

Quantum Mechanics and Molecular Spectroscopy
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Lecture – 15
Absorption Probability

Welcome to lecture number 15 of the course quantum mechanics and molecular spectroscopy. As usual, we will have a quick recap of the previous lecture and continue with the present lecture.

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$$P_f(t) = \frac{E_0^2}{\hbar^2 \omega^2} \left| \int_0^t dt e^{-i\omega_f t} \cos(\omega t) \langle f | \hat{\mu} | i \rangle \right|^2$$

Transition moment integral.

$$P_f(t) \propto \left| \langle f | \hat{\mu} | i \rangle \right|^2$$

operator $\hat{\mu} \Rightarrow$ electrical must be parallel to the dipole moment

$$\int_0^t dt e^{-i\omega_f t} \cos(\omega t) \quad \text{with } \cos(\omega t) = \frac{1}{2} \left[e^{i\omega t} + e^{-i\omega t} \right]$$

$$\frac{1}{2} \int_0^t dt \left[e^{i\omega_f t} \left[e^{i\omega t} + e^{-i\omega t} \right] \right]$$

$$= \frac{1}{2} \int_0^t dt \left\{ e^{i(\omega_f + \omega)t} + e^{i(\omega_f - \omega)t} \right\}$$

$\omega_f = -\omega_{fi}$
 $E_i - E_f \quad E_f - E_i$

In the last lecture, we looked at the transition probability to state f and this was given by E_0^2 square by \hbar^2 square ω_f^2 by ω^2 square modulus of integral 0 to t prime dt e to the power of $-i\omega_f t$ Cos ωt f dot mu i square, yeah so this is the integral and this I told you is the transition moment integral. So, one can think of P of t to be proportional to the square of the transition moment integral okay.

And in the transition moment integral, the operator says, E dot mu says that the dipole moment of the molecule or the atom okay, by the way dipole moment is not same as the permanent dipole moment okay μ_0 , dipole moment of the molecule and epsilon, this is epsilon is nothing but your electric field. So electric field must be parallel to the dipole moment or at least should have some projection, it cannot be perpendicular okay.

Now, I take the other, so apart from that there is this integral in this okay. That integral is 0 to

t prime dt e to the power of $-i$ ω if t $\cos \omega t$. I am going to manipulate this integral little bit. Now we know that $\cos \omega t = \frac{1}{2} (e^{i\omega t} + e^{-i\omega t})$. So if I substitute that in here, then what I will get is 0 to t prime dt , I will write two, so I will take the half outside.

This is nothing but e to the power of $i\omega t$, now ω if $= -\omega$ if. This is nothing but $E_i - E_f$ and this is nothing but $E_f - E_i$ so that is just a reverse of sign, ω if and ω if. So this can be written as $\frac{1}{2} (e^{i\omega t} + e^{-i\omega t})$ okay. So this will be nothing but half of 0 to t prime dt , then you will have 2 into e to the power of $i\omega t$ if $+ \omega$ into t $+ e$ to the power of $-i\omega t$ if $- \omega$ into t okay, so that is the integral.

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$$P_f(i) = \frac{E_0^2}{4\pi^2} \left(\frac{\omega_H}{\omega} \right)^2 \left| \left[\int_0^{t'} dt e^{i(\omega_H + \omega)t} + \int_0^{t'} dt e^{i(\omega_H - \omega)t} \right] \langle \hat{\epsilon} \cdot \hat{i} \rangle \right|^2$$

$(x,y) \rightarrow$ variables
 x and y are conjugate variables.
 x -length $y = \frac{1}{\text{length}}$ $(x,y) \rightarrow$ dimensionless quantity.

So expanding or continuing what we have P of f of $t = \frac{E_0^2}{4\pi^2} \left(\frac{\omega_H}{\omega} \right)^2 \left| \int_0^{t'} dt e^{i(\omega_H + \omega)t} + \int_0^{t'} dt e^{i(\omega_H - \omega)t} \right|^2$. Now the question is what does it say? It says that the ω or the electromagnetic field acts from time 0 to time t prime okay.

So think of it like this, so there is some perturbation, time-dependent perturbation that starts at 0 okay and ends at t prime, so 0 to t prime. So this is your limit of integration, so this will be nothing but your limit of integration okay, but if you consider the entire time of minus infinity to plus infinity okay, then what happens is from minus infinity to 0 , there is no light, so there is no perturbation okay and similarly after t prime there is no perturbation.

So the effect of perturbation before $t = 0$ and after $t = t$ prime is going to be non-existent. If that is non-existent, then without losing any physical concept, this integral can now be written as minus infinity to plus infinity $dt e$ to the power of $i \omega f_i + \omega t$ and this integral can be written as integral dt minus infinity to plus infinity e to the power of $i \omega f_i - \omega t$. Now, this is a case of adding zeros.

For example you want to know how much money you have in your bank within say some period of time okay, say first of a month to 15th of a month, but before that you neither you open the account on say first of the month and you have closed the account on 15th of the month and you want to know what is, but if I want to look at the entire time period before the first of the month, the previous month, and after 15th of that month okay you have the account is opened and closed.

So before the opening previous month and after 15th of month the money that account will have will have zeros, will add to zeros. So essentially, the entire transaction will try only between the first of the month to 15th of the month and transactions before that and transactions after that will not lead to any usefulness, so they are all zeros. So this is the similar scenario where you know the limit of integration is only between 0 and t prime.

But any extending the integration limits to minus infinity to plus infinity is like adding zeros and has no physical consequence except the fact that it turns out to be a standard integral okay. Now, one can write a standard integral between 2 variables let us say x and y okay such that x and y are conjugate. Now, what are conjugate variables? Variables that have inverse unit, so for example if x is length, then y will be 1 over length okay.

Such that product of x and y will give to dimensionless quantity okay, such variables are called conjugate variables. In quantum mechanics, x and momentum or position and momentum are conjugate variable, energy and time are conjugate variable and they are also related by the Schrodinger equation

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Fourier Integral of two conjugate variables. (x,y)

$$\delta(x-a) = \frac{1}{2\pi} \int_{-\infty}^{\infty} dy e^{i(x-a)y}$$

Kronecker delta

$$x=a \Rightarrow 1$$

$$\text{else} \Rightarrow 0$$

$$\int_{-\infty}^{\infty} dt e^{i(\omega_f + \omega)t} + \int_{-\infty}^{\infty} dt e^{i(\omega_f - \omega)t}$$

Now, one can write a Fourier integral of 2 conjugate variables, okay let us say x, y in such a way that $x-a = 1$ over 2π integral $-\infty$ to $+\infty$ $dy e$ to the power of i to $x-a$ y okay. Now, what is d ? This is nothing but your Kronecker delta participant that means when $x = a$ this will go to 1, else will go to 0 okay. Now let us look at the integrals that we had. What are the 2 integrals that we had in the previous case?

Those were minus infinity to plus infinity $dt e$ to the power of i $\omega_f + \omega$ $t + -\infty$ to $+\infty$ $dt e$ to the power of i $\omega_f - \omega$ t . Now you can look at these 2 integrals okay, you can compare this integral with these 2 integrals. They look very similar except the fact that the 2π is missing, but 2π can always be multiplied and divided.

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$$P_f(t) = \frac{4E_0^2 \pi^2}{\hbar^2} \left(\frac{\omega_f}{\omega}\right)^2 \left| \frac{1}{2\pi} \int_{-\infty}^{\infty} dt e^{i(\omega_f + \omega)t} + \frac{1}{2\pi} \int_{-\infty}^{\infty} dt e^{i(\omega_f - \omega)t} \right|^2 \langle f | \hat{\mu} | i \rangle^2$$

$\delta(\omega_f + \omega)$ $\delta(\omega_f - \omega)$

$$P_f(t) = \frac{4E_0^2 \pi^2}{\hbar^2} \left(\frac{\omega_f}{\omega}\right)^2 \left| \delta(\omega_f + \omega) + \delta(\omega_f - \omega) \right|^2 \langle f | \hat{\mu} | i \rangle^2$$

$\delta(\omega_f + \omega)$
 $\delta(\omega_f - \omega)$

E_f ————— $\hbar\omega_f$
 \uparrow
 $\Delta E = \hbar(\omega_f - \omega_i) = \hbar\omega_f - \hbar\omega_i$
 \downarrow
 E_i ————— $\hbar\omega_i$

So if I do that, then Pf of t will become E_0^2 square okay, so what I will do is I will multiply by

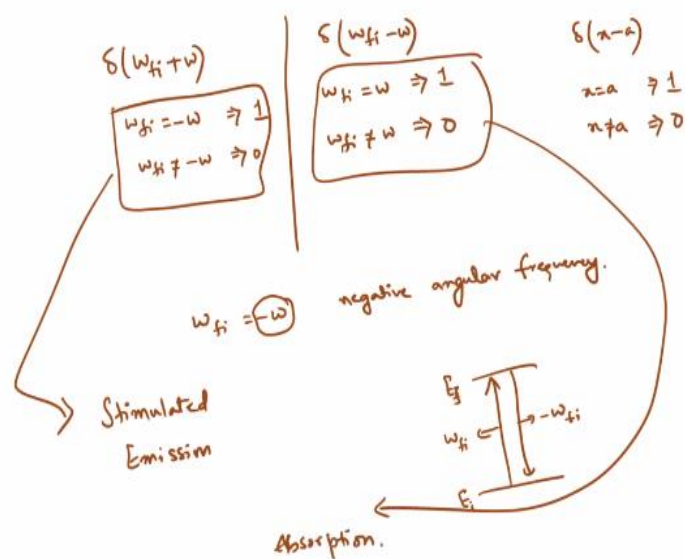
$\pi^2 \hbar^2 \omega_{fi}^2 / \omega^2 \left| \int_{-\infty}^{+\infty} dt e^{i(\omega_{fi} + \omega)t} \right|^2$ okay, sorry I will multiply it by $4\pi^2$ okay, $1 / 2\pi \int_{-\infty}^{+\infty} dt e^{i(\omega_{fi} + \omega)t}$ + $1 / 2\pi \int_{-\infty}^{+\infty} dt e^{i(\omega_{fi} - \omega)t}$ square $f \mu \cdot \epsilon_0 \hbar^2$ okay.

Now if you have this, now look at this, this function can be written as $\delta(\omega_{fi} + \omega)$ and this function can be written as $\delta(\omega_{fi} - \omega)$. So your P of t will now become $4 E_0^2 \pi^2 \hbar^2 \omega_{fi}^2 / \omega^2 \left| \delta(\omega_{fi} + \omega) + \delta(\omega_{fi} - \omega) \right|^2 f \epsilon_0 \mu \hbar^2$ okay. So this is your probability of transition from state i to state f .

Now this involves 2 factors that are the Kronecker deltas, $\omega_{fi} + \omega$ and $\omega_{fi} - \omega$. Now, let us consider a very simple scenario. The scenario is this okay, this is your E_i that will be nothing but $\hbar \omega_i$ and this is E_f this is the $\hbar \omega_f$ and this energy difference is ΔE , this is equal to $\hbar \omega_f - \hbar \omega_i$, this is equal to $\hbar \omega_{fi}$ that is the energy separation.

So ω_{fi} is nothing but the frequency of this energy separation okay. Now what is this, so which means in this case if you look at these, if you have $\delta(\omega_{fi} + \omega)$ and you have $\delta(\omega_{fi} - \omega)$.

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$\delta(\omega_{fi} + \omega)$, now remember we showed when you have $\delta(x-a)$ okay. So this implies if $x = a$, then this function will go to 1 and x is not equal to a , this function will go to

0. If I apply the same principle $\omega_f = \omega$, this means when $\omega_f = -\omega$, this function will go to 1. When ω_f is not equal to $-\omega$, this function will go to 0. Similarly if you have $\omega_f = \omega$, then $\omega_f = \omega$, this function will go to 1.

And ω_f is not equal to ω this function will go to 0. Now there is one issue is that what is this condition? ω is angular frequency, how can angular frequency be negative? Okay, so here I am looking at a condition $\omega_f = -\omega$ okay, which means this has to be negative and negative angular frequency does not exist okay. One way to look at the negative angular frequency is emission of light when positive.

So which means when you go from top to bottom that is E_i to E_f one can think of the frequency to be positive and when you go down the frequency, so this will correspond to ω_f and this will correspond to $-\omega_f$ okay. So when the frequency negative simply means it is a case of stimulated emission. So this corresponds to stimulated and this will correspond to absorption okay.

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$$P_f(t) = \frac{4\pi^2 E_0^2}{\hbar^2} \left(\frac{\omega_f}{\omega}\right)^2 \left| \frac{\delta(\omega_f + \omega)}{\delta(\omega_f - \omega)} \right|^2 |\langle f | \hat{\epsilon} \cdot \mu | i \rangle|^2$$

↑ Stimulated Emission ↑ Absorption.

$\omega_f = -\omega \quad \omega_f = \omega$

not simultaneously possible.

For absorption from initial state $|i\rangle$ to a final state $|f\rangle$

$$P_f(t) = \frac{4\pi^2 E_0^2}{\hbar^2} \left(\frac{\omega_f}{\omega}\right)^2 \left| \delta(\omega_f - \omega) \right|^2 |\langle f | \hat{\epsilon} \cdot \mu | i \rangle|^2$$

↑ Kramers delta. ↑ T_{fi}

Now, let us go back to our probability that is nothing but P_f of $t = 4\pi^2 E_0^2$ square by \hbar^2 square ω_f by ω square modulus of, now there are 2 terms. So there are 2 terms, let me write down $\omega_f = \omega + \delta$ $\omega_f = \omega - \delta$ square integral $f \epsilon \cdot \mu i$ whole square okay. Now you look at there are 2 terms, this one and this. So I told you this is nothing but stimulated emission and this is nothing but absorption.

Now it turns out that it is not possible for ω_{fi} to be positive of ω and ω_{fi} to be negative of ω simultaneously okay. So ω_{fi} can be equal to $-\omega$ and $\omega_{fi} = \omega$ is not simultaneous. That means stimulated emission and absorption cannot happen simultaneously, they have to happen one after the other okay. So in that case since they cannot happen simultaneously, you cannot have both the conditions.

You can have only one condition and if you are dealing with absorption, then this condition will happen and when you are dealing with stimulated emission this condition will be valid. So for absorption from initial state i to a final state f , the probability of this transition is equal to $4\pi^2 \epsilon_0^2 \hbar^2 \omega_{fi}^2 / \omega^2 \delta(\omega_{fi} - \omega)$ okay.

So, this is the probability of transition for the absorption from an initial state i to a final state f okay. Now, there are several factors here, one is the transient moment integral TMI and then there is this Kronecker delta and then there is a ratio of ω_{fi} to ω okay and many other factors are dependent, but this is the probability of transition and then we will look at this particular equation in the next lecture more carefully. I am going to stop it here and thank you very much.