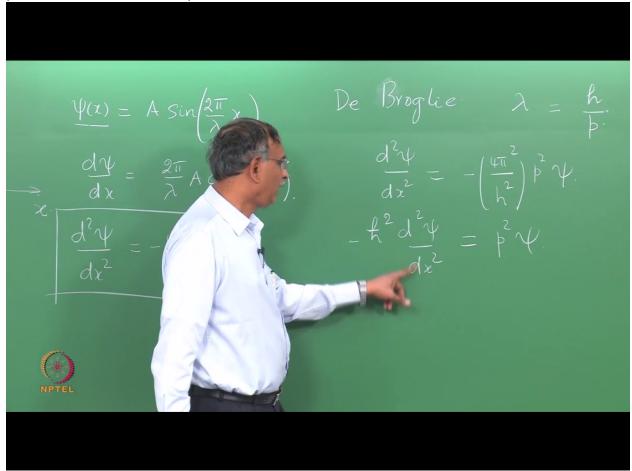
Welcome back to the lecture for the introductory chemistry using Schrodinger and quantum mechanical methods for the atomic structure. So what we would do in this and in the next segment is introduce the Schrodinger equation and also do a model problem using the particle in a one dimensional box model, this is one of the simplest models that we have. Let's take a guick look at the Schrodinger equation. In the lecture earlier I mentioned that I would be talking about the time independent Schrodinger equation in which this quantity was referred to as the Hamiltonian and this as a constant, but with dimensions of energy. And the function Psi is the function that we wanted to find out by solving an equation of this sort, but we don't know what this is right now, we have to introduce that to understand how this equation comes about or what is its origin. We can do a very simple example of a standing wave. And you know that a standing wave is something that happens between fixed points and the wave motion of a particle fixed to the end, something of that kind, and let me put it precisely so that the wave when it reflects it still follows and therefore the standing wave remains as a wave and the amplitudes don't cancel each other. So if you... if you want to look at the axis this is the coordinate or the x-axis that you might want to talk about and this is the axis for the amplitude of the wave at any position X between some fixed points. Obviously for this wave the length of the repeating unit is obviously called the wavelength lambda and here we have 1 2 yes 2 this is 1 and this is 2 and then you have 3 and 3 and a $\frac{1}{2}$, it has to be either exactly $\frac{1}{2}$ wavelength or a full wavelength for this to be a standing wave, okay. Now the equation for the standing wave for the amplitude a, or let us call that amplitude as Psi, in relation to what we have here, we will see later that this Psi is not necessarily the same as the Psi that we talked about, but for that Psi, if we have the maximum amplitude as A, this quantity is A, then the wave function Psi of X is written as A sine 2 Pi by lambda of X, this is something that you are familiar with for a standing wave. Now this quantity Psi, when you differentiate twice it satisfies the derivative equation, let's do that for the first derivative D Psi by DX as 2 Pi by Lambda times A sin cos 2 Pi by lambda X and the second derivative D square Psi by DX square = minus 4 Pi square by lambda square... square Psi of X because this will become sin 2 Pi by lambda of X and that's the same thing as Psi of X, therefore you see that the standing wave satisfies the differential equation D square Psi by DX square where Psi is the amplitude of the wave with lambda, the wave length associated with that. Now the Broglie if you remember in the lecture earlier, gave an expression for the matter waves lambda in terms of the momentum of the particle, in terms of momentum of the particle you have here and therefore if we write the wave equation it's Dr. square Psi by DX square, which is equal to minus 4 Pi square by H square multiplied by P square Psi or minus H bar square, we know that H by 2 Pi is H-bar, therefore if we bring that in, it is minus H bar square D square Psi by DX square = P squared Psi. This is the equation for the standing wave using the De Broglie idea and the quantization idea, namely that the energy quantum for material particles light etcetera given in terms of the Planck's constant, so the Planck's constant enters naturally here in describing what happens to the momentum square on the wave function is the same thing as the sec derivative on the wave function multiplied by minus H bar square.

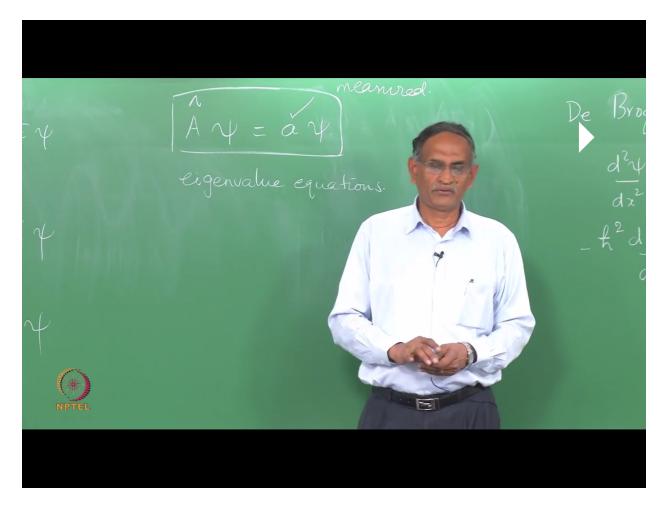
(Refer Slide Time: 6:00)



Therefore if we write the kinetic energy P square by 2 M Psi that turns out to be minus H bar square by 2M D square by DX square Psi. This being the kinetic energy, this is the difference between, if there is a potential energy V then it's a difference between the total energy E and the potential energy V, which may be a function of X for whatever, if there is a potential, we have to consider that, therefore what happens is P square by 2 M is nothing but E minus V on Psi, giving you minus H bar square by 2 M D square by DX square Psi. Now one last step and then you see the equation H Psi is equal to E Psi making sense to us, because now if you bring the V here, just rewrite the equation you have minus H bar square by 2 M D square Psi by DX square plus V of Psi is equal to E of Psi, please remember we had already written this as the kinetic energy and this is on Psi. This is the potential energy on Psi and therefore you see that this is nothing but kinetic energy plus potential energy on Psi, giving you a constant times E Psi and so you see that this is nothing but the Hamiltonian 1 Psi giving you E Psi. This is a very simple

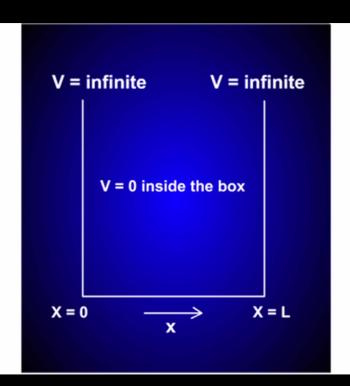
justification. I don't think we can really say that we have derived it from any fundamental principles or whatever, it's a justification to see from a simple standing wave picture and using the de Broglie principle or the proposition with the Planck's constant, it looks like the particle wave function satisfies the equation Hamiltonian. But the Hamiltonian looks somewhat hard. It has a derivative instead of the P square by 2M that we have, now we have put a derivative here and therefore the Hamiltonian is a derivative acting on the wave function and the potential, which is of course a function of the position of whatever particle or the system that you talk about. The potential generally multiplies the wave function but the two together is actually an operator acting on Psi. The Hamiltonian operator acting on Psi giving you a constant time Psi. Schrodinger equation is a very specific equation for the Hamiltonian operator and such equations in mathematics are known as Eigenvalue equations for whatever quantities that appear here. instead of H, it's any other operator that we are going to look at A Psi any operator giving some constant times Psi. Please remember this constant has to have the same dimension as the operator A here in the same way that this constant has the energy dimension for the Hamiltonian operator which is also energy. Any such equation in which A can be measured experimentally such equations are called Eigen value equations... Eigen value equations and the Schrodinger equation, the time-independent Schrodinger equation is the Eigen value equation for the Hamiltonian or the energy operator. This is the picture that you have to. So let me give you some small problems associated with whatever we have done right after this, but then we will go to the next part namely how do we solve this for the specific case of a simple model.

(Refer Slide Time: 10:20)



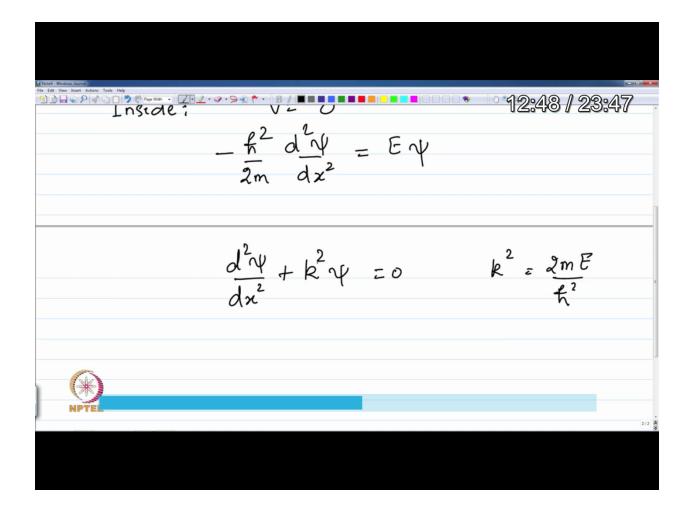
Now what's the model? let's look at the model now of the particle in a one dimensional box. I have a small drawing here that tells you that we have a particle in a finite region, the potentials are in finite at two points namely points with X equal to 0 and the point X is equal to L meaning that the particle is confined to a region of a box of length L and the particle motion or the particle coordinate is only one coordinate or one variable namely X. Let's assume for the time being that the potential inside the box is 0, so this is what we call as the particle in a one dimensional box with infinite barriers and what does this particle give you.

(Refer Slide Time: 11:26)



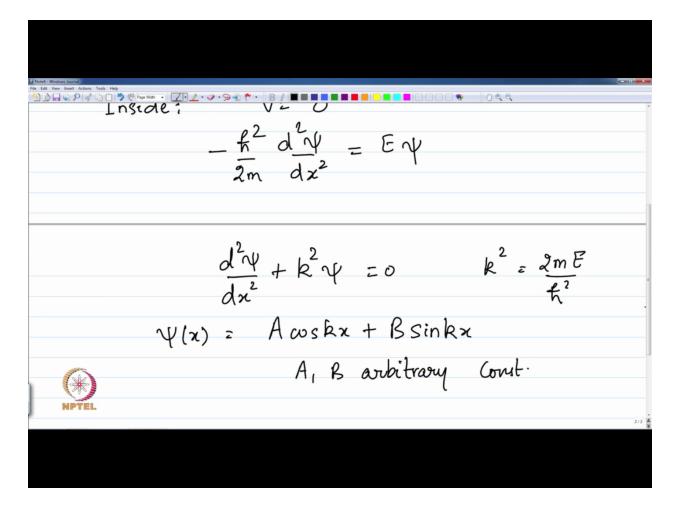
Now let's look at the equations, we have minus H bar square by 2M D square Psi by D X square plus V of Psi is equal to P of Psi. if the potential is infinite then Psi has to be 0 in order to satisfy that therefore at the boundaries X is equal to 0, X is equal to L, the wave function Psi FX is 0. Inside the box we have V is 0, therefore what we have is minus H bar square by 2M D square Psi by DX square is equal to E Psi. The total energy, because there is no potential inside the box.

(Refer Slide Time: 12:52)



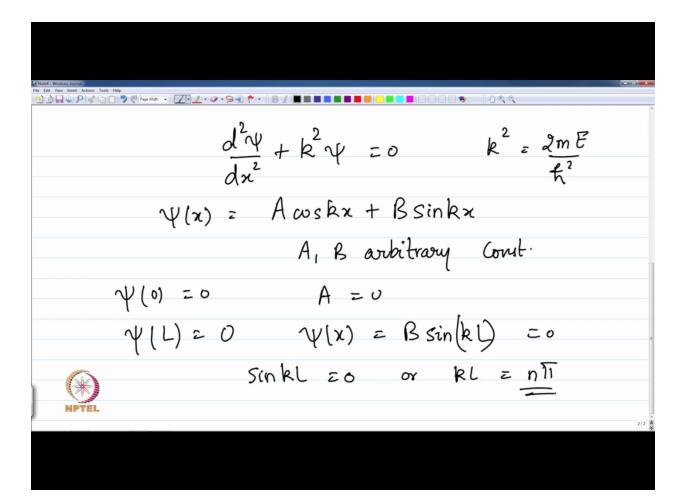
We shall solve this in a very quick manner namely D square Psi by DX square plus a constant... a positive constant K squared Psi is equal to 0 where K square is 2 M E by H bar square. This is... the K square is positive obviously and therefore what you have here is a simple derivative equation for second order and you know such functions can be obtained, the solutions can be obtained from either trigonometric function or the exponential with imaginary argument. Let's use the trigonometric function namely a sin, let's write that to be consistent we have a cos KX plus B sin KX, where A and B are arbitrary constants... arbitrary constants.

(Refer Slide Time: 13:42)

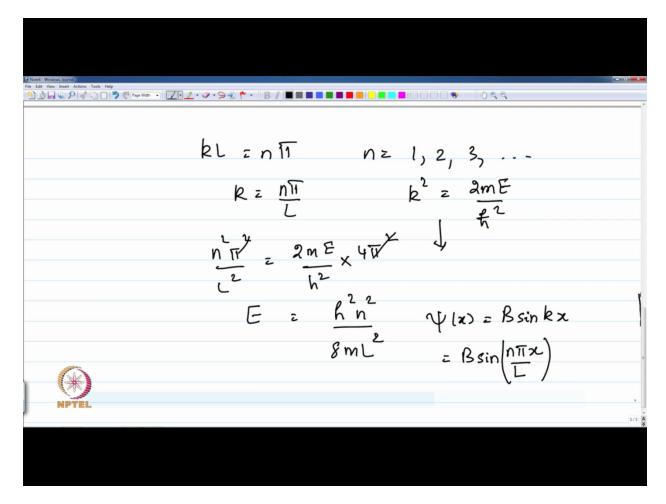


Now if you look at that solution with the boundary condition that you have namely Psi of 0 is 0, immediately you have A is equal to 0, because cos KX is 1 and sin KX goes to 0, therefore A is equal to 0. If you have Psi at Left, which is the other extreme of the box, please remember this model at X is equal to L at this point okay. Therefore we have Psi of L is 0, which implies that since A is already 0, Psi of X is B sine K L and that's equal to zero. We don't want B to be 0, because if A and B are 0, that's anyway it's a trivial solution for any such differential equation, doesn't give you anything of interest, I think there's no meaning, there's no interpretation. Therefore we are going to consider the case obviously a non trivial solution, with B not equal to 0, which means sin KL has to be 0 or KL has to be an integer times Pi, N is an integer.

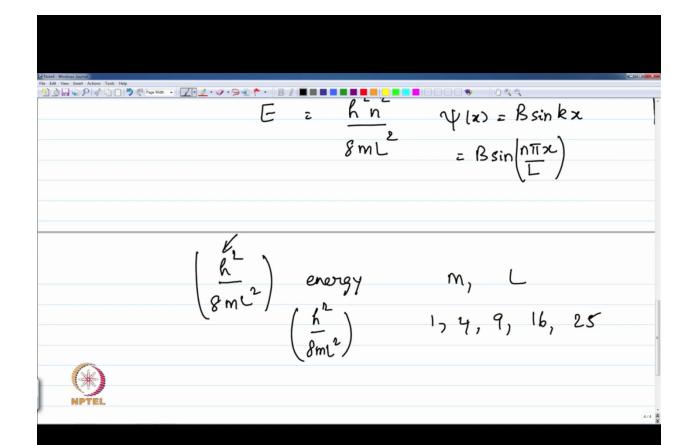
(Refer Slide Time: 15:00)

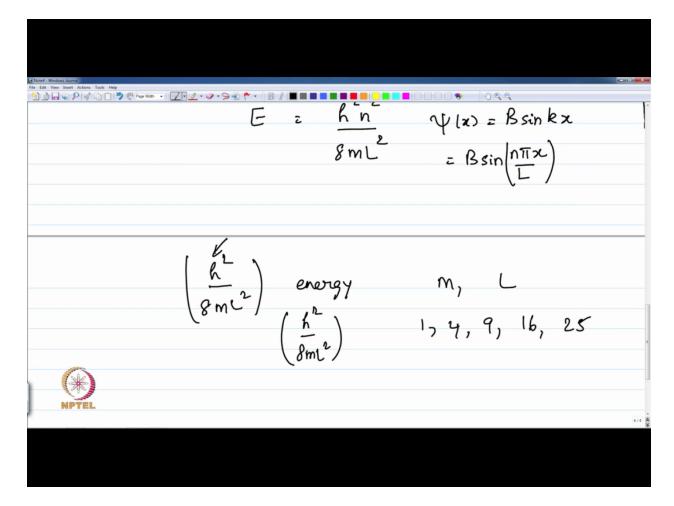


KL is equal to N Pi and N has to be, obviously we don't want N equal to 0, which is also the case of triviality, and so what we have is N is equal to 1 2 3 etcetera integers or please remember K is equal to N Pi by L, look at this K square if you recall is 2ME by H bar square, therefore this gives you immediately that M square Pi square by L square is equal to 2M E by H square times the 4 Pi square that we have, cancel things off and you immediately get the solution namely E is equal to H square M square by 8 M L square and what is the solution for the wave function Psi of X is B sin KX, which is B sin N Pi X by L. Because K is N Pi by L okay. So this is the simplest solution, but two important results. One is that the energy for the particle in the box which is subject to boundary conditions that the wave function vanishes at some boundaries, subject to that the particle energy appears to be quantized, is not arbitrary. You recall the dimension, the quantity H square by M square, H square by 8 ML square the quantity. (Refer Slide Time: 16:45)

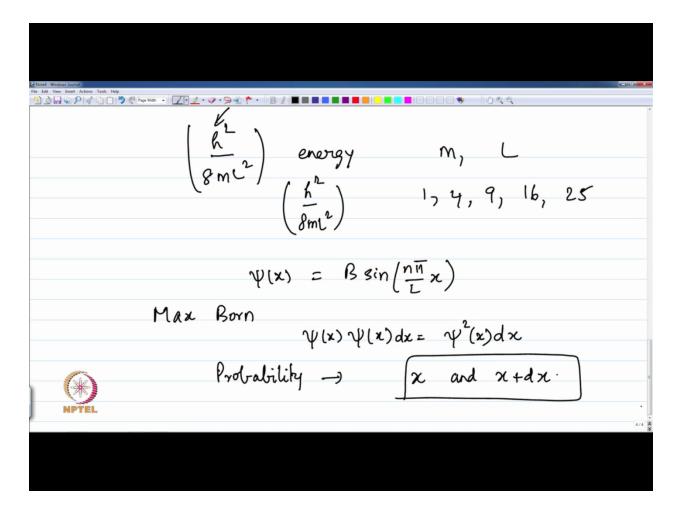


Has the dimension of the energy and it has the only two inputs, which is... which are the inputs for this problem, namely the mass of the particle M and the length of the box Left, and the other constant is of course Planck's constant. So now the energy seems to be quantized in terms of the, the two physical parameters that we introduced, which particle, a larger particle, a heavier particle or a lighter particle, in a smaller box or in the larger box, but with all the other conditions being the same namely potentials being zero inside the potentials being in finite, given that you see that the energy is discretized and the energy is in the units of H square by 8 ML square this is the fundamental unit for this box and then it is 1 4 9 16 25, as the value of N becomes 1 2 3 4 etc. Therefore particle particle energies are discretized. (Refer Slide Time: 18:00)





The second part is the other, namely the wave function is given in terms of B sine N Pi X by L. Now what is this wave function, from the beginning of this lecture you might think that this wave function is essentially a function telling you how the particle is oscillating, that's not true okay. That picture was a starting point for us to get an idea that the Schrodinger equation is like this, the wave function that we have here is not a function representing how the particle is moving, it is just a function associated with that particle. What's the meaning of it, Max Bourne gave the interpretation namely that wave function by itself does not have any meaning, but Psi of X-square, Psi Star Psi, in this case Psi is real, therefore Psi of X Psi of X are Psi squared of X. In a small interval DX gives the probability of the particle being in the position between X and X plus DX. The probability of locating the particle between X and X plus DX, that's the number given by the product of the wave function with itself in this case, because it's real that Max Bourne suggested that Psi squared X DX gives the probability that the system we found in the interval X and X plus DX, that's all there is to it. (Refer Slide Time: 19:58)



Therefore let me conclude immediately what D should be, because if Psi Star X Psi X, which is the same as Psi of X square, with a DX is a probability, then if you add all the probabilities from 0 to L, because the particle can have any position between the endpoint, but not at the endpoint from anywhere as close to the endpoint as possible, but as close to the other endpoint. Therefore if you integrate the total probabilities, this being a continuous function you have 0 to L Psi X square DX, that probability has to add to 1, because we have made sure that the potentials are infinite in our model, therefore the partner cannot we found outside of that region. Therefore the probability that the particle stays inside the box is 1. This gives you immediately a value for B, because you have B square sin square N Pi X by L, DX between 0 and L, that's equal to 1, which gives you the value B = root 2 by L, okay.

(Refer Slide Time: 21:15)

The fit two last states that how

$$\psi^{4}(x) \psi(x) dx = 1 \\
\psi(x)^{2} dx = 1$$

$$\psi(x)^{2} dx = 1$$

Therefore you have got two results for the particle in the box namely the wave function is root 2 by L sin N Pi X by L and E, the particles energy is given by H square N square by 8 M L square. Now because the energy is given by the quantum number N, let me use a highlighter here, because it's given by N and N can take any number of values and for that N, the corresponding wave function is sin N Pi X by L, we see that there are many solutions to the wave function and many solutions to the energy, this will also turn out to be a general property when we solve the Hamiltonian equation... the Schrodinger equation for the systems in all the other models, that in one step, you will get all the different types of... all the possible energies and all the possible wave functions and the best way to... I mean a convenient way, I wouldn't call it the best way, a convenient way is to label the wave function with the quantum number Psi N of X and E N for a given quantum number N. So let me summarize and then stop for this lecture namely the particle in the 1D box has two results, A quantization of energy or discretization due to boundary conditions and of energy E and a probability statement for determining the position of the particle in the box at various locations, okay. Let's continue this in the next part and complete the remaining that we needed to do in terms of what are called the measurables

and then how do we interpret this probability and so on for various values,

we will do that in the second part, until then thank you.

