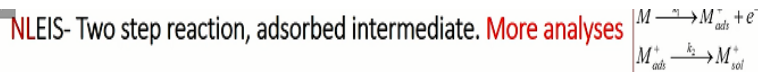


Electrochemical Impedance Spectroscopy
Prof. S. Ramanathan
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Indian Institute of Technology - Madras

Chapter – 45
Rt and Rp Estimation

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1. How long does it take to stabilize?
2. What happens if we take Z data before stabilization?
3. Can we estimate the value of θ_{av} during steady periodic oscillation?
4. Can we estimate nonlinear charge transfer resistance and polarization resistance?



Next can we estimate them Nonlinear charge transfer resistance $R_{t,NL}$ and the polarization resistance.

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NLEIS- Two step reaction. $R_{t,NL}$

Fathima Fasmin, S. Ramanathan, J. Solid State Electrochem. 19 (6) (2015) 1833-1847

$$(R_t)^{-1} = \lim_{\omega \rightarrow 0} \lim_{E \rightarrow E_{ac0}} \frac{i_{E-ac0}}{E_{ac0}}$$

$$(R_{t,NL})^{-1} = \lim_{\omega \rightarrow 0} \lim_{E \rightarrow E_{ac0}} \frac{i_{E-ac0}}{E_{ac0}}$$

$$(R_t)^{-1} = F (b_1 k_1 (1 - \theta_{av}))$$

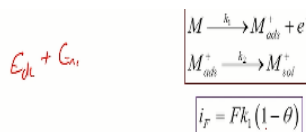
$$\theta = \theta_{av} + \frac{d\theta}{dE} E_{ac0} \sin(\omega t) + \dots$$

$$\lim_{\omega \rightarrow 0} \frac{d\theta}{dE} \rightarrow 0$$

$$\lim_{\omega \rightarrow 0} \theta \rightarrow \theta_{av}$$

$$\lim_{\omega \rightarrow 0} i_F = F k_1 e^{\frac{b_1(E_{ac0} + E_{ac0} \sin(\omega t))}{RT}} (1 - \theta_{av})$$

$$\lim_{\omega \rightarrow 0} i_F = F k_1 e^{\frac{b_1 E_{ac0}}{RT}} [I_0 (b_1 E_{ac0}) + 2I_1 (b_1 E_{ac0}) \sin(\omega t) + \dots] (1 - \theta_{av})$$



So this work is published by our group in 2015 just to refresh your memory charge transfer resistance is R_t and the definition is given here with the inverse of charge transfer resistance right here their limit of E_{ac} tending to 0 and ω tending to infinity that is high frequency limit and small amplitude perturbation the faraday current that will not have any phase offset with respect to the potential.

And that impedance faraday impedance is inverse of the charge transfer resistance I am sorry faraday impedance is the charge transfer resistance under those conditions or faraday admittance is the inverse of charge transfer system. So we write it as Z_f limit $E_{ac} \rightarrow 0$ tending to 0 ω tending to infinity that is charge transfer resistance I can write it as Z_f I can write it sometime as $\frac{dE}{dI}$ I can write it $\frac{dI}{dE}$ for the faraday component.

I can write it as $E_{ac} \rightarrow 0$ $iF_{ac} \rightarrow 0$ they all mean the same thing. Now same faraday impedance at the limit of ω tending to 0 under those conditions also it will be a real value it will not have any imaginary component and that limit the faraday impedance is called polarization resistance Now if I get rid of this constraint and say $E_{ac} \rightarrow 0$ need not be small if I say $E_{ac} \rightarrow 0$ can be an arbitrary value large or small it does not matter.

I would call that as a nonlinear charge transfer resistance. For example this mechanism we have an expression for faraday current only the first reaction contributes to the electrons and therefore we write it as $Fk_1(1-\theta)$ that is the rate of the first reactions multiplied by faraday constant and at infinite frequency we know that $\frac{d\theta}{dE}$ tends to 0 therefore its going to remain at θ_{ss} or actually with θ_{AB} .

But when it is small θ_{ss} and θ we are more or less the same. The rate constant k_1 if I take the derivative its going to be $b_1 k_1$ at that dc potential. So when I vary the potential close to the dc potential I will say $E_{dc} + E_{ac}$ and when it is small amplitude perturbation we know that we can raise the charge transfer resistance as R_t inverse of R_t as $f b_1 k_1 E_{dc} (1-\theta_{ss})$. This we have seen before.

The theta we will write it as average value steady state value + d theta / dE and so on and we know that when omega tends to infinity d theta / dE will tend to 0. Basically omega comes in the denominator of the expression for d theta / dE and theta tends to theta average when omega tends to infinity what happens to IF we can write f k1 1 - theta and expand the k1 as k1 dc multiplied by E power b1 Eac yeah.

Now here I want you to try this instead of writing this in Taylor series and truncating after the first term. Let us write it in Fourier series so we keep the k1 0 e power b1 Edc and call that as k1 dc e power b1 Eac 0 sin omega t that I would like you to expand in Fourier series. So first term it is going to be i0 b1 Eac 0 second term its going to be 2i1 b1 Eac 0 sin omega t + and so on. Right sin 2 omega t or cosine 2 omega t sin 3 omega etc +- all those things will come.

Multiplied by 1 - theta average now we want to look at the IFAC 0 the coefficient of the response at fundamental that is at omega. So we want to look at the coefficient of this and the coefficient of this is 2i1 b1 Eac 0.

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NLEIS- Two step reaction. $R_{t,NL}$

$$\lim_{\omega \rightarrow \infty} i_F = Fk_{1-dc} \left[I_0(b_1 E_{ac0}) + 2I_1(b_1 E_{ac0}) \sin(\omega t) + \dots \right] (1 - \theta_{av})$$

$$\lim_{\omega \rightarrow \infty} i_F = Fk_{1-dc} 2I_1(b_1 E_{ac0}) (1 - \theta_{av})$$

$$(R_{t,NL})^{-1} = \lim_{\omega \rightarrow \infty} \frac{i_F - i_{F0}}{E_{ac0}} = \lim_{\omega \rightarrow \infty} \frac{i_F}{E_{ac0}} = F \left(\frac{2I_1(b_1 E_{ac0})}{E_{ac0}} k_{1-dc} (1 - \theta_{av}) \right)$$



$$(R_{t,NL})^{-1} = F \left(\frac{2I_1(b_1 E_{ac0})}{E_{ac0}} k_{1-dc} (1 - \theta_{av}) \right)$$

Heuristic

Use a correction factor of $[2I_1(b_1 E_{ac0}) / E_{ac0}]$ for each k_{1-dc} and use θ_{1-2V} instead of θ_{1-2S}

$$\begin{array}{c} M \xrightarrow{k_1} M_{adl}^+ + e^- \\ M_{adl}^+ \xrightarrow{k_2} M_{red}^+ \end{array}$$

$$(R_t)^{-1} = F(b_1 k_{1-dc} (1 - \theta_{av}))$$

So if I write this I just copied from the previous page the iF value given by the Fourier series the iF coefficient corresponding to omega is given by Fk1dc multiplied by 2I1modified Bessel function of the first order b1 Eac 0 and then multiply by 1 - theta average this is going to be the iF

E_{ac0} if I write iF as $iF_{dc} + iF_{ac0}$ or $iF_1 \sin \omega t + iF_2 \sin 2\omega t$ sin or cosine it does not matter $\phi_2 \phi_1$ etc.

Of course once they know the value of this coefficient I can take the ratio and call that as the charge transfer resistance or inverse of the charge transfer resistance because I do not make the assumption that E_{ac0} is small I am just using the Fourier series of the full e power a sin theta I call this as the nonlinear charge transfer resistance and that is given by $2I_1 b_1 E_{ac0}/E_{ac0}$. So earlier we would have gotten an expression saying R_t inverse is $f b_1 k_1 dc 1 - \theta_{ss}$.

Now when E_{ac} is small θ_{avg} is going to be roughly θ_{ss} and when E_{ac} is small this value I_1 of x under the limit extending to 0 is going to be $x/2$ this value when E_{ac} is small it is going to be 2 multiplied by $b_1 E_{ac0}/2$ and of course if you take the denominator this term given in the blue colour will tend to b_1 at the limit of E_{ac0} being a small value and θ_{avg} will tend to θ_{ss} when E_{ac0} is small.

So you can see this expression will reduce to the familiar expression when E_{ac0} is small. When E_{ac0} is large this is what you should use to predict the charge transfer resistance. So what that means is you can write the expression for charge transfer resistance as you would do normally okay just replace θ_{ss} with θ_{avg} and replace all the kinetic parameters b_i with E_{ac0}/E_{ac0} .

So 2 times I_1 of $b_1 E_{ac0}/E_{ac0}$ I want you to remember this this is heuristic meaning its based on empirical evidence and based on some expansions for certain mechanisms and in our experience it seems to work but this is not proven analytically. Now basically what it means is use a correction factor of $2 I_1 b_i E_{ac0}/E_{ac0}$ this tells how it is going to depend on the E_{ac0} for each of these b_i okay and use θ_i average θ_1 average θ_2 average.

You may have multiple intermediate states for all of them you have to use the θ_{avg} in sort of θ statistics. Then so far we have been able to predict the RTNL correctly for practically all the mechanisms we have tried. For all the kinetic parameters we have tried.


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Previous class

- More analyses of the two step reaction with an adsorbed intermediate
 - How long should we wait for steady periodic response?
 - Where does θ_{av} settle? How does it depend on E_{act} , f ?
 - What do we get, if we do not wait for steady periodic response?
 - Can we estimate the values of $R_{t,NL}$ without solving the ODE?

Today

- Two step reaction with an adsorbed intermediate
 - Can we estimate the values of $R_{p,NL}$ without solving the ODE?





So earlier we saw that you can estimate the value of theta average when we apply sinusoidal potential it settles after some time we found that the wait time necessarily its independent of frequency it depends on the amplitude of perturbation it of course depends on the kinetic parameters also. We found that we can estimate the average value of the theta its independent of the frequency again.

But at least under high frequency limit we can show that it can be returned in terms of modified Bessel functions we also saw that if you do not wait for sufficient time and acquire data then the results are going to be different especially at larger amplitudes and we just mentioned that the results taken under those conditions will not be KKT complaint. We did not show that but we just mentioned that.

And we also saw that its possible to estimate the charge transfer resistance and that large amplitude perturbation using a heuristic method that is not proven to be valid under all conditions but it seems to be working correctly for at least the mechanism so you tried and the parameters we have tried. Now can we estimate the polarization resistance under large amplitude perturbation again without solving the ODE.

Of course if you can generate the spectrum then we can just look at what happens to the faraday impedance when we go to lower and lower frequencies and when it settles we say that is a

polarization resistance. But if I have current if I have dc current as a function of dc potential it is possible to estimate the polarization resistance. If I have an expression for dc current it is possible to estimate the polarization resistance. And I want to describe that part.

(Refer Slide Time: 11:37)

NLEIS- Two step reaction. $R_{p,NL}$

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$$(R_p)^{-1} = \lim_{\omega \rightarrow 0} \lim_{E_{ac} \rightarrow 0} \frac{j\omega}{F_{ac}(\omega)}$$

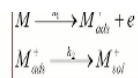
$$(R_{p,NL})^{-1} = \lim_{\omega \rightarrow 0} \frac{j\omega}{E_{ac}(\omega)}$$

Linearize: i.e. Write $E = E_{dc} + E_{ac}$, Expand i_f in Taylor series, truncate after the first term ...

$$i_f = i_{f,dc} + i_{f,ac} \approx Fk_1(1-\theta_{ss}) + \left[Fb_1k_1(1-\theta_{ss}) - Fk_1 \frac{d\theta}{dE} \right] E_{ac} + \dots$$

$$\frac{d\theta}{dE} = \frac{b_1k_1(1-\theta_{ss})}{k_1 + k_2 + j\omega\Gamma}$$

$$i_{f,ac} \approx \left[Fb_1k_1(1-\theta_{ss}) - Fk_1 \frac{d\theta}{dE} \right] E_{ac}$$



$$i_f = Fk_1(1-\theta)$$



Now charge transfer resistance is faraday impedance at high frequencies and under small amplitude perturbation. Polarization resistance is faraday impedance at low frequencies at the limit of omega tending to 0 that impedance faraday impedance is called polarization resistance. Now if I get rid of the constraint saying $E_{ac} \rightarrow 0$ has to be very small if I say $E_{ac} \rightarrow 0$ can be large I still look at the faraday impedance at omega tending to 0.

That we can call us nonlinear polarization resistance when we want to calculate the polarization resistance we can write it as $E_{dc} + E_{ac} \sin \omega t$ that is the potential we have an expression for current expanding Taylor series truncate after the first term because it is a small amplitude perturbation and these are actually approximate signs. Faraday current is split into dc and ac component and that I can write as dc component here and ac component here.

We had done this before and we know the expression for $d\theta/dE$ comes as $b_1k_1(1-\theta_{ss})$ for this particular mechanism and the denominator has $k_1 + k_2 + j\omega\Gamma$ when omega tends to 0 $d\theta/dE$ we can write it easily and $i_{f,ac}$ is the second term here so we take

the second term this actually the charge transfer resistance or the inverse of charge transfer resistance this is the addition to estimate the value of polarization resistance.

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NLEIS- Two step reaction. $R_{p,NL}$

$$i_{F-\omega} = F \left[b_1 k_{1-\omega} (1 - \theta_{ss}) - F k_{1-\omega} \frac{d\theta}{dE} \right] E_{\omega}$$

$$(R_p)^{-1} = \lim_{\omega \rightarrow 0} \lim_{E \rightarrow 0} \frac{i_{F-\omega}}{E} = F \left(b_1 k_{1-\omega} (1 - \theta_{ss}) - k_{1-\omega} \lim_{\omega \rightarrow 0} \frac{d\theta}{dE} \right)$$

$$\lim_{\omega \rightarrow 0} \frac{d\theta}{dE} = \frac{b_1 k_{1-\omega} (1 - \theta_{ss})}{k_{1-\omega} + k_2}$$


$$(R_p)^{-1} = F \left(b_1 k_{1-\omega} (1 - \theta_{ss}) - k_{1-\omega} \frac{b_1 k_{1-\omega} (1 - \theta_{ss})}{k_{1-\omega} + k_2} \right)$$

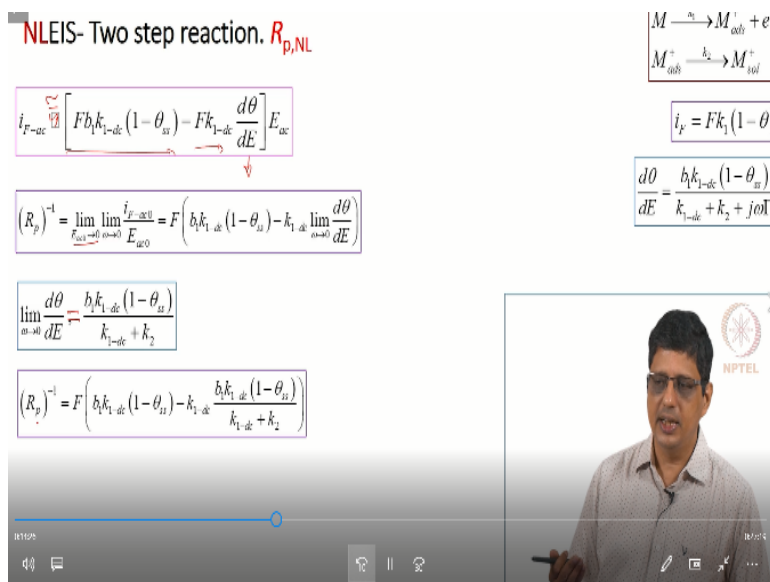
$$M \xrightarrow{k_1} M_{ads} + e$$

$$M_{ads} \xrightarrow{k_2} M_{red}$$

$$i_p = F k_1 (1 - \theta)$$

$$\frac{d\theta}{dE} = \frac{b_1 k_{1-\omega} (1 - \theta_{ss})}{k_{1-\omega} + k_2 + j\omega\tau}$$





So we can find the inverse of polarization resistance as R_p inverse and under the assumption that this is small signal condition ω tending to 0 this is independent of ω this is independent of ω only this depends on ω and at ω tending to 0 this can be simplified easily the $j\omega\tau$ goes out therefore we can write an expression for polarization resistance for a given reaction with the assumption that mass transfer is.

And solution resistance is negligible. Now for small amplitude perturbation when polarization resistance is defined it is easy for us to calculate for any given mechanism. We can calculate the impedance at any frequency therefore we can calculate the importance at 0 frequency or infinite frequency no problem.

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NLEIS- Two step reaction. $R_{p,NL}$

Calculate R_p and $R_{p,NL}$ from Potentiodynamic polarization curve: i_{dc} vs. E_{dc}

$i_{F-dc} = f(E_{dc})$

$\frac{di_{F-dc}}{dE_{dc}} = f'(E_{dc}) = \lim_{\Delta E_{dc} \rightarrow 0} \lim_{\Delta i_F \rightarrow 0} \frac{\Delta i_F}{\Delta E_{dc}} = \lim_{E_{dc} \rightarrow 0} \lim_{i_F \rightarrow 0} \frac{i_{F-dc0}}{E_{dc0}} = (R_p)^{-1}$

- i.e. Take derivative of i_{dc} vs. E_{dc} , and inverse of that is R_p
- This can even be a numerical derivative,
- The functional form of i_{dc} vs. E_{dc} need not be known

Diard, J.-P., Le Gorrec, B., & Montella, C. (1997) *J Electroanal Chem*, 432, 27
 Diard, J.-P., Le Gorrec, B., & Montella, C. (1997) , 41-52.
 Diard, J.-P., Le Gorrec, B., & Montella, C. (1997) 53-62.


$\lim_{\omega \rightarrow 0} \frac{i_{F-dc0}}{E_{dc0}} = (R_{p,NL})^{-1} = \sum_{k=0}^{\infty} b_k f^{(2k+1)} E_{dc0}^{(2k)}$

$b_0 = 1$
 $b_k = -\frac{b_{k-1}}{4k(k+1)}$, for $k \geq 1$

$f^{(2k+1)}$ means $\frac{d^{(2k+1)} f}{dE^{(2k+1)}}$

$M \xrightarrow{k_1} M_{ad} + e$
 $M_{ad} \xrightarrow{k_2} M_{red}$

$i_F = Fk_1(1-\theta)$



What happens when you have a large amplitude perturbation okay we can estimate the dc current for a given dc potential we can estimate the dc current value the faraday current. Okay so basically its this expression with the concern that potential e is E_{dc} . So in general for any mechanism you would rate as a function of E_{dc} in this case we know what the function is what the expression is.

When I make small changes I can say that result in change small changes ΔE the resulting changes ΔI and if you make the change slowly that is ω tending to 0 that is inverse of polarization resistance. I can write it in different notations I can write it as $\Delta I / \Delta E$ I can write it as I_{FAC0} by E_{ac0} sometimes people also write it as prime indicating it is a derivative. If I have experimental data I vs E_{dc} let us see we do not even know the mechanism.

We do not have an expression for this still if we take the derivative here using numerical method I can take 2 data points find the $\Delta I / \Delta E$ that this find the value of the slope that slope is going to be di/de and that is going to be inverse of the polarization resistance. So this curve which we call as potential dynamic polarization the slope of the curve at any given point that has to be equal to the inverse of polarization resistance.

So if I acquire impedance data here that data may come like this it does not matter the low frequency limit there is R_p and that should be equal to the inverse of the slope of that curve. So

you can take numerical derivative if you know the form of this faraday impedance if you know the dependence of faraday current on the potential you can take derivative and at the limit of ω tending to 0 you can get the expression anyway.

So it does not require us to solve ODE anyway for linearized cases we do not need to solve ODEs for this kind of expressions. In 1997 there were a series of publications by professor Diard Gorrec where they have shown even for large amplitude perturbation one can calculate the polarization resistance as long as we know the expression iF relating to E_{dc} so when we get rid of the constraint that ΔE has to be is very small.

We can defend that as R_{pNL} and the inverse of R_{pNL} is given by this series this is the series that they have suggested and if i remember right they have used a notation like this. It does not matter we use a notation saying $R_{subscript NL}$ and they have given it by a series and out here f' prime indicates the first derivative. f'' double prime will indicate this secondary derivative and $f^{(n)}$ indicates the n th derivative with respect to potential okay.

So $f^{(2k+1)}$ means it is actually not $f^{(n)}$ derivative or $2k+1$ th derivative b_0 is given as 1 and subsequent b values b_1 b_2 b_3 etc can be obtained using this formula. So if I know b_0 I can calculate b_1 and once I know b_1 I can calculate b_2 and so on. So this is the expression they have derived so which means I need to know the functional form I need to have the ability to get n number of derivatives from this.

And when I take more and more numbers hopefully I will get more and more accurate estimate and beyond certain limit you may not need this anymore. Because this coefficient b_k will become smaller and smaller when we go to larger and larger values of k .

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We want to say upper case D with subscript e would represent d/dE and that is the derivative. If I write dE power n that is the nth derivative. So I want to use this expression operative notation here and substitute here second I want you to recognize or remember this the modified Bessel function of nth order is written like this $1/2 Z$ whole power n multiplied by this summation and this of course depends on k value and n value.

You can expand this series and you can compare with the terms here they will basically mean that you write R_{pNL} as $2/E_{ac}$ 11 of the operator DE E_{ac} 0 is multiplying the DE applied on to

the function f of E . So we have a function $I_{\text{faraday current}} = f$ of E_{dc} if I use this notation it's a compact notation this will tell me the nonlinear polarization resistance. So this part tells them low frequency limit current.


So this tells me $i_{\text{fac } 0}$ when ω tends to infinity I am sorry when ω tends to be 0 and $E_{\text{ac } 0}$ is arbitrary. When $E_{\text{ac } 0}$ is small I_1 of $E_{\text{ac } 0}$ DE is approximately $E_{\text{ac } 0} \text{ DE}/2$ and of course when I multiply by $2/E_{\text{ac } 0}$ I just get the derivative DE f polarization resistance inverse of that is di/dE this simplifies to the known expression when $E_{\text{ac } 0}$ is small. So if you have I_{dc} at various E_{dc} s then you can calculate the polarization resistance.

Nonlinear polarization resistance either by taking n number of derivatives first derivative second derivative and so on around the DC potential where we want to estimate the polarization resistance. So in corrosion polarization resistance basically means at the open circuit potential but in general polarization resistance it can be acquired at any DC bias. So it means the low frequency limit of the impedance spectra with or without dc bias.

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NLEIS- Two step reaction. $R_{p,NL}$

Analytical Calculations: Example



C_1, C_2, \dots

$$\begin{aligned} M &\xrightarrow{k_1} M_{\text{ad}} + e \\ M_{\text{ad}} &\xrightarrow{k_2} M^+ \end{aligned}$$

$$i_f = Fk_1(1-\theta)$$

$$\tau \frac{d\theta}{dt} = k_1(1-\theta) - k_2\theta$$

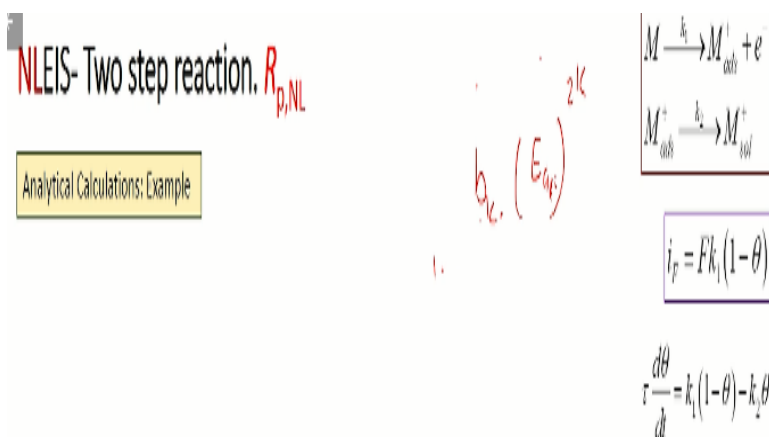
Now I have an expression i_f is given as $Fk_1(1-\theta)$ at very low frequencies it is still valid just that I had to replace the k_1 with k_1^{DC} θ with θ_{ss} and i_f with i_f^{DC} . For this it's possible to an analytical expression for the nonlinear polarization resistance we do not need to actually

generate various i_F values at various DC potentials and take the numerical derivative it is possible to do it you can take the numerical derivative.

But numerical derivative tends to be noisy and especially when you go for multiple order derivative 5th order 10th order 20th order when you go to a higher order derivatives you would have incorrect result. Roundoff error will play a role now if we take E_1 E_2 E_3 and make them very small if we make the interval very small I get a little more accuracy. But again if you make it extremely small I will have round off errors.

So numerical derivative is something we should avoid if you can if it is possible to get an analytical expression it is better to use it. Now I also want you to notice the following.

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We have the expression which goes E_{ac} 0 power $2k$ and then b_{k_2} when you use large amplitude perturbation you need to use a summation $k=0$ to some value let us say 10 beyond that you can neglect the terms. If I use larger perturbation I should go to larger values here if you go to even larger E_{ac} we can go to larger value here before I can neglect the remaining terms. So what level do we stop we cannot go to infinite summation.

Well actually there are some software which allow you to calculate those values or estimate those values. But in general if you are just writing your own program it is easy to calculate up to some level and then stop.

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NLEIS- Two step reaction. $R_{p,NL}$

Analytical Calculations: Example

$$\theta = \theta_{|E=E_0} + \frac{\Delta E}{1!} \left(\frac{d\theta}{dE} \right)_{|E=E_0} + \frac{(\Delta E)^2}{2!} \left(\frac{d^2\theta}{dE^2} \right)_{|E=E_0} + \frac{(\Delta E)^3}{3!} \left(\frac{d^3\theta}{dE^3} \right)_{|E=E_0} + \dots$$

$$\theta = \theta_{|E=E_0} + \frac{E_{ac}}{1!} \left(\frac{d\theta}{dE} \right)_{|E=E_0} + \frac{(E_{ac})^2}{2!} \left(\frac{d^2\theta}{dE^2} \right)_{|E=E_0} + \frac{(E_{ac})^3}{3!} \left(\frac{d^3\theta}{dE^3} \right)_{|E=E_0} + \dots$$

$$M \xrightarrow{k_1} M_{ad}^+ + e^-$$

$$M_{ad}^+ \xrightarrow{k_2} M_{red}^+$$

$$i_f = Fk_1(1-\theta)$$

$$\tau \frac{d\theta}{dt} = k_1(1-\theta) - k_2\theta$$

I want to expand the theta in Taylor series theta around E_0 $\frac{d\theta}{dE}$ $\frac{d^2\theta}{dE^2}$ $\frac{d^3\theta}{dE^3}$ this is how we would expand in Taylor series we would expand around x_0+h this is similar to that. Now $\frac{d\theta}{dE}$ we can replace by E_{ac} so I can rewrite that as E_{ac} E_{ac}^2 E_{ac}^3 and so on. Now I cannot say that E_{ac} is small I cannot say that I can neglect these terms.

I actually want to use this to calculate the nth derivative of current with respect to potential in order to calculate this I need to be able to calculate the nth derivative of theta with respect to potential and we are doing some algebra here just so that we can get the expression for nth derivative it is lengthy but again its not particularly complex its not particularly difficult.

(Refer Slide Time: 27:21)

NLEIS- Two step reaction. $R_{p,NL}$

Analytical Calculations: Example

$$\theta = \theta_{ss} + \frac{E_{ac}}{1!} \left(\frac{d\theta}{dE} \right) + \frac{(E_{ac})^2}{2!} \left(\frac{d^2\theta}{dE^2} \right) + \frac{(E_{ac})^3}{3!} \left(\frac{d^3\theta}{dE^3} \right) + \dots$$

$$k_i = k_{i,dc} + \frac{E_{ac}}{1!} \left(\frac{dk_i}{dE} \right) + \frac{(E_{ac})^2}{2!} \left(\frac{d^2k_i}{dE^2} \right) + \frac{(E_{ac})^3}{3!} \left(\frac{d^3k_i}{dE^3} \right) + \dots$$



$$i_F = Fk_1(1-\theta)$$

$$\tau \frac{d\theta}{dt} = k_1(1-\theta) - k_2\theta$$

I can write the k_1 in this case k_2 is a constant. In general I can write k_i as k_i at dc and then the derivative 1st derivative 2nd derivative nth derivative multiplied correspondingly by del E del E square and del E cube divided by 1 factorial 2 factorial 3 factorial and so on.

(Refer Slide Time: 27:21)

NLEIS- Two step reaction. $R_{p,NL}$

Analytical Calculations: Example

$$\theta = \theta_{ss} + \frac{E_{ac}}{1!} \left(\frac{d\theta}{dE} \right) + \frac{(E_{ac})^2}{2!} \left(\frac{d^2\theta}{dE^2} \right) + \frac{(E_{ac})^3}{3!} \left(\frac{d^3\theta}{dE^3} \right) + \dots$$

$$k_i = k_{i,dc} + \frac{E_{ac}}{1!} \left(\frac{dk_i}{dE} \right) + \frac{(E_{ac})^2}{2!} \left(\frac{d^2k_i}{dE^2} \right) + \frac{(E_{ac})^3}{3!} \left(\frac{d^3k_i}{dE^3} \right) + \dots$$

$$1 \frac{d\theta}{dt} = 1 \frac{d\theta}{dE} \frac{dE}{dt} = 1 \frac{d\theta}{dE} \frac{d(E_{ac} e^{j\omega t})}{dt} = j\omega \frac{d\theta}{dE} (E_{ac} e^{j\omega t}) = j\omega (E_{ac}) \frac{d\theta}{dE}$$



$$i_F = Fk_1(1-\theta)$$

$$\tau \frac{d\theta}{dt} = k_1(1-\theta) - k_2\theta$$

Again I can be replace that by $k_{i,dc}$ and the 1st derivative is going to give me b_1 2nd derivative is going to give me b_2 square this is actually simple. So basically I want to take this expression I find the value of $d\theta/dE$ I want to find the value of $d^2\theta/dE^2$ I want to find the value of $d^3\theta/dE^3$. So I can write this easily expanded I can explain this in this Taylor series.

I can expand this in the Taylor series this is the constant on the left side we have seen this before very early when we look at the linear decision. We would have written $\gamma \frac{d\theta}{dt}$ as $\gamma \frac{d\theta}{dE}$ multiplied by $\frac{dE}{dt}$ and we would get $j\omega \gamma E_{ac} \frac{d\theta}{dE}$. We have done this n number of times so we will just continue with this.

(Refer Slide Time: 28:43)

NLEIS- Two step reaction. $R_{p,NL}$

Handwritten notes: $\frac{d\theta}{dt} E$, $\frac{d^2\theta}{dt^2} E^2$, $\frac{d^3\theta}{dt^3} E^3$

Analytical Calculations: Example

$$\theta = \theta_{ss} + \frac{E_{ac}}{1!} \left(\frac{d\theta}{dE} \right) + \frac{(E_{ac})^2}{2!} \left(\frac{d^2\theta}{dE^2} \right) + \frac{(E_{ac})^3}{3!} \left(\frac{d^3\theta}{dE^3} \right) + \dots$$

$$k_1 = k_{10} + \frac{E_{ac}}{1!} b_1 k_{10} + \frac{(E_{ac})^2}{2!} b_1^2 k_{10} + \frac{(E_{ac})^3}{3!} b_1^3 k_{10} + \dots$$

$$j\omega \tau \frac{d\theta}{dt} = j\omega \tau \left(E_{ac} \frac{d\theta}{dE} \right)$$

$$j\omega \tau \left(E_{ac} \frac{d\theta}{dE} \right) = \left(k_{10} + \frac{E_{ac}}{1!} b_1 k_{10} + \frac{(E_{ac})^2}{2!} b_1^2 k_{10} + \frac{(E_{ac})^3}{3!} b_1^3 k_{10} + \dots \right) \times \left(1 - \theta_{ss} \right) \left(\frac{E_{ac}}{1!} \frac{d\theta}{dE} \right) + \frac{(E_{ac})^2}{2!} \frac{d^2\theta}{dE^2} + \frac{(E_{ac})^3}{3!} \frac{d^3\theta}{dE^3} + \dots$$

$$-k_2 \times \left(\theta_{ss} + \frac{E_{ac}}{1!} \left(\frac{d\theta}{dE} \right) + \frac{(E_{ac})^2}{2!} \left(\frac{d^2\theta}{dE^2} \right) + \frac{(E_{ac})^3}{3!} \left(\frac{d^3\theta}{dE^3} \right) + \dots \right)$$

Coefficient of E_{ac} $\frac{d\theta}{dE} = \frac{b_1 k_{10} (1 - \theta_{ss})}{k_{10} + k_2 + j\omega \tau}$

Chemical reactions:

$$M \xrightarrow{k_1} M_{adi} + e$$

$$M_{adi} \xrightarrow{k_2} M_{tot}$$

Equation: $i_F = Fk_1(1 - \theta)$

Derivative: $\tau \frac{d\theta}{dt} = k_1(1 - \theta) - k_2 \theta$

NPTEL logo

Now if I take this equation and substitute without linearization. So I substitute for k_1 from the expression here I substitute for θ from the expression here k_2 is left as it is multiplied by θ . Now this expression is valid under all E_{ac} that means I should be able to isolate the coefficients of E_{ac} and equate them. Coefficients of E_{ac} square and equate them E_{ac} cube and equate them.

So I had to find what terms here will give me E_{ac} square what terms will give me E_{ac} cube right now I am just focusing on E_{ac} square now look at this this coefficient its a coefficient of E_{ac} square this is a coefficient of E_{ac} square this is a coefficient of E_{ac} square. All the cubic terms are given in blue colour but not only that if I multiply term here with term here I will get E_{ac} square.

If you multiply this term with this term I will get cube if you multiply this term with the term here I will get a cube so I had to look at it carefully coefficient of E_{ac} alone it is pretty straight

forward it can be taken from this multiplied by 1 -theta this multiplied by factor here this multiplied by the factor here and of course the factor here.

And it would give us what we got earlier d theta/ dE is going to be given by the expression we are familiar with j omega gamma.

(Refer Slide Time: 30:30)

NLEIS- Two step reaction. $R_{p,NL}$

Analytical Calculations: Example

Coefficient of $(E_{ac})^2$

$$j\theta\tau(E_{ac})\frac{d\theta}{dE} = \left[\left(\frac{k_{1,ss}}{1!} + \frac{E_{ac}}{1!} \frac{b_1 k_{1,ss}}{1!} + \frac{(E_{ac})^2}{2!} b_1^2 k_{1,ss} + \frac{(E_{ac})^3}{3!} b_1^3 k_{1,ss} + \dots \right) \right. \\ \left. \times \left(1 - \theta_{ss} - \frac{E_{ac}}{1!} \left(\frac{d\theta}{dE} \right) - \frac{(E_{ac})^2}{2!} \left(\frac{d^2\theta}{dE^2} \right) - \frac{(E_{ac})^3}{3!} \left(\frac{d^3\theta}{dE^3} \right) + \dots \right) \right. \\ \left. - \frac{k_2}{1!} \times \left(\theta_{ss} + \frac{E_{ac}}{1!} \left(\frac{d\theta}{dE} \right) + \frac{(E_{ac})^2}{2!} \left(\frac{d^2\theta}{dE^2} \right) + \frac{(E_{ac})^3}{3!} \left(\frac{d^3\theta}{dE^3} \right) + \dots \right) \right]$$

$$\theta = \left[\frac{-b_1 k_{1,ss}}{1!} \left(\frac{d\theta}{dE} \right) - \frac{k_2}{2!} \left(\frac{d^2\theta}{dE^2} \right) \right] + \left[(1 - \theta_{ss}) \frac{b_1^2 k_{1,ss}}{2!} - \frac{k_2}{2!} \left(\frac{d^2\theta}{dE^2} \right) \right]$$

$$\left[\frac{k_{1,ss} + k_2}{2!} \right] \left(\frac{d^2\theta}{dE^2} \right) = \left[\frac{-b_1 k_{1,ss}}{1!} \left(\frac{d\theta}{dE} \right) \right] + \left[(1 - \theta_{ss}) \frac{b_1^2 k_{1,ss}}{2!} \right] - \left[\frac{k_2}{2!} \left(\frac{d^2\theta}{dE^2} \right) \right]$$

Rt and Rp estimation-mod1lec45


Chemical reactions:

$$M \xrightarrow{k_1} M_{ox}^+ + e$$

$$M_{ox}^+ \xrightarrow{k_2} M_{red}^+$$

Steady state approximation:

$$i_F = Fk_1(1 - \theta)$$

$$\tau \frac{d\theta}{dt} = k_1(1 - \theta) - k_2\theta$$


Now if you take the coefficient of Eac square we saw that we had to look at it carefully now I have arranged them or coloured them so that the same coloured numbers if you multiply it you would get Eac square the pink colour, blue colour factors and violet colour factors they would give you Eac square here the red colour factor would give you the Eac square so I can group them here other times it will give you higher order terms.

So we are not going to worry about them on the left side there is no Eac square so it is going to be a 0 on the right side I would get b1 k1 dc multiplied by d theta/d E we are taking the Eac square term out and then grouping the remaining terms k1 dc multiplied by 1/2 factorial d square theta/d E square with of course the negative side and the violet colour b1 square k1 dc and 1-theta SS and of course k2 multiplied by d square theta/d E square.

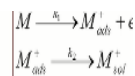
Now I can rearrange you can rearrange and write basically d square theta/d E square in terms of other parameters known parameters.

(Refer Slide Time: 31:54)

NLEIS- Two step reaction. $R_{p,NL}$

Analytical Calculations: Example

Coefficient of $(F_{20})^2$



$$i_F = Fk_1(1-\theta)$$

$$\tau \frac{d\theta}{dt} = k_1(1-\theta) - k_2$$

$$\theta = \left[\frac{-b_1 k_{12}}{1!} \left(\frac{d\theta}{dE} \right) + \left[\frac{-k_{22}}{2!} \left(\frac{d^2\theta}{dE^2} \right) + \left(1-\theta \right) \frac{b_1^2 k_{12}}{2!} \right] - \left[\frac{k_2}{2!} \left(\frac{d^2\theta}{dE^2} \right) \right] \right]$$

$$\left[\frac{k_{12} + k_2}{2!} \right] \left(\frac{d^2\theta}{dE^2} \right) = \left[\frac{-b_1 k_{12}}{1!} \left(\frac{d\theta}{dE} \right) + \left(1-\theta \right) \frac{b_1^2 k_{12}}{2!} \right] -$$

$$\frac{d^2\theta}{dE^2} = \frac{2!}{(k_{12} + k_2)} \left[\frac{b_1^2 k_{12}}{2!} - \left(\frac{b_1 k_{12}}{2!} \right) \frac{1}{0!} + \left(\frac{b_1 k_{12}}{1!} \right) \frac{d\theta}{dE} \right]$$

$$\frac{d^n \theta}{dE^n} = \frac{n!}{(k_{12} + k_2)} \left[\frac{(b_1^n) k_{12}}{n!} - \sum_{m=0}^{n-1} \left[\frac{(b_1^{(m+1)} k_{12})}{(m+1)!} \frac{d^{(n-m-1)} \theta}{dE^{(n-m-1)}} \right] \frac{1}{(n-m-1)!} \right] \text{ for } n > 0$$



I can write it as $d^2\theta/dE^2$ is $2 \text{ factorial}/k_1+k_2$ it is going to depend on $d\theta/dE$. Now if you look at it is independent of ω well it is not really independent of ω . ω does not appear explicitly here $d\theta/dE$ depends on ω and of course. We are writing all this under the limit of ω tending to 0. Now this depends on θ it depends on the $d\theta/dE$ and the remaining terms are arranged.

Here I am going to regroup them slightly differently but it is possible to show that in general you can do this for the 3rd order once you do it for 1st 2nd 3rd you can look at that and tell that for any n th order I can write it as $n \text{ factorial}/k_1 + k_2$. b_1 raised to the power n $k_1 d^n \theta / 2 \text{ factorial}$ for this particular mechanism and this entire term on the bracket can be written as a summation with m going from 0 to $n-1$ etc b_1 power $n+1$ $k_1 d^n \theta$ this derivative basically goes from a constant.

So if I say dn/dE of θ $n=0$ it means there is no derivative just functioning θ $n=1$ it is 1st derivative $n=2$ it is 2nd derivative and so on. So this expression for the mechanism given here for any kinetic parameter as ω tends to 0. This expression is valid and its not that difficult to calculate once you write a program for this it is easy to calculate only problem is when you want to go to many terms then $n \text{ factorial}$ will become a problem each operating system.

And each software for example if I want to get 100 factorial in 64 bit it is possible if you want to get 1000 factorial it would give me a number in scientific notation it would not be accurate. Okay so it will have to truncate after certain number of decimals and it will write in scientific notation it cannot give me the entire number as an integer unless we have some special program for that.

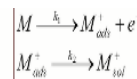
There are some programs which will allow you to change the accuracy but if I use for example Matlab typically there is a limit on what factorial can I get without loss of accuracy. So basically what it means is. I cannot just blindly write a program in C or Fortran or any of such languages and then say go to 1000 factorial or 1000 terms I will have to be aware of the limitation. I have to know under what condition I can get accurate answers and under what conditions we are going to face truncation errors.

(Refer Slide Time: 35:10)

NLEIS- Two step reaction. $R_{p,NL}$

Analytical Calculations: Example

$$\frac{d^n \theta}{dE^n} = \frac{n!}{(k_{-1} + k_2)} \left[\frac{(b_1^n) k_{-1}}{n!} - \sum_{i=0}^{n-1} \left[\frac{(b_1^{(n-i)} k_{-1})}{(n+1)!} \frac{d^{(n-i)} \theta}{dE^{(n-i)}} \frac{1}{(n-m-1)!} \right] \right] \text{ for } n \geq 0$$



$$i_F = Fk_1(1 - \theta)$$

$$F \frac{d\theta}{dE} = k_1(1 - \theta) - k_2 \theta$$

$$\frac{d\theta}{dE} = Fk_1 \left[b_1(1 - \theta) - \frac{d\theta}{dE} \right] = Fk_1 \left[b_1 - \frac{d\theta}{dE} - b_1 \theta \right]$$

$$\frac{d^2 \theta}{dE^2} = Fk_1 \left[\frac{d}{dE} \left(\frac{d\theta}{dE} \right) - Fk_1 b_1 \frac{d\theta}{dE} - Fk_1 \theta \right] = Fk_1 \left[\frac{d^2 \theta}{dE^2} - b_1 \frac{d\theta}{dE} - \theta \right]$$

$$\frac{d^2 \theta}{dE^2} = Fk_1 \left[b_1^2 - \left(\frac{d^2 \theta}{dE^2} + 2b_1 \frac{d\theta}{dE} + b_1^2 \theta \right) \right]$$

$$\frac{d^2 \theta}{dE^2} = Fk_1 \left[b_1^2 - \sum_{n=0}^{\infty} \frac{2}{n!} \left(b_1 \right)^n \left(\frac{d^{(2-n)} \theta}{dE^{(2-n)}} \right) \right] \quad \left(\frac{n}{m} \right) = \frac{n!}{m!(n-m)!}$$

$$\frac{d^n \theta}{dE^n} = Fk_1 \left[b_1^n - \sum_{n=0}^{\infty} \left(\frac{n}{m} \right) \left(b_1 \right)^n \frac{d^{(n-m)} \theta}{dE^{(n-m)}} \right]$$



So in this we can write the expression for the nth derivative of theta now look at the current value di/dE 1st derivative I can easily write we had done this before Fk1dc b1-1- theta d theta/dE. Now I just arranged it slightly differently so that its easy to generalize later instead of rating b1 1-theta I will write it as b1 -d theta/dE-b1 theta. I can take the 2nd derivative if I take the 2nd derivative.

This gives rise to I can say this is $u \cdot v$ so I take the $u \cdot v + u \cdot v$ right k1 dc if I take the derivative again I will write $b_1 \cdot k_1 \cdot dc$ and this of course remains the same here u it remains the same this if I take the derivative b_1 becomes 0 this becomes $d^2 \theta / dE^2$ term is missing and of course this becomes this $-b_1 \cdot d \theta / dE$. I can rearrange it I can multiply and then rearrange $f \cdot k_1 \cdot dc$ I can keep it outside retain it.

B_1 square then b_1 square θ then I have $1 \cdot b_1 \cdot d \theta / dE$ with a negative sign I have another $d \cdot b_1 \cdot d \theta / dE$ with a negative thing so it becomes $-2 \cdot b_1 \cdot d \theta / dE$ and of course I have $-d^2 \theta / dE^2$ and likewise I can write for 3rd derivative 4th derivative etc. But I have rearranged it such that b_1 square is here with a negative sign $d^2 \theta \cdot d \theta \cdot b_1$ square θ so if you look here this is b_1 square this is $2 \cdot b_1$ no coefficient before that it is just 1.

And this can be written as a binomial series $m=0$ to 2 to m b_1 power m and 2 -mth derivative of θ of course here $n \cdot m$ we call it as $n \cdot m$ is $n \text{ factorial} / m \text{ factorial} \cdot n-m \text{ factorial}$ and I can write the n th derivative as $f \cdot k_1 \cdot dc \cdot b_1$ power n . So essentially I can write the n th derivative of the current in terms of the n th and lower derivatives of θ with respect to potential and for each of them we have got an expression before.

So we should be able to calculate and actually if you calculate it for any given $E_{ac} = 0$ you would be able to get the polarization resistance value.

(Refer Slide Time: 38:29)

NLEIS- Two step reaction. $R_{p,NL}$

Analytical Calculations: Example

$$\frac{d^2\theta}{dE^2} = \frac{n!}{(k_{ox} + k_r)} \left[\frac{(b_1^n)k_{ox}}{n!} - \sum_{n=2}^{n-1} \left[\frac{(b_1^{(n-1)}k_{ox})}{(m+1)!} \frac{d^{(n-1)}\theta}{dE^{(n-1)}} \frac{1}{(n-m-1)!} \right] \right] \text{ for } n > 0$$

$$\frac{di}{dE} = Fk_{ox} \left[b_1(1-\theta) - \frac{d\theta}{dE} \right] = Fk_{ox} \left[b_1 - \frac{d\theta}{dE} - b_1\theta \right]$$

$$\tau \frac{d\theta}{dt} = k_1(1-\theta) - k_2\theta$$

$i_f = Fk_1(1-\theta)$

$\frac{2}{E_{ox}} \frac{I_1(E_{ox}, D_F)}{I_f}$

$M \xrightarrow{k_1} M^{+}_{ads} + e$
 $M^{+}_{ads} \xrightarrow{k_2} M^{+}_{sol}$

So essentially we saw that the polarization resistance can be written in terms of the I1 Eac0 de of the I faraday or we can call it as f of Edc. So in order to calculate this we need to be able to take the derivative 1st order 2nd order 3rd order and so on. And in order to calculate those derivatives we need to be able to take the derivative of theta with respect to potential once you have analytical expression for this for any given Eac 0.

We can substitute in that formula and it does match with the polarization resistance that you would get by solving the ODE and calculating the impedance what we had done before for the nonlinear impedance.

(Refer Slide Time: 39:30)

NLEIS- Two step reaction. $R_{p,NL}$

Analytical Calculations: Example

$$\frac{d^2\theta}{dE^2} = \frac{n!}{(k_{ox} + k_r)} \left[\frac{(b_1^n)k_{ox}}{n!} - \sum_{n=2}^{n-1} \left[\frac{(b_1^{(n-1)}k_{ox})}{(m+1)!} \frac{d^{(n-1)}\theta}{dE^{(n-1)}} \frac{1}{(n-m-1)!} \right] \right] \text{ for } n > 0$$

$$\frac{di}{dE} = Fk_{ox} \left[b_1(1-\theta) - \frac{d\theta}{dE} \right] = Fk_{ox} \left[b_1 - \frac{d\theta}{dE} - b_1\theta \right]$$

$$\tau \frac{d\theta}{dt} = k_1(1-\theta) - k_2\theta$$

$$\frac{d^2i}{dE^2} = Fk_{ox} \left[b_1 - \left(\frac{d^2\theta}{dE^2} + 2b_1 \frac{d\theta}{dE} + b_1^2\theta \right) \right]$$

$$\frac{d^2i}{dE^2} = Fk_{ox} \left[b_1^2 - \sum_{n=0}^2 \left(\frac{2}{m} \right) (b_1)^n \left(\frac{d^{(2-n)}\theta}{dE^{(2-n)}} \right) \right]$$

$$\left(\frac{n}{m} \right) = \frac{n!}{m!(n-m)!}$$

$$\frac{d^2i}{dE^2} = Fk_{ox} \left[b_1^2 - \sum_{n=0}^2 \left(\frac{n}{m} \right) (b_1)^n \left(\frac{d^{(2-n)}\theta}{dE^{(2-n)}} \right) \right]$$

$$(R_{p,NL})^{-1} = \frac{2}{E_{ox}} \left[I_1(E_{ox}, D_F) \right] (f(E))$$

$M \xrightarrow{k_1} M^{+}_{ads} + e$
 $M^{+}_{ads} \xrightarrow{k_2} M^{+}_{sol}$

$i_f = Fk_1(1-\theta)$

$\tau \frac{d\theta}{dt} = k_1(1-\theta) - k_2\theta$

So the point is it lengthy but its doable and once you write a program for it its actually pretty fast so we do not need to solve the ODE. You do not need to write a program for calculating these values.

(Refer Slide Time: 39:45)

NLEIS- Summary. $R_{p,NL}$

- In case of simple electron transfer reaction limited by kinetics, $R_p = R_{p,NL}$
- For other reactions, as long as reaction is kinetic limited, and Langmuir isotherm model is used, it is possible to obtain analytical expression for
 - Derivatives of fractional surface coverage with potential
 - Derivatives of current with potential
- Estimate $R_{p,NL}$
- Otherwise, numerical differentiation of i_{dc} vs. E_{dc} can be used
 - Use with care!

$$(R_{p,NL})^{-1} = \frac{2}{RT} \left[\frac{d}{dE} \left[\frac{1}{1 + \exp(-\alpha F(E - E_0)/RT)} \right] \right] f(E)$$



So in summary in case of simple electron transfer reaction charge transfer resistance and polarization resistance are the same as long as its limited the kinetics assuming that mass transfer is rapid and solution resistance is not high. For other reactions as long as we can say that the kinetics is limited and Langmuir isotherm model is applicable we can estimate the polarization resistance.

Using analytical expression basically we can calculate the change in fractional surface coverage with respect to potential the derivatives and we can calculate the nth derivative of current with respect to potential and then using this expression or the equivalent of this expression we can calculate the polarization resistance and in case solution resistance is important we have to write this as $R_{p,NL} + R_{solution}$ inverse of this okay.

Now what happens if mass transfer is not rapid even in that case as long as we can get the nth derivative of current with respect to potential perhaps using numerical method we can still calculate the polarization resistance we can still calculate the nonlinear polarization resistance it

is just that when you take numerical derivatives you have to be very careful. We will stop here today.