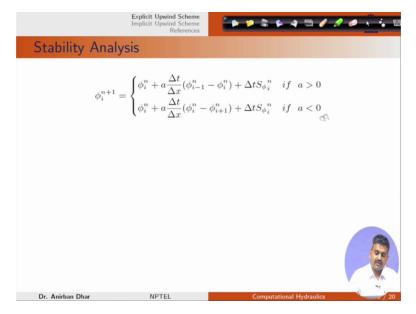
So we have this I n and A n I var phi x corresponding to that direction.

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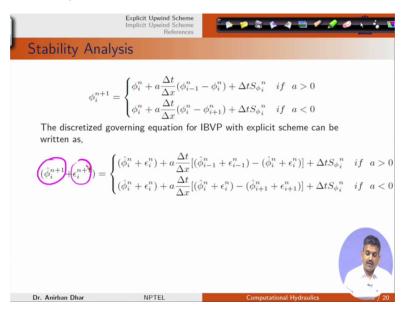


Now with this information we can proceed for stability analysis. Now we can bring that A term out of this bracket and we can write like this for A greater than zero and less than zero.

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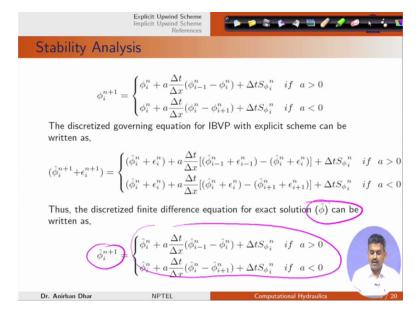


Now discretize governing equation for this initial boundary value problem. Butin this case boundary initial condition is a main thing because one side only it is defined. So it is like initial value problem. Now this is the solution which is coming from infinite precision computer and this is the amount of error involved in the particular step and particular discretization level. (Refer Slide Time 13:41)



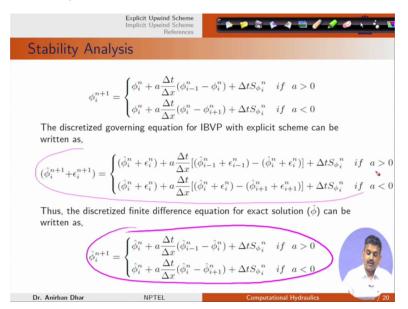
With this we can also write the discretized finite difference equation for exact solution. This is the exact solution think and exact solution should satisfy our original finite difference discretization. So we have written phi hat I n plus 1 and others are at known time level.

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Now if we subtract our exact solution equation discretized form, from this one, we can get the corresponding error equation.

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So error equation can be returned like this where this one is unknown.

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	Explicit Upwind Scheme Implicit Upwind Scheme References	► ■ ≈ ► 4 = Ø	🖉 🥔 🔪 🍬 🛯
Stability	Analysis		
Thus the	error equation can be represent		
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And these are known level things.

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	Explicit Upwind Scheme Implicit Upwind Scheme References	<mark>← > ≈ + + = /</mark> .	s 🖉 🔬 🕯 🛛
Stability Ana	lysis		
	equation can be represent $ \underbrace{\epsilon_{i}^{n+1}}_{i} = \begin{cases} \epsilon_{i}^{n} + a \frac{\Delta t}{\Delta x} (a + b) \\ \epsilon_{i}^{n} + a \frac{\Delta t}{\Delta x} (a + b) \end{cases} $		
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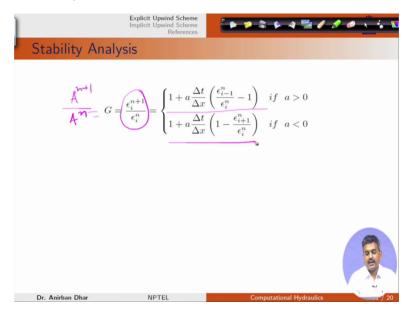
Now as per our definition we can write this epsilonI n plus 1 where only change in the amplitude, no change in this index for I. I n this is A and I varphi x minus 1, only change in index for I. I plus 1 only change in the index for I.

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		cit Upwind Scheme cit Upwind Scheme References		* • • =	1130	• • • •
Stability	Analysis					
Thus the	error equation $\epsilon_i^{n+1} =$	can be represent $\begin{cases} \epsilon_i^n + a \frac{\Delta t}{\Delta x} (\\ \epsilon_i^n + a \frac{\Delta t}{\Delta x} (\\ \end{cases}$		f a > 0 $f a < 0$		
		$ \begin{array}{c} \overbrace{\epsilon_{i}^{n+1}}=A^{n} \\ \hline \epsilon_{i}^{n}=A^{n} \\ \hline \epsilon_{i-1}^{n}=A^{n} \end{array} $	$e^{\sqrt{-1}i\varphi_x}$ $e^{\sqrt{-1}i\varphi_x}$ $e^{\sqrt{-1}(i-1)\varphi_x}$			
		$\underline{\epsilon_{i+1}^n} = A^n$	$e^{\sqrt{-1}(i+1)\varphi_x}$	6		
Dr. Anirban Dha	ar	NPTEL	Cor	nputational Hydr	aulics	20

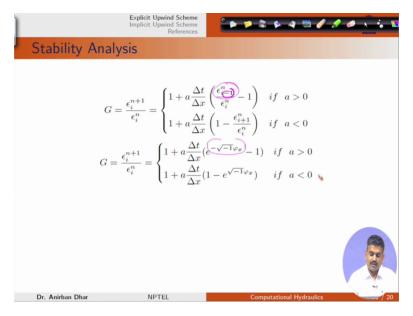
So with this expressions if we write our growth factor by dividing epsilon on both sides. This is epsilon I n plus 1 divided by epsilon I n and interestingly this is equal to A n plus 1 divided by A n. Now for A positive we have this form, A negative we have the second equation.

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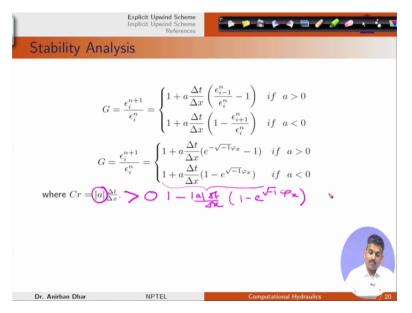
Now if you further simplify this one, we can write like this where this e to the power minus part is there. So minus I the minus this imaginary number into var phi x. And this is 1 plus 1 minus e to the power, imaginary number var phi x again for negative value.

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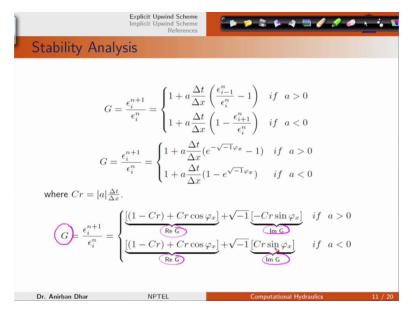
Now in this casewe can define a number which is Cr. And Cr is defined like this. Modulus of A into delta t divided by delta x. So always Cr is positive. Now the problem with this one is that in this case if A is negative we can write this equation as, 1 minus modulus of A into deltat divided by delta x 1 minus e to the power minus 1 var phi x. Only for negative value.

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Now with this Cr we can proceed and we can define our amplification factor in terms of real and imaginary part of G in simplified form.

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Now as per our Von Neumann stability criteria, this modulus of G should be less than equals to 1.

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	Explicit Upwind Scheme Implicit Upwind Scheme References		<u>i</u> 🔽
Stability Analy	ysis		
	$\frac{\epsilon_i^{n+1}}{\epsilon_i^n} = \begin{cases} 1 + a\frac{\Delta t}{\Delta x}(1 + a\frac{\Delta t}{\Delta x}) \\ 1 + a\frac{\Delta t}{\Delta x}(1 + a\frac{\Delta t}{\Delta x}) \end{cases}$	$ \begin{pmatrix} \frac{\epsilon_{i-1}^n}{\epsilon_i^n} - 1 \end{pmatrix} if a > 0 \\ \left(1 - \frac{\epsilon_{i+1}^n}{\epsilon_i^n}\right) if a < 0 \\ e^{-\sqrt{-1}\varphi_x} - 1) if a > 0 \\ 1 - e^{\sqrt{-1}\varphi_x}) if a < 0 $	
		$\begin{aligned} \varphi_x \\ \varphi_x \\ + \sqrt{-1} \underbrace{\left[-Cr \sin \varphi_x \right]}_{\text{Im G}} & if a > 0 \\ \varphi_x \\ + \sqrt{-1} \underbrace{\left[Cr \sin \varphi_x \right]}_{\text{Im G}} & if a < 0 \\ \ell \\ \varphi_x \\ \end{bmatrix} \end{aligned}$	
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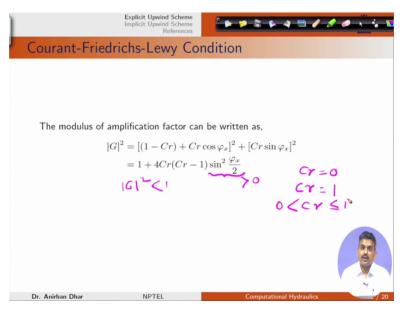
Now we need to evaluate the modulus of G. So if we evaluate the modulus of G in this case so basically we are taking square of real and imaginary part. Now this comes to this simplified form, 1 plus 4 Cr, Crminus 1 sin square phi x by 2.

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		Explicit Upwind Scheme Implicit Upwind Scheme References		🖉 🥔 🕺 🙀
1	Courant-Fried	Irichs-Lewy Co	ndition	
	The modulus of a	mplification factor car	n be written as,	
	1		$\cos\varphi_x]^2 + [Cr\sin\varphi_x]^2$	
		= 1 + 4Cr(Cr - 1)	$1)\sin^2\frac{\varphi_x}{2}$	
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In this caseimportant point is that this quantity is always positive. Now as per our condition G Mod square should be less than equal to 1. If Cr equals to zero then we have 1. If Cr equals to 1, againthis is equals to 1. And for values less than1 and greater than zero. For Cr we can get G modulus square less than equal to 1. So we can say that for this condition where Cr is greater than zero, less than equal to 1. This scheme is stable. This is known as Courant Friedrich Lewy condition or CFL condition.

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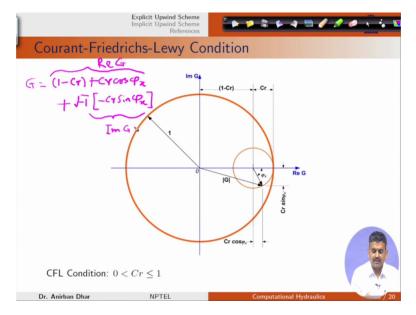
Thus we can say that explicit scheme is conditionally stable. And this is the actual condition.

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		Explicit Upwind Scheme Implicit Upwind Scheme References	**************************************	2 <u>)</u>		
1	Courant-Fried	richs-Lewy Co	ndition			
	The modulus of a	mplification factor car	n be written as,			
	0		$\cos\varphi_x]^2 + [Cr\sin\varphi_x]^2$			
		= 1 + 4Cr(Cr - 1)	$1)\sin^2\frac{\varphi_x}{2}$			
	CFL Condition: $0 < Cr \leq 1$ Explicit Scheme is Conditionally Stable.					
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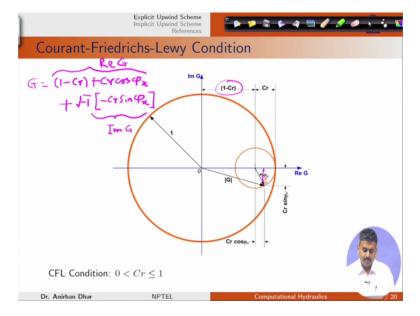
Now from geometrical point of view if we see, we can represent our G where G is equal to 1 minus Cr plus Cr cos var phi x plus square root minus 1 and within bracket minus Cr sin var phi x, this quantity. So this part is represented in the real axis or ReG. And this part is represented in the imaginary axis.

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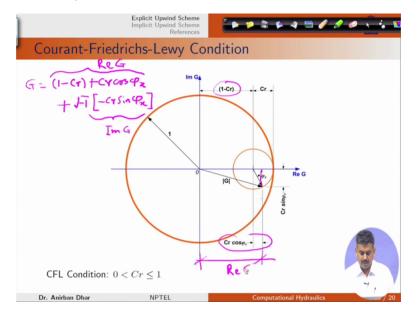
So this part is 1 minus Cr and this one is the projected one.

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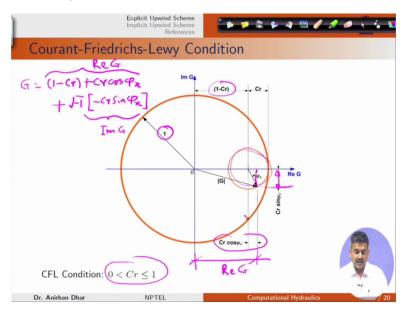
So this length is Cr into cos phi. So total length, this part is actually our real of G.

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Now and in this case, this part is imaginary component of G. So this thing, this modulus of G should be within this circle with unit radius for stability. So we can say that in this case if Cr value is within 0 to 1, this condition will be satisfied. That means the circle will be within this bigger circle. So our stability criterion will be satisfied.

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Now if we consider implicit upwind equation, we are discretizing it at n plus 1 level.

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	Explicit Upwind Scheme Implicit Upwind Scheme References	`►▼≈►∢≡ ∥	🖉 🥔 🛴 🐜 🛯
Implicit Upwind	d Scheme		
Governing Equa			
	$\frac{\partial \phi}{\partial t} \Big _{i}^{n+1} + \frac{\partial \mathcal{F}_{\phi}}{\partial x}$	$\int_{i}^{n+1} = S_{\phi} \Big _{i}^{n+1}$	(5)
Dr. Anirban Dhar	NPTEL	Computational Hydraulics	14 / 20

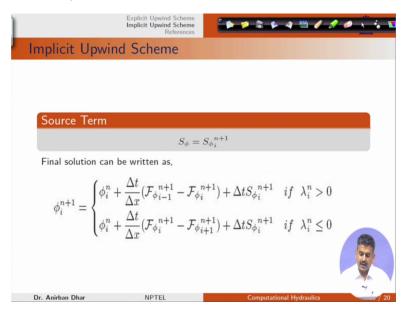
Nowif we consider our time discretization, this is our time discretization. We have truncation error of orderdel t. For space discretization again values are at nplus 1 level. For lambda I n plus 1 greater than zero, lambda I n plus 1 less than equal to zero. This is again of order delta x. Overall accuracy of the scheme is delta x delta t.

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	Explicit Upwind Scheme Implicit Upwind Scheme References	-	4 🛎 🦉 🖋	🥔 🔪 i. 🛚
Implicit Upwir	nd Scheme			
Governing Equ	uation			
	$\left. \frac{\partial \phi}{\partial t} \right _{i}^{n+1} + \frac{\partial \mathcal{F}_{q}}{\partial x}$	$\left \int_{i}^{n+1} = S_{\phi} \right _{i}^{n+1}$		(5)
Time Discretiz	zation			
	$\frac{\partial \phi}{\partial t} = \frac{\phi_i^{n+1}}{\Delta t}$	$\frac{-\phi_i^n}{2} + \mathcal{O}(\Delta t)$	0 (32)	\$
Space Discreti	zation			
$rac{\partial \mathcal{F}_{\phi}}{\partial x}$	$= \begin{cases} \frac{\mathcal{F}_{\phi_i}^{n+1} - \mathcal{F}_{\phi_i}^{n+1}}{\Delta x} \\ \frac{\mathcal{F}_{\phi_{i+1}}^{n+1} - \mathcal{F}_{\phi_i}^{n+1}}{\Delta x} \end{cases}$	$ \begin{array}{c} \begin{array}{c} & \\ + & \mathcal{O}(\Delta x) \end{array} & if \left(\lambda \right) \\ - & + & \mathcal{O}(\Delta x) \end{array} & if \left(\lambda \right) \end{array} $	$n+1 > 0$ $n+1 \le 0$	
				~ ~ ~
Dr. Anirban Dhar	NPTEL	Computati	onal Hydraulics	/ 20

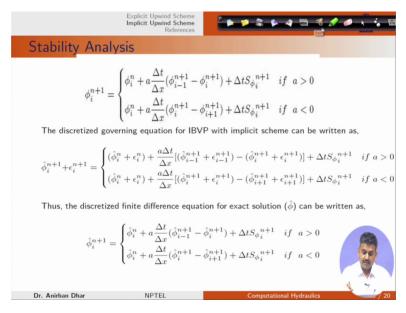
This is for source term and this source term is discretized at n plus 1 level. Now with this we can get the final solution which is in terms of values at nth level. And flux terms at nplus 1 level, source sink terms are source sink terms at n plus 1 level. This should be evaluated at nplus 1 level or at the nth level.

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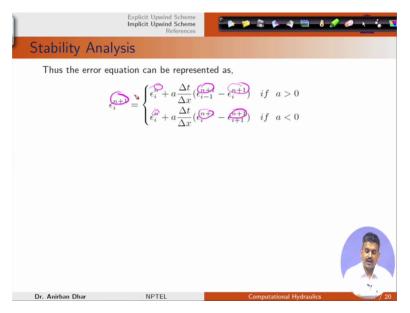
So stability analysis, again if we consider this, we need to see our final solution. We can again discretize the governing equation with a implicit scheme. And in this casewe can write the exact discrete equation for the finite difference solutions. So in this casewe are getting the equation in terms of errors.

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For epsilon I n plus 1 and epsilon I minus 1, n plus 1. All these are at nplus 1 levels. Only the first term is at nth level.

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So we can write our (api) amplification factor by dividing it with epsilon I. Again defining the Courant number like this where this is always positive. So we can write G like this.

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			cit Upwind Scheme cit Upwind Scheme References	î 🕨 📁 🕸	\$ 4 🛎 🖉 s	/ / x i. 🛛
St	ability	Analysis				
	Thus the	error equation	can be represer	ited as,		
		$\epsilon_i^{n+1} =$	$\begin{cases} \epsilon_i^n + a \frac{\Delta t}{\Delta x} (\epsilon_i^n + a \frac{\Delta t}{\Delta x}) \\ \epsilon_i^n + a \frac{\Delta t}{\Delta x} (\epsilon_i^n + a \frac{\Delta t}{\Delta x}) \end{cases}$	$e^{i+1}_{-1} - \epsilon^{n+1}_i$ if $e^{i+1}_{i+1} - \epsilon^{n+1}_{i+1}$ if	a > 0 a < 0	
($Cr = a _{\frac{1}{2}}$	Δt	•	$ \begin{split} & \frac{\epsilon_{i-1}^{n+1}}{\epsilon_i^n} - \frac{\epsilon_i^{n+1}}{\epsilon_i^n} \\ & \frac{\epsilon_i^{n+1}}{\epsilon_i^n} - \frac{\epsilon_{i+1}^{n+1}}{\epsilon_i^n} \\ \end{split} $		
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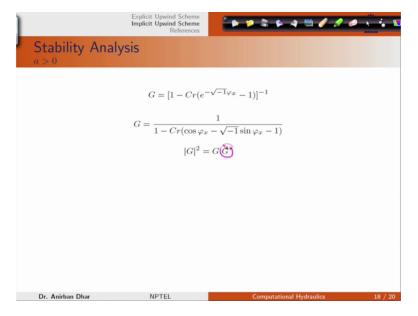
If we further simplify, we can write G in this format. Now we need to define the modulus for this G to evaluate the stability of the scheme.

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	Explicit Upwind Scheme Implicit Upwind Scheme References	*****
Stability Analy	/sis	
Thus the error equ	ation can be represer	ited as,
ϵ^n_i	$^{+1} = \begin{cases} \epsilon_i^n + a \frac{\Delta t}{\Delta x} (\epsilon_i^n) \\ \epsilon_i^n + a \frac{\Delta t}{\Delta x} (\epsilon_i^n) \end{cases}$	
$G = \frac{1}{2}$	$\frac{a^{n+1}_{i}}{\epsilon_{i}^{n}} = \begin{cases} 1 + a \frac{\Delta t}{\Delta x} \left(1 + a \frac{\Delta t}{\Delta x} \left(1 + a \frac{\Delta t}{\Delta x} \right) \right) \end{cases}$	$ \begin{array}{l} \displaystyle \frac{\epsilon_{i-1}^{n+1}}{\epsilon_i^n} - \frac{\epsilon_i^{n+1}}{\epsilon_i^n} \\ \displaystyle \epsilon_i^{n+1} - \frac{\epsilon_{i+1}^{n+1}}{\epsilon_i^n} \end{array} \right) if a > 0 \\ \end{array} \\ \left. \begin{array}{l} \displaystyle if a > 0 \\ \displaystyle if a < 0 \end{array} \right. $
$Cr = a \frac{\Delta t}{\Delta x}$	$G = \begin{cases} 1 + CrG(e^{-\sqrt{2}}) \\ 1 + CrG(e^{\sqrt{-1}}) \end{cases}$	$\overline{\varphi_x} - 1$ if $a > 0$ $\overline{\varphi_x} - 1$ if $a < 0$
		$ e^{\varphi_x} - 1)]^{-1} if a > 0 $ $ e^{-1} - 1)]^{-1} if a < 0 $
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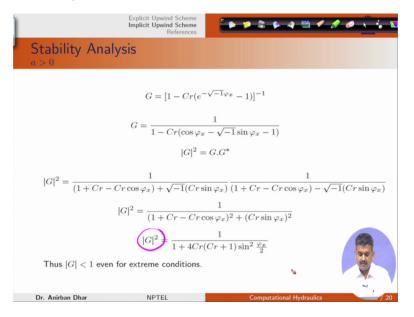
In this case considering the thing for A greater than zero, we can write like this. And with this we need to evaluate the G square. This is the conjugate complex number in this case.

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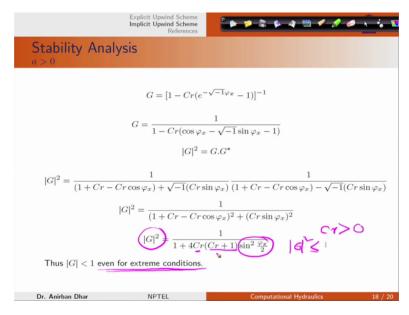
So with this if you proceed and write this thing for (mod) modulus of G square, we can get this one and further simplification will give this G square condition.

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In this case we can see that Cr is always positive. In sin square phi x by 2 is always positive, so Cr plus 1 is positive, CR is positive. So 1 by 1plus is something positive. This will give always this G Square or modulus of G square less than equal to 1. For extreme condition, this will be less than equal to 1.

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So we can see that for implicit scheme it is unconditionally stable. Now we have covered the finite difference thing. In the next lecture onwards we will be starting finite volume. Thank you.