

**Chemical Principles II**  
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**Module 03**  
**Lecture 18**  
**Adiabatic Reversible Work**

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The image shows a blackboard with handwritten mathematical equations. The first equation is  $dU = dQ + dW$ , with a zero written above the  $dQ$  term and an arrow pointing to it from above. The second equation is  $dU = dW$ . The third equation is  $dU = C_v dT$ . The fourth equation is  $\Delta U = \int_{T_1}^{T_2} C_v dT = C_v (T_2 - T_1) = W$ . In the bottom right corner of the blackboard area, there is a small inset video frame showing a man in a red shirt.

So now we are going to talk about adiabatic work, so what is an adiabatic work? Adiabatic work is where you have no exchange of heat and I am going to just derive that adiabatic work here, so  $dU$  equal to  $dQ + dW$ , in case of adiabatic  $dQ$  is equal to 0 so we are going to have  $dU$  equal to  $dW$  for adiabatic cases. Now we know that for an ideal gas we can write  $dU$  as  $C_v dT$  now given temperature difference so in adiabatic case of course if there is a work done that means some value of  $W$  is there, then you see that  $C_v$  because  $C_v$  is taken to be constant, temperature is going to change and that is the reason that in adiabatic case the temperature changes.

How much is the work done that will depend on how much change in the temperature is happening. Let us say the temperature is changing from  $T_1$  to  $T_2$  so in that case total change in  $U$  will be integration of  $C_v dT$  from  $T_1$  to  $T_2$  giving us  $C_v T_2$  minus  $T_1$ . And since this is the amount of work done so the work done in adiabatic case is nothing but  $C_v T_2$  minus  $T_1$  that is when you write the work in terms of temperature. However if you want to write that work in terms of pressure and volume you are going to do a little bit more work okay and that is what you are going to do now so let us go back again.

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$P = \frac{nRT}{V}$   
 $dU = dW$   
 $C_v dT = -P dV$   
 $C_v dT = -\frac{nRT}{V} dV$   
 $C_v \frac{dT}{T} = -\frac{nR}{V} dV$   
 $C_v \int_{T_1}^{T_2} \frac{dT}{T} = -\frac{nR}{V} \int_{V_1}^{V_2} dV$   
 $C_v \ln\left(\frac{T_2}{T_1}\right) = -R \ln\left(\frac{V_2}{V_1}\right)$   
 $\ln\left(\frac{T_2}{T_1}\right) = -\frac{R}{C_v} \ln\left(\frac{V_2}{V_1}\right)$   
 $\ln\left(\frac{T_2}{T_1}\right) = -(\gamma - 1) \ln\left(\frac{V_2}{V_1}\right)$

$C_p - C_v = R$   
 $\frac{C_p}{C_v} - 1 = \frac{R}{C_v}$   
 $\gamma - 1 = \frac{R}{C_v}$

So now we know that  $dU$  equal to  $dW$  for ideal gas or  $C_v dT$  equal to minus  $PdV$ . Now for ideal gas we know that  $PV$  is equal to  $nRT$  and therefore I am going to replace  $P$  by  $nRT/V$  by  $V$ , so  $C_v dT$  is equal to minus  $nRT/V dV$  but left-hand side I have a  $dT$  you know which takes care of the temperature change so I have to bring this  $T$  below here because there is a function of temperature  $C_v dT/T$  is equal to minus  $nR/V dV$  okay. Now I can write  $C_v dT/T$  is minus  $R/V$  bar,  $V$  bar is  $V$  by  $n$  molar volume  $dV$ .

Now I integrate both sides, so when I am integrating from  $T_1$  to  $T_2$  what I am going to get is  $C_v \ln T$  so it was current actually  $T_2$  by  $T_1$  it was correct, right-hand side is minus  $R \ln V$  by  $V_1$  we will put the bar. So now you see now I equate them both so in that case I am bringing  $C_v$  down below. So  $\ln T_2$  by  $T_1$  is equal to minus  $R$  by  $C_v \ln V_2$  by  $V_1$ , and what is  $R$  by  $C_v$ ? So are by  $C_v$  is, you know  $C_p$  minus  $C_v$  is equal to  $R$  right divide both sides by  $C_v$ , you will get  $C_p$  by  $C_v - 1$  equal to  $R$  by  $C_v$  right, so  $R$  by  $C_v$  is  $C_p$  by  $C_v - 1$ , and  $C_p$  by  $C_v$  is typically called  $\gamma$  so  $\gamma - 1$  is  $R$  by  $C_v$ . So we can now write that  $\ln T_2$  by  $T_1$  is equal to minus  $\gamma - 1 \ln V_2$  by  $V_1$  bar ok.

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$$\ln\left(\frac{T_2}{T_1}\right) = -\frac{R}{C_v} \ln\left(\frac{V_2}{V_1}\right)$$

$$\ln\left(\frac{T_2}{T_1}\right) = \frac{R}{C_v} \ln\left(\frac{V_1}{V_2}\right)$$

$$\ln\left(\frac{T_2}{T_1}\right) = \ln\left(\frac{V_1}{V_2}\right)^{R/C_v}$$

$$\frac{T_2}{T_1} = \left(\frac{V_1}{V_2}\right)^{\gamma-1} \quad \text{--- (1)}$$

$$\frac{P_2 V_2}{P_1 V_1} \times \left(\frac{V_1}{V_2}\right)^{\gamma-1} \Rightarrow P_2 V_2 V_2^{\gamma-1} = P_1 V_1 V_1^{\gamma-1}$$

$$P_2 V_2^\gamma = P_1 V_1^\gamma \quad [PV^\gamma = \text{const}]$$

$$PV^\gamma = \text{const.}$$

$PV = nRT$   
 $P\bar{V}_1 = RT_1$   
 $P_2\bar{V}_2 = RT_2$   
 $\frac{T_2}{T_1} = \frac{P_2\bar{V}_2}{P_1\bar{V}_1} \quad \text{--- (2)}$

Refer all v as  $\bar{V}$

If I will have a little bit more space here I will write it so how much we got? We got that  $\ln T_2$  by  $T_1$  was minus  $R$  by  $C_v \ln V_2$  by  $V_1$ , this was we had and we show that this is nothing but are by  $C_v \ln V_1$  by  $V_2$  okay, which means  $\ln T_2$  by  $T_1$  is equal to  $\ln V_1$  by  $V_2$  to the power  $R$  by  $C_v$  and we know that  $R$  by  $C_v$  okay we will do that  $R$  by  $C_v$  later on. So now when we have this equation both sides  $\ln$  is there then we can write  $T_2$  by  $T_1$  is equal to  $V_1$  by  $V_2$  to the power  $R$  by  $C_v$  which we have calculated to be  $\gamma - 1$ , so this is the equation we have but we cannot forget the 1<sup>st</sup> law of thermodynamics sorry we cannot forget the equation of ideal gas.

What was that?  $PV$  is equal to  $nRT$  that is always there or  $PV$  power equal to  $RT$  which means if I have  $T_1$  then I have  $P_1 V_1$  bar and if I have  $T_2$  then I have  $P_2 V_2$  right. Now let us take the ratio of  $T_2$  by  $T_1$  what I am going to get,  $P_2 V_2$  bar by  $P_1 V_1$  bar okay, so let us put equation 2 in equation 1 so  $T_2$  by  $T_1$  I am going to replace there so  $P_2 V_2$  bar by  $P_1 V_1$  bar is equal to  $V_1$  bar by  $V_2$  bar to the power  $\gamma - 1$ . Now let us take  $V_2$  to the power  $\gamma - 1$  on the left and now let us cross multiply ok. What we are going to get,  $P_2 V_2$  bar into  $V_2$  to the power  $\gamma - 1$  is equal to  $P_1 V_1$  bar  $V_1$  to the power  $\gamma - 1$  or  $\gamma - 1$  correspondence like  $P_2 V_2$  to the power  $\gamma$  equal to  $P_1 V_1$  to the power  $\gamma$ .

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Adiabatic Work (No Heat Exchange Allowed)


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$$P_1 V_1^\gamma = P_2 V_2^\gamma = PV^\gamma \Rightarrow P = \frac{P_1 V_1^\gamma}{V^\gamma} \quad (1)$$

$$W = - \int_{V_1}^{V_2} P dV$$

$$= - P_1 V_1^\gamma \int_{V_1}^{V_2} \frac{dV}{V^\gamma}$$

$$= - \frac{P_1 V_1^\gamma}{-\gamma + 1} \left[ V_2^{1-\gamma} - V_1^{1-\gamma} \right]$$

$$= \frac{P_1 V_1^\gamma}{1-\gamma} \left[ V_2^{1-\gamma} - V_1^{1-\gamma} \right]$$


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So for an isothermal case remember what was the case,  $PV$  is equal to constant but for an adiabatic case we get  $PV$  to the power gamma equal to constant that is the difference between isothermal and adiabatic case. So now that we have derived that let us write it here, so what we have seen that  $P_1 V_1$  to the power gamma is equal to  $P_2 V_2$  to the power gamma equal to  $PV$  to the power gamma in general. Once that is the case we can always write then pressure as, in general pressure we can express as  $P_1 V_1$  to the power gamma divided by  $V$ . Now we are going to calculate the  $PV$  way the work done so which is minus  $PdV$  integration  $V_1$  to  $V_2$  and we are going to replace the pressure from this expression which is.

Now  $P_1 V_1$  to the power gamma is a constant right, and this expression is wrong so  $P V$  to the power gamma so it will be gamma ok. You see from this one I am getting it is  $P_1 V_1$  to the power gamma by  $V_2$  the power gamma that was the thing so now I am getting the correct one  $dV$  by  $V_2$  the power gamma that is I now. Now minus  $P_1 V_1$  to the power gamma when I integrate the Gamma what I am going to get, minus Gamma plus 1 by  $V_2$  the power 1 minus Gamma going from  $V_1$  to  $V_2$ , so minus  $P_1 V_1$  to the power gamma by 1 minus Gamma  $V_2$  to the power 1 minus Gamma minus 1 minus gamma  $V_1$  to the power 1 minus Gamma.

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### Adiabatic Work (No Heat Exchange Allowed)

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<p>For adiabatic process, <math>dq = 0</math>,</p> <p>So, <math>dU = dW</math></p> <p>or, <math>dU = -p dV</math></p> <p>Integrating both sides, we get</p> $\int_{U_1}^{U_2} dU = - \int_{V_1}^{V_2} p_{ext} dV$ <p>or, <math>\Delta U = U_2 - U_1 = w</math></p> <p>But, <math>\int_{U_1}^{U_2} dU = C_V \int_{T_1}^{T_2} dT = C_V(T_2 - T_1)</math></p> <p>So, <math>w = C_V(T_2 - T_1)</math></p>	$dU = -p dV$ <p>Or, <math>C_V dT = -p dV = -\frac{RT}{V} dV</math> [ since <math>p dV = nRT</math> or <math>pV = RT</math> for <math>n = 1</math> ]</p> <p>Or, <math>\frac{C_V dT}{T} = -\frac{R}{V} dV</math></p> $\int_{T_1}^{T_2} \frac{C_V dT}{T} = -R \int_{V_1}^{V_2} \frac{dV}{V}$ $C_V \ln\left(\frac{T_2}{T_1}\right) = -R \ln\left(\frac{V_2}{V_1}\right) = R \ln\left(\frac{V_1}{V_2}\right)$ <p>Or, <math>\ln\left(\frac{T_2}{T_1}\right)^{C_V} = \ln\left(\frac{V_1}{V_2}\right)^R</math></p>
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And we can see the same expression is written here so we can see that what is the correct one?

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### Adiabatic Work (Derivation Continued)

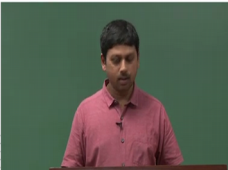
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Or,  $\left(\frac{T_2}{T_1}\right) = \left(\frac{V_1}{V_2}\right)^{R/C_V} = \left(\frac{V_1}{V_2}\right)^{\gamma-1}$ , where  $\gamma = C_P/C_V \rightarrow \gamma = 1 + R/C_V$

Or,  $\left(\frac{P_2 V_2}{P_1 V_1}\right) = \left(\frac{V_1}{V_2}\right)^{\gamma-1}$  Or,  $P_1 V_1^\gamma = P_2 V_2^\gamma = pV^\gamma$  or  $p = CV^{-\gamma}$ , where  $C = p_1 V_1^\gamma$ .

So,  $pV^\gamma = \text{constant}$  and  $p \propto 1/V^\gamma$

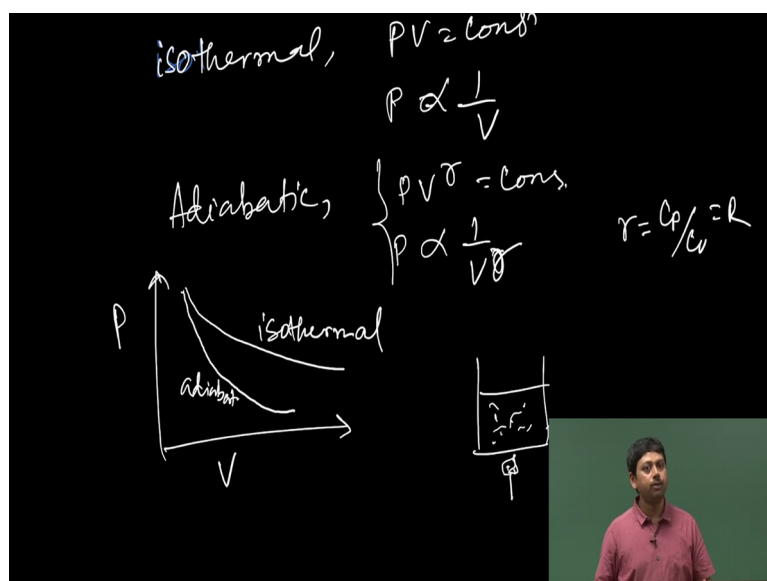
$$W_{adiabatic} = \int_{V_1}^{V_2} -pdV = -C \int_{V_1}^{V_2} V^{-\gamma} dV = -C \left[ \frac{V^{1-\gamma}}{1-\gamma} \right]_{V_1}^{V_2} = -\frac{p_1 V_1^\gamma}{1-\gamma} [V_2^{1-\gamma} - V_1^{1-\gamma}] = C_V (T_2 - T_1)$$



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This is the derivation that we have already done. Okay so this is the now expression for so you see this is same as that of  $C_V T_2$  minus  $T_1$ , do not forget that that is same as  $C_V T_2$  minus  $T_1$ . So if you have this information of temperature you use this  $C_V T_2$  minus  $T_1$  formula, if you do not have the information of temperature you have to use this complicated formula ok.

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So now obviously when you compare between isothermal and adiabatic work, for an isothermal  $PV$  is equal to constant and therefore  $P$  goes as  $1$  by  $V$  right. For adiabatic case adiabatic processes  $PV$  to the power gamma is constant so therefore  $P$  goes as  $1$  by  $V^2$  the power gamma and Gamma is  $C_p$  by  $C_v$  okay which is  $R$  in case of ideal gas and it will be like more than  $1$ , values will be more than  $1$  positive quantity. So when I plot isothermal and adiabatic what do you think the profile will look like? Let us say this is isothermal, will the adiabatic be higher or lower? Just compare these 2 expressions and tell me, in one case it is just going  $1$  by  $V$  and another case it is going  $1$  by  $V^2$  the power gamma.

Since gamma is larger than  $1$ , imagine Gamma is  $2$  so it is decreasing like you know much faster so if gamma is greater than  $1$  then it will be always lower, so this is isothermal and this is adiabatic. Adiabatic work will be always the pressure falls in adiabatic processes will be always lower than isothermal processes. So this is the mathematical reason that one is going by  $1$  by  $V$  and another is going by  $1$  by  $V^2$  the power gamma, what will be the physical reason for this?

The physical reason for this is that in isothermal case let us say you have a gas system and that is going to expand but in order to keep the temperature you have a burner kept or a heat reservoir is kept there so heat is being supplied continuously so because of expansion, pressure falls, but because heat goes inside it does not fall as much. But in adiabatic case system there is no heat going inside so pressure is falling according to the expansion of volume alone, there is no compensation coming from heat outside and that is the reason for

isothermal processes the pressure is less compared to that of adiabatic processes and that is shown by this particular nice looking graph where you have not only 2 dimensions but also 3 dimensions.

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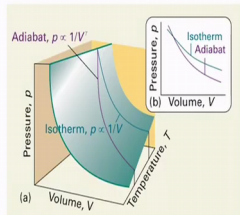
### Comparison between Adiabatic and Isothermal Processes

In isothermal process,  $P \propto \frac{1}{V}$

In adiabatic process,  $P \propto \frac{1}{V^\gamma}$  and  $\gamma > 1$

Therefore pressure falls more rapidly in adiabatic process than in the isothermal process.

The reason is that in the isothermal process, both internal energy and heat do the work (expand the volume). But in adiabatic process, the work is done by the internal energy alone. Therefore, temperature drops.



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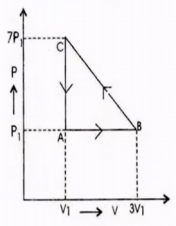
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The 3<sup>rd</sup> dimension says that you know as if there is small change in temperature then how it look like so that is going to different so that one can study, and the reason I have already mentioned.

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### Application of isothermal and adiabatic work: TQ1

Calculate the work done for this cyclic process.



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Okay so now I think we have crossed the time but we are going to discuss a few problems or calculations about how to do sudden works and things like that in tomorrow's class, so we will stop.