

**Chemical Principles II**  
**Professor Dr. Arnab Mukherjee**  
**Department of Chemistry**  
**Indian Institute of Science Education and Research, Pune**  
**Tutorial Problem - 13**

So we have learned how to change one (thermo dynic vareb) one thermodynamic derivative to another thermodynamics derivative using Jacobian technique and I have shown some usefulness in terms of finding out CP minus CV, a general expression of CP minus CV and also an arbitrary partial derivatives of del G by Del H at a constant V, I guess. So the point is that we can do whatever partial derivative it is using the Jacobian technique and that will make our life much simpler and will help us also to find out useful quantities.

For example you know joule coefficient or joule Thomson coefficient and things like that. So we are going to go and do some more thermodynamic derivatives you know it is called reduction of thermodynamic derivative, so we are going to go and do some more reduction of thermodynamic derivatives and see that how we can easily you know use them and you know convert them and we get more accustomed to that.

So we will do some more problems to help you understand that things better.

(Refer Slide Time: 01:28)

Problem 1

---

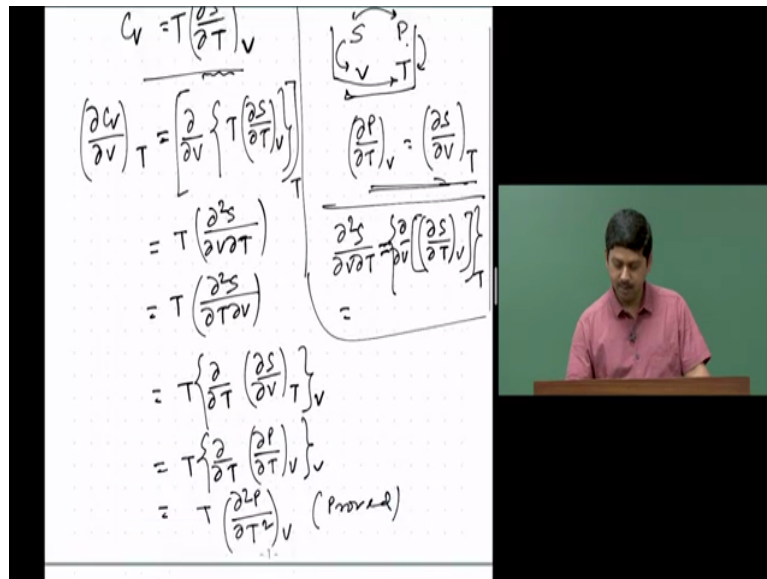
Derive the following equation

$$(i) \left( \frac{\partial C_V}{\partial V} \right)_T = T \left( \frac{\partial^2 P}{\partial T^2} \right)_V \quad \text{and prove that for ideal gas it is constant}$$



So today is problem that we are going to do is that problem number 1 has these derivations that derive this following Del C V by Del V so I am going to do it here.

(Refer Slide Time: 01:37)



So the first one is prove  $\Delta C_V$  by  $\Delta V$  at a constant  $T$  is equal to  $T \Delta^2 P / \Delta T^2$  by  $V$ , ok. So we have to do in this (pra) problem  $\Delta C_V$  by  $\Delta V$  at a constant  $T$  is equal to  $\Delta^2 P$  by  $\Delta T^2$  at a constant  $V$ . Now we will start with the definition of  $C_V$ ,  $C_V$  is  $\Delta S$  by  $\Delta T$ ,  $T \Delta S$  by  $\Delta T$  at a constant  $V$ , so we know that. So once we start with that we will have to see whether we can get the right hand side expression or not but before we do that let us also understand that if we get these things from somewhat from the Maxwell relation or not.

So we write Maxwell relation  $S$   $P$   $T$   $V$  and we see that we can get either  $\Delta S$  by  $\Delta P$  or we can get  $\Delta S$  by  $\Delta V$  or  $\Delta P$  by  $\Delta T$ , so  $\Delta P$  by  $\Delta T$  we can get though ok, that is a important thing, so  $\Delta P$  by  $\Delta T$  you can get and we know that from Maxwell relation that  $\Delta P$  by  $\Delta T$  at a constant  $V$  is we have to do this way is  $\Delta S$  by  $\Delta V$  at a constant  $T$ , ok.

So that is fine we have that however we see in this particular formula it is just interchange it is  $\Delta S$  by  $\Delta T$  at constant  $V$  but we need  $\Delta S$  by  $\Delta V$  at a constant  $T$  so they are not exactly same, so we will try to see whether we can get that or not. So now before you get that so  $C_V$  definition is this one, so therefore if I take one more derivative  $\Delta C_V$  by  $\Delta V$   $T$  if I do that then essentially I have to take the derivative of this entire quantity at a constant temperature.

Now if I want to do that then you see that  $T$  becomes a constant, so  $T$  will come out and it will give me  $\Delta^2 S$  by  $\Delta V \Delta T$  which by notation means this is a shorthand notation  $\Delta^2 S$  by  $\Delta V \Delta T$ , the shorthand notation is  $\Delta S$  by  $\Delta T$  at a constant  $V$  and then the whole thing taking a derivative with respect to  $V$  at a constant  $T$ , so that is equivalent to writing as  $\Delta S \Delta^2 S$  by  $\Delta V \Delta T$ .

Now we know that S is exactly differential, so therefore we can write that as Del 2 S by Del T del V as well, now once we do that we can write T Del Del T Del S by Del V at a constant T and this whole thing at a constant V. now we know that Del S by Del V T from Maxwell relation is nothing but Del P by Del T V and whole thing is at a V constant V which then proves our point that it is nothing but Del 2 P by Del T 2 at a constant V and that is what we wanted to show.

You see the Jacobian cannot be used in this particular case because Jacobian will take the derivative to the heat capacities like C P, C V and then alpha and Kappa T but here we wanted a conversion of relation and the way I understood that we have to do the this formula for exact differential is when I look into Del P Del T as this one then I realize that which means that the V and T are interchange and then I have done that, so you have to do that as well.

(Refer Slide Time: 06:08)

$$= T \left( \frac{\partial^2 S}{\partial T^2} \right)_V \quad (\text{proved})$$

$$(ii) \left( \frac{\partial C_p}{\partial p} \right)_T = -T \left( \frac{\partial^2 P}{\partial T^2} \right)_p$$

$$C_p = T \left( \frac{\partial S}{\partial T} \right)_p$$

$$\left( \frac{\partial C_p}{\partial p} \right)_T = T \left( \frac{\partial^2 S}{\partial p \partial T} \right)_p$$

$$= T \frac{\partial^2 S}{\partial T \partial p}$$

$$= T \left\{ \frac{\partial}{\partial T} \left( \frac{\partial S}{\partial p} \right)_T \right\}_p$$

$$= T \left\{ \frac{\partial}{\partial T} \left( \frac{-\partial V}{\partial T} \right)_p \right\}_p$$

$$\left( \frac{\partial S}{\partial p} \right)_T = \left( \frac{\partial V}{\partial T} \right)_p$$

S P  
 V T

$$C_p = T \left( \frac{\partial S}{\partial T} \right)_P$$

$$\left( \frac{\partial C_p}{\partial P} \right)_T = T \left( \frac{\partial^2 S}{\partial P \partial T} \right)$$

$$= T \frac{\partial^2 S}{\partial T \partial P}$$

$$= T \left\{ \frac{\partial}{\partial T} \left( \frac{\partial S}{\partial P} \right)_T \right\}_P$$

$$= T \left\{ \frac{\partial}{\partial T} \left( -\frac{\partial V}{\partial T} \right)_P \right\}_P$$

$$= -T \left( \frac{\partial^2 V}{\partial T^2} \right)_P$$

$$\left( \frac{\partial S}{\partial P} \right)_T = - \left( \frac{\partial V}{\partial T} \right)_P$$

So now go to the second problem, we have to show that  $\Delta C_p$  by  $\Delta P$  at a constant  $T$  is minus  $T \Delta^2 P$  by  $\Delta T^2$  at a constant  $P$ . so again I think it is very similar to the problem above just let us write  $C_p$  as  $T \Delta S$  by  $\Delta T$  at constant  $P$  that is the  $C_p$  and we need to do the derivative of this particular one, so we need to do  $\Delta C_p$  by  $\Delta P$  at a constant  $T$  which means we have to take the derivative of this quantity at a constant  $T$ , so  $T$  will come out and we will get  $\Delta^2 S$  by  $\Delta P \Delta T$ , ok.


So again it is a shorter notation  $\Delta^2 S$  by  $\Delta P \Delta T$  meaning first we will take the derivative with respect to  $T$  at a constant  $P$  then the whole thing at a constant with respect to  $P$  at a constant  $T$ . So this is same as since it is exact differential as  $\Delta S \Delta P$  which again means that first derivative with respect to  $P$  at a constant  $T$  then with respect to  $T$  at a constant  $P$ . So again we just going to look at the Maxwell relation  $S P T V$  and we know that  $\Delta S$  by  $\Delta P$  is  $\Delta P$  by  $\Delta T$  so  $\Delta S$  by  $\Delta P$  is  $\Delta V$  by  $\Delta T$  we know that, so but there is a plus and minus thing as you can see from my arrows, so  $\Delta S$  by  $\Delta P$  at a constant  $T$  is  $\Delta V$  by  $\Delta T$  at a constant  $P$  but with a minus.

So it will be  $T \Delta \Delta T \Delta S$  by  $\Delta P$  at a constant  $T$  and the whole thing it is at the constant  $P$  which is  $T \Delta \Delta T$  and  $\Delta S$  by  $\Delta P$  at a constant  $T$  is minus  $\Delta V$  by  $\Delta T$  at a constant  $P$  which is and the whole thing is at a constant  $P$  which is minus  $T \Delta^2 V$  by  $\Delta T^2$  at a constant  $P$ . Now but I think it is written question is given wrongly it should be  $V$  not  $P$ , ok.

(Refer Slide Time: 08:56)

Problem 2


A measure of the results of an adiabatic Joule free expansion is provided by the Joule coefficient  $\mu_J = (\partial T / \partial V)_U$   
Show that

$$\mu_J = -\frac{1}{C_V} \left( \frac{\alpha T}{\kappa_T} - P \right)$$
$$\mu_J = -\frac{1}{C_V} \left( \frac{\alpha T}{\kappa_T} - P \right)$$


So we will let us go to the next problem, so next problem says that a measure of the results of adiabatic joule free expansion is provided by the joule coefficient  $\mu_J$ , show that  $\mu_J$  is equal to  $1/C_V$ , so this is nothing but  $\alpha$  because  $\alpha$  is the coefficient of thermal expansion and this is a cup and this is nothing but  $\kappa_T$  minus  $P$ , so I will just write it clearly, so  $\mu_J$  show that  $\mu_J$  is  $1/C_V$   $\alpha T$  by  $\kappa_T$  minus  $P$ .

So let us see if that is the thing we can show or not.

(Refer Slide Time: 09:37)

$$\begin{aligned} \mu_J &= \left( \frac{\partial T}{\partial V} \right)_U \\ &= \frac{\partial(\tau, U)}{\partial(V, U)} \\ &= \frac{\partial(\tau, U)}{\partial(\tau, V)} \frac{\partial(\tau, V)}{\partial(V, U)} \\ &= \left( \frac{\partial U}{\partial V} \right)_\tau \left[ -\frac{\partial(V, T)}{\partial(V, U)} \right] \end{aligned}$$


$$\begin{aligned}
 &= -\frac{\partial(T, V)}{\partial(U, V)} \left[ -\frac{\partial(U, T)}{\partial(U, V)} \right] \\
 &= -\left(\frac{\partial U}{\partial V}\right)_T \left(\frac{\partial T}{\partial U}\right)_V \\
 &= -\left(\frac{\partial U}{\partial V}\right)_T \times \left[\left(\frac{\partial U}{\partial T}\right)_V\right]^{-1} \\
 &= -\frac{1}{C_V} \left(\frac{\partial U}{\partial V}\right)_T
 \end{aligned}$$

-3-

So by definition you see  $\mu_J$  is  $\Delta T$  change in temperature with respect to change in volume at a constant  $U$ , so you know in this particular case we have to use the Jacobian because we have to bring the  $U$  in the numerator, so it is  $\Delta T \Delta U$  by  $\Delta V \Delta U$ , I can write that and then I can take as  $\Delta T \Delta U$  by  $\Delta T \Delta P$  and I have to multiply by  $\Delta T \Delta P$  and  $\Delta V \Delta U$ .

Now this one the left hand side gives us  $\Delta U$  by  $\Delta P$  at constant  $T$ , so instead of  $T \Delta P$  let us change it to  $V$  because  $V$  is good for always  $U$ , so use a natural variable is  $V$  and  $T$ , so let us change it to  $V$  so instead of  $P$  I will just put  $V$ , so yeah and that also is good because here also I am getting  $V$  common. So now we can write that as  $\Delta U$  by  $\Delta V$  at a constant  $T$  multiplied by I take a minus and then I  $\Delta V \Delta T$  by  $\Delta V \Delta U$  and you know why the minus is because I switched the numerator positions  $T$  and  $V$ .

Now  $\Delta U$  by  $\Delta V \Delta T$  do you remember what it is? So  $U$  is a natural variable is  $PdV$  ok, so  $U$  is  $TdS$  minus  $PdV$  so that does not help us because a  $\Delta U$  by  $\Delta V$  at a constant  $S$  only we would have got  $T$  but  $\Delta U$  by  $\Delta V$  at constant  $T$  is given so that does not help us, so we will deal with that later on. So now just let us continue  $\Delta U$  by  $\Delta V \Delta T$  or I will put a minus here and that gives us  $\Delta T$  by  $\Delta U$  at a constant  $V$ , so which is  $\Delta U$  by  $\Delta V$  at a constant  $T$  multiplied by  $\Delta U$  by  $\Delta T$  at a constant  $V$  inverse which as you know is  $C_V$ , so this is minus 1 by  $C_V \Delta U$  by  $\Delta V$  at a constant  $T$  that is what we got, that is what the Joule coefficient is.

(Refer Slide Time: 12:31)

$dU = Tds - PdV$   
 divide both sides by  $dV$  at -  
 const  $T$   

$$\left(\frac{\partial U}{\partial V}\right)_T = T \left(\frac{\partial S}{\partial V}\right)_T - P \left(\frac{\partial V}{\partial V}\right)_T$$

$$= T \left(\frac{\partial S}{\partial V}\right)_T - P$$

$$= T \left(\frac{\partial P}{\partial T}\right)_V - P$$

$\begin{matrix} P \\ \swarrow \\ V \end{matrix} \begin{matrix} T \\ \rightarrow \end{matrix}$

$$\left(\frac{\partial T}{\partial V}\right)_U = \frac{-1}{C_V} \left[ T \left(\frac{\partial P}{\partial T}\right)_V - P \right]$$

$$\left(\frac{\partial P}{\partial T}\right)_V = \frac{\alpha T}{K_T}$$
 To prove,  $\left(\frac{\partial P}{\partial T}\right)_V = \frac{\alpha}{K_T}$

$\begin{matrix} P \\ \swarrow \\ V \end{matrix} \begin{matrix} T \\ \rightarrow \end{matrix}$

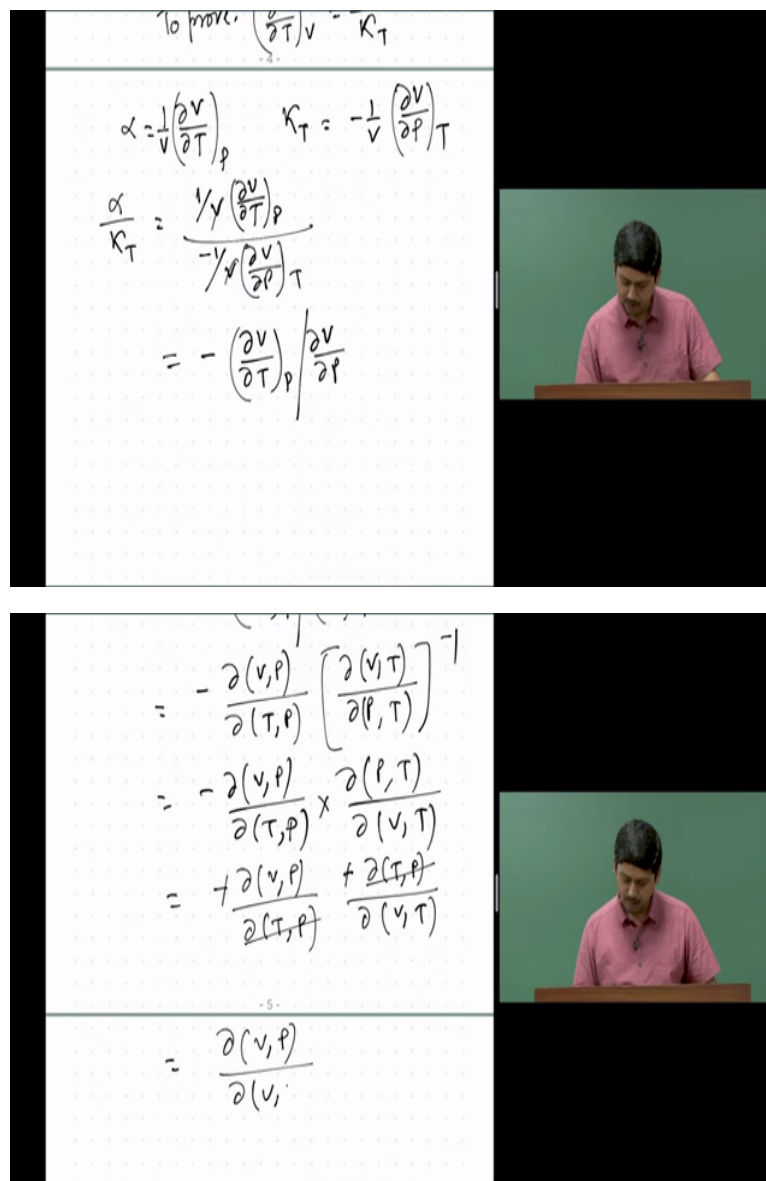
Now we have to do Del U by Del V at a constant T, so we will start with this expression dU equal to TdS minus PdV. Now we cannot do directly from here Del U by Del V at constant T because T is not a variable here, so T is just a parameter so therefore we have to divide both sides by dV and now at constant T, so what you are going to get from there is Del U by Del V at a constant T is T Del S by Del V at a constant T minus P Del V by Del V at a (cons) Del V by Del V at constant T which is T Del S by Del V at a constant T minus P, say I am getting back minus P.

So now what else see we also have got now I have to just do Del S by Del V and that I will now use formula for or Maxwell relations so S P T V, so you see Del S by Del V is at a constant T is nothing but Del P by Del T at constant V minus P. So this also we got but that

does not help us in getting the Alpha and Kappa T because in alpha and Kappa T we know, so this much we got.

So now if I just write my Del T by Del V at a constant U I have already got minus 1 by C V T Del P del T V minus P, now I have to prove that this quantity T Del P del T V is equal to alpha T by Kappa, so in fact T is also common, so I have to prove Del P Del T sorry it was T Del P Del T V is Kappa T I will just rewrite is it again I have to prove Del P Del T V is alpha by Kappa T, this what I have to proof then I can prove the whole relation because everything else is fine with this.

(Refer Slide Time: 15:08)



10 prove:  $\left(\frac{\partial T}{\partial V}\right)_U = K_T$

$$\alpha = -\frac{1}{V} \left(\frac{\partial V}{\partial T}\right)_P \quad K_T = -\frac{1}{V} \left(\frac{\partial V}{\partial P}\right)_T$$

$$\frac{\alpha}{K_T} = \frac{\frac{1}{V} \left(\frac{\partial V}{\partial T}\right)_P}{-\frac{1}{V} \left(\frac{\partial V}{\partial P}\right)_T}$$

$$= - \left(\frac{\partial V}{\partial T}\right)_P \bigg/ \left(\frac{\partial V}{\partial P}\right)_T$$
  

$$= - \frac{\partial(V, P)}{\partial(T, P)} \left[ \frac{\partial(V, T)}{\partial(P, T)} \right]^{-1}$$

$$= - \frac{\partial(V, P)}{\partial(T, P)} \times \frac{\partial(P, T)}{\partial(V, T)}$$

$$= + \frac{\partial(V, P)}{\partial(T, P)} \frac{\partial(T, P)}{\partial(V, T)}$$
  

$$= \frac{\partial(V, P)}{\partial(V, T)}$$



$$\begin{aligned}
&= -\frac{\partial(V,T)}{\partial(T,P)} \times \frac{\partial(V,T)}{\partial(V,T)} \\
&= +\frac{\partial(V,P)}{\partial(T,P)} + \frac{\partial(T,P)}{\partial(V,T)} \\
&= \frac{\partial(V,P)}{\partial(V,T)} \\
&= \frac{\partial(P,V)}{\partial(T,V)} \\
&= \left(\frac{\partial P}{\partial T}\right)_V
\end{aligned}$$

Now we know what is alpha, alpha is coefficient of thermal expansion, so  $\Delta V / \Delta T$  at a constant pressure  $1/V$  and we know what is Kappa T right minus  $1/V \Delta V / \Delta P$  at a constant T, so alpha by Kappa T is  $1/V \Delta V / \Delta T$  at constant P by minus  $1/V \Delta V / \Delta P$  at a constant T,  $1/V$   $1/V$  cancels each other there is a minus sign we will have to deal with that, now  $\Delta V / \Delta T$  at a constant pressure  $\Delta V / \Delta T$  no I will just take by sign  $\Delta V / \Delta P$  at the constant T.

Now I am going to use a Jacobian technique minus  $\Delta V / \Delta P$  by  $\Delta T / \Delta P$  and divided by  $\Delta V / \Delta T$  is not a good way of writing it, so I will put an inverse  $\Delta V / \Delta T$  by  $\Delta P / \Delta T$  inverse meaning just I will just change it  $\Delta V / \Delta P$   $\Delta T / \Delta P$  multiplied by  $\Delta P / \Delta T$  by  $\Delta V / \Delta T$ . Now there is a  $P / T$  so I am not going to change that  $V / P$   $\Delta T / \Delta P$  I will going to change take a minus and say  $T / P$   $\Delta V / \Delta T$ , now  $\Delta T / \Delta P$   $\Delta T / \Delta P$  cancels giving us and minus minus cancels giving us  $\Delta V / \Delta P$  by  $\Delta V / \Delta T$  which is nothing but  $\Delta P / \Delta V$  by  $\Delta T / \Delta V$  which is nothing but  $\Delta P / \Delta T$  at constant V.

(Refer Slide Time: 17:21)

$$\left(\frac{\partial T}{\partial V}\right)_U = -\frac{1}{C_V} \left[ T \left(\frac{\partial P}{\partial T}\right)_V - P \right]$$

$$= -\frac{1}{C_V} \left[ T \frac{\alpha}{K_T} - P \right] \text{ (proved)}$$

To prove,  $\left(\frac{\partial P}{\partial T}\right)_V = \frac{\alpha}{K_T}$

$\alpha = \frac{1}{V} \left(\frac{\partial V}{\partial T}\right)_P$       $K_T = -\frac{1}{V} \left(\frac{\partial V}{\partial P}\right)_T$

$\alpha = \frac{1}{V} \left(\frac{\partial V}{\partial T}\right)_P$

So now I am coming back and I see that I can write this whole expression I am going to just remove this part, now I am going to whole expression I am going to write as T alpha by Kappa T minus P and proved, ok.

(Refer Slide Time: 17:47)

**Problem 2**

---

A measure of the results of an adiabatic Joule free expansion is provided by the Joule coefficient  $\mu_J = (\partial T/\partial V)_U$   
 Show that

$$\mu_J = -\frac{1}{C_V} \left( \frac{\beta T}{K_T} - P \right) \quad \mu_J = -\frac{1}{C_V} \left( \frac{\alpha T}{K_T} - P \right)$$

and for Joule-Thomson coefficient

$$\mu_{JT} = \frac{V}{C_p} (\beta T - 1)$$

$-(\partial T)_V$


$$\mu_{JT} = -\frac{V}{C_p} (\alpha_T - 1)$$

$$\mu_{JT} = \left(\frac{\partial T}{\partial P}\right)_H \quad \left. \begin{array}{l} \\ \\ \\ \end{array} \right\} C_p = T \left(\frac{\partial S}{\partial T}\right)_P$$

$$= \frac{\partial(\tau, H)}{\partial(P, H)}$$

$$=$$

-6-



$$\mu_{JT} = \left(\frac{\partial T}{\partial P}\right)_H \quad \left. \begin{array}{l} \\ \\ \\ \end{array} \right\} C_p = T \left(\frac{\partial S}{\partial T}\right)_P$$

$$= \frac{\partial(\tau, H)}{\partial(P, H)} \quad \left. \begin{array}{l} \\ \\ \\ \end{array} \right\} \begin{array}{l} dH = dU + PdV \\ + v dP \\ = T ds + v dP \end{array}$$


$$\left(\frac{\partial H}{\partial T}\right)_P = T \left(\frac{\partial S}{\partial T}\right)_P + v \left(\frac{\partial P}{\partial T}\right)_P$$

$$\mu_{JT} = \frac{\partial(\tau, H)}{\partial(P, H)} \quad \left(\frac{\partial H}{\partial P}\right)_T = T \left(\frac{\partial S}{\partial P}\right)_T + v$$

$$= \frac{\partial(H, T)}{\partial(H, P)} \quad \begin{array}{l} S \ P \\ \sim T \end{array}$$

$$= \frac{\partial(H, T)}{\partial(T, P)} \frac{\partial(T, P)}{\partial(H, P)}$$

-6-



$$= \frac{\partial(H, T)}{\partial(T, P)} \frac{\partial(H, P)}{\partial(H, P)}$$


$$= -\frac{\partial(H, T)}{\partial(P, T)} \times \left[\frac{\partial(H, P)}{\partial(T, P)}\right]^{-1}$$

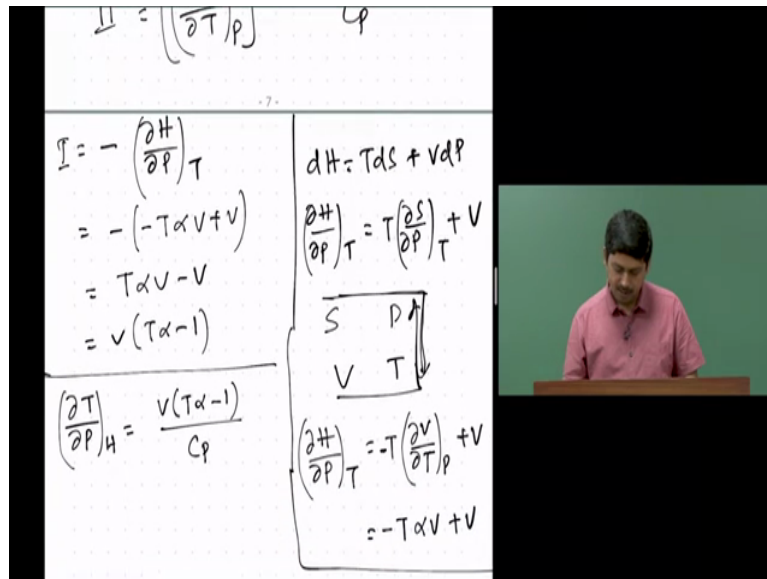
$$= -\left(\frac{\partial H}{\partial P}\right)_T \times \left[\left(\frac{\partial H}{\partial T}\right)_P\right]^{-1} \quad \text{--- (1)}$$

$$= \text{I} \times \text{II}$$

$$\text{II} = \left[\left(\frac{\partial H}{\partial T}\right)_P\right]^{-1} = \frac{1}{C_p}$$

-7-

$$\text{I} = -\left(\frac{\partial H}{\partial P}\right)_T \quad dH = T ds + v dP$$




Now for Joule Thomson coefficient we have to show that  $\mu_{JT}$  is equal to  $-V/C_p$ , so this  $\beta$  is  $\alpha$  actually it suggests a notational issue  $V/C_p \alpha T - 1$ . Now Joule Thomson coefficient is as you know that it is change in temperature with respect to pressure at a constant enthalpy, so this one we have to show as the one that is mentioned here since  $H$  is held as constant we have to bring the  $H$  in the numerator and we can do that by taking this and immediately we need something.

So for example we need the  $C_p$  and you know just we can just do a back calculation and just find out that what  $C_p$  is  $C_p$  is  $\Delta S / \Delta T$  at constant  $P$ , so we need some  $S$  somewhere and  $\alpha$  is also there, so yeah but we need the  $S$  and  $T$  and pressure so that you have to just remember and also we have to remember our  $H$ ,  $H$  is nothing but  $H$  is  $U + PV$  so it is  $dU + PdV + VdP$  and  $dU + PdV$  as you know is nothing but  $TdS$ , so  $TdS + VdP$  is  $\Delta H$ .

So we just keep that those things in mind and continue our derivation, so  $\mu_{JT}$  is  $\Delta T / \Delta P$  we can take the  $H$  on the left hand side, so  $\Delta H / \Delta T$  by  $\Delta H / \Delta P$ . So now what do you want as a derivative so for example if I do  $\Delta H / \Delta T$  we can do  $\Delta H / \Delta T$  at constant  $P$  yeah that we can do here so for example we do from this above expression  $\Delta H / \Delta T$  at a constant  $P$  will give us  $T \Delta S / \Delta T$  at constant  $P$  plus, so I am doing with respect to  $T$  right so  $V \Delta P / \Delta T$  at constant  $P$ , so that goes 0, so  $\Delta H / \Delta T$  at constant  $P$  is nothing but that, so we need that.

So we can do this  $\Delta H / \Delta T$  and so we need with respect to temperature at a held at constant pressure, so that means we need this so  $\Delta H / \Delta P$  by  $\Delta T / \Delta P$  we need that, so  $H / P$  is here for

sure so  $T P$  is not here, so let us let us write that  $\Delta T P$  multiplied with  $\Delta T P$  and  $\Delta H P$ , so you see I wanted to bring  $\Delta S$  by  $\Delta T$  at constant  $P$  and I just that is that is why I multiplied with  $T P$  in the numerator and denominator.

So now what I am going to see is what is the left hand side? So I can just take a minus sign and rewrite it  $\Delta H T$  by  $\Delta P T$  multiplied by  $\Delta H P$  by  $\Delta T P$  inverse, left hand side will give us  $\Delta H$  by  $\Delta P$  at a constant  $T$  and this is going to give us  $\Delta H$  by  $\Delta T$  at a constant  $P$  and inverse of that. Now this one what is  $\Delta H$  by  $\Delta P$  at constant  $T$ ? So let us figure out from that, so  $\Delta H$  by  $\Delta P$  at a constant  $T$  I have to take the differentiation again so if I do that this equation from this equation  $\Delta H$  by  $\Delta P$  let me see a  $\Delta H$  by  $\Delta P$ ,  $\Delta P$  at a constant temperature is  $T \Delta S$  by  $\Delta P$  at a constant temperature plus  $V$  because  $\Delta P \Delta P$  will cancel.

So now when we do  $SP T V$  ok, so this so this will be a little bit longer not just simple once we will do that more carefully in the denominator ok so first just write this one. So we are going to do individually all that ok, so let us do both of them separately, so let us write that as equation 1 and just write that as term 1 into term 2, now term 2 is by definition is very simple that is  $\Delta H$  by  $\Delta T P$  inverse we have shown that that is nothing but  $\Delta H$  by  $\Delta T P$  is nothing but  $T \Delta S$  by  $\Delta T P$  which is nothing but  $C P$ , so this is nothing but  $C P$  so we can write as  $1$  by  $C P$  so you got as you can see we got our  $1$  by  $C P$  right there from the term 2.

Now term 1 term one is minus  $\Delta H$  by  $\Delta P T$ , so for that what we need is we need to write again  $dH$  as  $TdS$  plus  $VdP$  take derivative with respect to  $P$  on both sides, so  $\Delta H$  by  $\Delta P$  at a constant  $T$  is  $T \Delta S$  by  $\Delta P$  at constant  $T$  plus  $V$  because  $\Delta P$  by  $\Delta P$  will be  $1$ . Now once we have that we have to use Maxwell relation  $SP T V$  and we know  $\Delta S$  by  $\Delta P$  at a constant  $T$  is  $\Delta V$  by  $\Delta T$  at a constant  $P$ .

So  $\Delta H$  by  $\Delta P$  at a constant  $T$  is  $T \Delta V$  by  $\Delta T$  at a constant  $P$  plus  $V$ , now what is  $\Delta V$  by  $\Delta T$  at constant  $P$ ? So yeah so this and ok and it will be a minus sign because as you can see that this is going down and this going up so there will be a minus here, so that will be minus  $T$  now what is  $\Delta V$  by  $\Delta T$ ? So we can say coefficient of thermal expansion multiplied by  $V$ , so this will be  $\alpha V$  Plus  $V$  so we got the whole thing so that means this term 2 our term 1 which is this is equal to minus of this quantity minus  $T \alpha V$  Plus  $V$  which is going to give us  $T \alpha V$  minus  $V$  or  $V T \alpha$  minus  $1$ , ok.

So now we put together whatever our answer was our answer that we wanted is  $\Delta T$  by  $\Delta P$  at a constant  $H$  is equal to  $V T \alpha \text{ minus } 1$  divided by  $C P$  did you get that? So the way it is written is  $\text{minus } V \text{ by } C P$ , so  $\text{minus } V \text{ by } C P$  1 minus yeah so I think that we have made a mistake with minus somewhere let us see where did I let us see we start with this  $\Delta T$  by  $\Delta P$  at a constant  $H$  this is not a problem this part is fine here I have interchanged  $P$  and  $T$ , so this this part is fine so 1 is having the minus, now 2 the side 2 that is  $T P H P H P T P$ , so 1 by  $C P$  is fine.

Now  $\text{minus } \Delta H \text{ by } \Delta P T$ , now  $\Delta H \text{ by } \Delta P T$  is written here, so I guess we got the thing right although there is there difference with just a minus, so the question was little bit wrong so this is correct the question was wrong, so this should be plus and not minus, so I got the answer so we will stop here and we will continue.