




**Mass transfer phenomenon in polymers: Laminar flow and boundary layer conditions**

Hello friends, welcome to the next episode of Mass Transfer Operations. In this particular segment, we are going to discuss the laminar flow and boundary conditions in mass transfer operations. So, let us have a brief outlook on what we discussed in the previous lecture, we discussed the mass transfer coefficient in the previous lecture, then different types of mass transfer coefficient, and then we discussed the equimolar counter diffusion in two components. We established the relationship between the mass transfer coefficient and we discussed the mass transfer coefficient and film thickness. Apart from this, we discussed the dimensionless numbers in mass transfer operations. In this particular segment, we are going to discuss the mass transfer coefficient in laminar flow.

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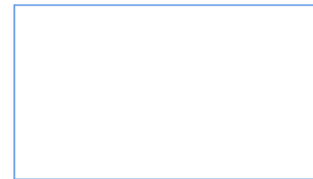


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We are going to discuss the mass transfer in falling film. Apart from this, we will discuss the laminar falling film in the inclined surface. We will also discuss the mass transfer coefficient in turbulent flow. Apart from this, we will discuss boundary clear theory and film theory, all these things we are going to discuss in this particular section.

## Mass transfer coefficient in laminar flow

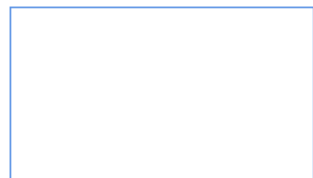
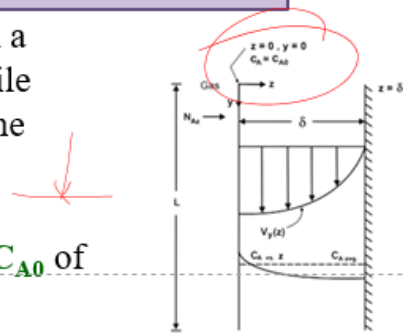
- **In principle** we do not need to study **the mass transfer coefficient in laminar flow conditions**.
- **A uniform method** of dealing both **laminar and turbulent flow** is nevertheless desirable.
- We shall **choose one relatively simple situation** to illustrate **the general technique** and to provide some basis for considering **turbulent flow**.



Now, let us talk about the mass transfer coefficient in laminar flow. In principle, rather we did not, we do not need to study the mass transfer coefficient in laminar flow conditions. A uniform method dealing this both laminar flow and turbulent flow is nevertheless desirable. So, we shall choose one relatively simple situation to illustrate the general technique and to provide some basis for considering turbulent flow. Now, this is the particular figure.

## Mass transfer coefficient in laminar flow

- Figure shows **a liquid falling in a thin film** in a laminar flow **down a vertical flat surface** while being exposed to a gas A, which dissolves in the liquid.
- The liquid contains a **uniform concentration  $C_{A0}$**  of A at the top.
- At the liquid surface **the concentration of the dissolved gas is  $C_{Ai}$** , in equilibrium with the pressure of A in the gas phase, since  **$C_{Ai} > C_{A0}$**  gas dissolves in the liquid.





This particular figure shows a liquid falling in a thin film in a laminar flow and down a vertical flat surface while being exposed to gas A, in that case, which dissolved in that particular liquid. So, this is a very common phenomenon when a gas is usually dissolved in a particular liquid like in a scrubbing operation. So, this liquid contains a uniform concentration of  $C_A$ , not of A at the top. At the liquid surface, the concentration of the dissolved gas is  $C_{A1}$ .

In the equilibrium with the pressure A in the gas phase, now since  $C_{A1}$  is greater than  $C_{A0}$ , gas dissolves in the liquid. The problem is to obtain the mass transfer coefficient  $k_l$  with which the amount of gas dissolved after the liquid falls a certain distance  $l$ , this can be computed. The problem is solved by the simultaneous solution of the equation of continuity for component A with the equation which usually describes the liquid motion in the Navier-Stokes equation. The simultaneous solution of a formidable set of partial differential equations becomes possible only when several simplifying assumptions are made. Now, let us consider this particular equation which is derived for unsteady state mass transfer.

### Mass transfer coefficient in laminar flow

- Consider the following **equations of continuity** derived for **unsteady state mass transfer**:
 

$$V_x \frac{\partial C_A}{\partial x} + V_y \frac{\partial C_A}{\partial y} + V_z \frac{\partial C_A}{\partial z} + \frac{\partial C_A}{\partial t} = D_{AB} \left( \frac{\partial^2 C_A}{\partial x^2} + \frac{\partial^2 C_A}{\partial y^2} + \frac{\partial^2 C_A}{\partial z^2} \right) + R_A \quad \dots (1)$$
- For the present purpose, **assume that**
  - a) there is **no chemical reactions** in the systems,  **$R_A = 0$**
  - b) Conditions do not change in the x direction (perpendicular to the plane of the paper. All derivatives with respect to **x should be 0**.
  - c) **Steady state condition** prevail,  **$\frac{\partial C_A}{\partial t} = 0$**



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#### Continuity equation

$$V_x \frac{\partial C_A}{\partial x} + V_y \frac{\partial C_A}{\partial y} + V_z \frac{\partial C_A}{\partial z} + \frac{\partial C_A}{\partial t} = D_{AB} \left( \frac{\partial^2 C_A}{\partial x^2} + \frac{\partial^2 C_A}{\partial y^2} + \frac{\partial^2 C_A}{\partial z^2} \right) + R_A$$

At steady state condition;

$$\frac{\partial C_A}{\partial t} = 0$$

Here,  $v_x \frac{\partial C}{\partial x} + v_y \frac{\partial C}{\partial y} + v_z \frac{\partial C}{\partial z} + \frac{\partial C}{\partial t}$  is equal to  $D_{AB} \left( \frac{\partial^2 C}{\partial x^2} + \frac{\partial^2 C}{\partial y^2} + \frac{\partial^2 C}{\partial z^2} \right) + r_A$ . This is the equation number 1. For the present purpose, assume that there is no chemical reaction within the system. So, we can take  $r_A$  is equal to 0 and condition do not change in the x direction perpendicular to the plane of the paper. All derivatives with respect to x should be 0.

## Mass transfer coefficient in laminar flow

- **Other assumptions are:**

d) **The rate of absorption** of gas is very small this means that  $V_z$  in equation 1 **due to diffusion** of A is essentially **zero**.

e) **Diffusion of A** in the y direction is **negligible** in comparison with the movement of A outward due to bulk flow,

$$D_{AB} \frac{\partial^2 C_A}{\partial y^2} = 0$$

f) **The physical properties** in this case ( $D_{AB}$ ,  $\rho$ ,  $\mu$ ) are **constant**. So, equation 1 reduces to:

$$V_y \frac{\partial C_A}{\partial y} = D_{AB} \frac{\partial^2 C_A}{\partial z^2} \quad \dots (2)$$



$$D_{AB} \frac{\partial^2 C_A}{\partial y^2} = 0$$

$$V_y \frac{\partial C_A}{\partial y} = D_{AB} \frac{\partial^2 C_A}{\partial z^2}$$

Now, if we take the steady state condition previous, in that case,  $\frac{dC_A}{dt}$  is equal to 0. Now, we can take the other assumptions too. One is that the rate of absorption of gas is very small. This means that  $v_z$  in equation 1, this is equation 1, due to the diffusion of A is essentially 0. A diffusion of A in the y direction is negligible in comparison with the movement of A outward due to bulk flow and that is  $D_{AB} \frac{\partial^2 C_A}{\partial y^2}$  is equal to 0 and the physical properties in this case, that is  $D_{AB}$ ,  $\rho$ ,  $\mu$  are constant. So, equation 1 which can this equation can be reduced to  $v$  at  $y$   $\frac{\partial C_A}{\partial y}$  is equal to  $D_{AB} \frac{\partial^2 C_A}{\partial z^2}$  and we can say that this is equation number 2. Now, if we talk about the mass transfer in falling film, now, this is states that A added to the liquid running down at any location  $z$  over an increment of  $\delta$ .

$$\mu \frac{\partial^2 V_y}{\partial z^2} + \rho \cdot g = 0$$

$$V_y = \frac{\rho \cdot g \cdot \delta^2}{2 \cdot \mu} \left[ 1 - \left( \frac{z}{\delta} \right)^2 \right]$$

## Mass transfer in falling film

- This states that A added to the **liquid running down at any location z** over an increment in y, got there by **diffusion in the z direction**.
- The equation of motion under this condition will be again reduces to

$$\mu \frac{\partial^2 V_y}{\partial z^2} + \rho \cdot g = 0 \quad \dots (3)$$

- The solution to equation 3 with the conditions  $V_y = 0$  at  $z = \delta$  and that  $dV_y/dz = 0$  at  $z = 0$ , is well known

$$V_y = \frac{\rho \cdot g \cdot \delta^2}{2 \cdot \mu} \left[ 1 - \left( \frac{z}{\delta} \right)^2 \right] \quad \dots (4)$$



$$V_{y,max} = \frac{\rho \cdot g \cdot \delta^2}{2 \cdot \mu}$$

$$V_{y,avg} = \frac{1}{A} \int V_y dA = \frac{1}{W\delta} \int_0^W \int_0^\delta V_y dx dz = \frac{W}{W\delta} \int_0^\delta V_y dz = \frac{1}{\delta} \int_0^\delta \frac{\rho \cdot g \cdot \delta^2}{2 \cdot \mu} \left[ 1 - \left( \frac{z}{\delta} \right)^2 \right] dz$$

$$V_{y,avg} = \frac{\rho \cdot g \cdot \delta^2}{3 \cdot \mu}$$

So, we have to be the diffuser, we have to be called as the diffusion in the z direction and the equation of motion under this condition can be reduced to  $\mu \frac{\partial^2 v_y}{\partial z^2} + \rho \cdot g = 0$  and that is the equation number 3. Now, the solution of equation 3 with the condition  $v_y$  is equal to 0 at  $z$  is equal to  $\delta$  and that  $dv_y/dz$  is equal to 0 at  $z$  is equal to 0 is well known and this can be represented like this  $v_y$  is equal to  $\frac{\rho \cdot g \cdot \delta^2}{2 \cdot \mu} \left[ 1 - \left( \frac{z}{\delta} \right)^2 \right]$  and that can be represented as equation number 4. So, the maximum velocity that occurs at  $z$  is equal to 0 in this particular equation  $v_{y,max}$  is equal to  $\frac{\rho \cdot g \cdot \delta^2}{2 \cdot \mu}$  and that is equation number 5.

## Mass transfer in falling film

- The **maximum velocity** which occurs at  $z = 0$  in equation 4

$$V_{y,\max} = \frac{\rho \cdot g \cdot \delta^2}{2 \cdot \mu} \quad \dots (5)$$

- The **bulk average velocity** can be obtained as follows:

$$V_{y,\text{avg}} = \frac{1}{A} \iint V_y \, dA = \frac{1}{W\delta} \int_0^W \int_0^\delta V_y \, dx \, dz = \frac{W}{W\delta} \int_0^\delta V_y \, dz = \frac{1}{\delta} \int_0^\delta \frac{\rho \cdot g \cdot \delta^2}{2 \cdot \mu} \left[ 1 - \left( \frac{z}{\delta} \right)^2 \right] dz$$

$$V_{y,\text{avg}} = \frac{\rho \cdot g \cdot \delta^2}{3 \cdot \mu} \quad \dots (6)$$



So, the bulk average velocity can be obtained like  $v_y$  average this can be represented as a mathematical formula like this and  $v_y$  average can be represented if we adopt all kind of integration and another thing can be put as  $\rho \cdot g \cdot \delta^2$  over  $3 \mu$  and this can be written as equation number 6.

$$\delta = \left( \frac{3 \cdot V_{y,\text{avg}} \mu}{\rho \cdot g} \right)^{\frac{1}{2}}$$

$$\frac{3}{2} V_{y,\text{avg}} \left[ 1 - \left( \frac{z}{\delta} \right)^2 \right] \frac{\partial C_A}{\partial y} = D_{AB} \frac{\partial^2 C_A}{\partial z^2}$$

## Mass transfer in falling film

- The **film thickness** is then  $\delta = \left( \frac{3 \cdot V_{y,\text{avg}} \mu}{\rho \cdot g} \right)^{\frac{1}{2}} \quad \dots (7)$

- Substituting equation 4 into equation 2 and then by using equation 6, we get:

$$\frac{3}{2} V_{y,\text{avg}} \left[ 1 - \left( \frac{z}{\delta} \right)^2 \right] \frac{\partial C_A}{\partial y} = D_{AB} \frac{\partial^2 C_A}{\partial z^2} \quad \dots (8)$$

- Which is to be solved under **the following conditions**

- At  $z = 0$ ,  $C_A = C_{A,i}$  at all the values of  $y$ .
- At  $z = \delta$ ,  $\frac{\partial C_A}{\partial z} = 0$  at all the values of  $y$ , since no diffusion takes place into the solid wall.
- At  $y = 0$ ,  $C_A = C_{A,0}$  at all values of  $z$



So, the film thickness can be represented as  $\delta$  is equal to  $3 \sqrt{\frac{\nu}{g}}$  by average  $\mu$  over  $\rho g$  to the power half and that is equation number 7. So, if we substitute equation 4 into equation 2, which we derived previously, and then by equation 6, we get this final equation  $\frac{\partial c}{\partial y} = \frac{D_A b}{\nu} \frac{\partial^2 c}{\partial z^2}$ . So, which is to be solved under the different conditions like at  $z$  is equal to 0,  $c_a$  is equal to  $c_{a,i}$  at all the values of  $y$  and at  $z$  is equal to  $\delta$ ,  $c_a$  over  $\partial z$  is equal to 0 at all the values of  $y$  since no diffusion takes place into the solid wall and at  $y$  is equal to 0,  $c_a$  is equal to  $c_{a,0}$  at all values of  $z$ . So, the solution in the result in the general expression and infinite series gives  $c_a$  for any for any  $z$  and  $y$  therefore, the providing the distribution that  $c_a$  at  $z$  at  $y$  is equal to solid wall  $L$ .

### Mass transfer in falling film

$c_A$  for any  $z$  &  $y$

Eq (4) & (5) into consideration

$v_y = v_{max} \left[ 1 - \left( \frac{z}{\delta} \right)^2 \right]$

very small distance

Eq (2) becomes

$c_A(z)$  at  $y=L$

where  $v_{max} = \frac{3}{2} v_{y,avg}$

$t = \frac{y}{v_{max}}$

$\frac{\partial c_A}{\partial (y/v_{max})} = D_{AB} \frac{\partial^2 c_A}{\partial z^2}$  --- (9)

$c_A = c_{a,i}$  at  $z=0$

$c_A = 0$  at  $y = \infty$

$c_A = 0$  at  $z = \infty$

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So, if we take equation 4 and 5 into consideration then  $v_y$  is equal to  $v_{y, max} \left[ 1 - \frac{z}{\delta} \right]^2$  where  $v_{y, max}$  is equal to  $\frac{3}{2} v_{y, average}$ . So, if the solute is penetrated only a very small distance, very small distance into the fluid and that is a short contact time that is  $t$  is equal to  $y$  over  $v_{y, max}$ , then the solute  $a$  that has diffused this has been carried out along with the velocity  $v_{y, maximum}$ . So, the previous equation 2 which becomes  $\frac{\partial c_a}{\partial y} = \frac{D_A b}{\nu} \frac{\partial^2 c_a}{\partial z^2}$  this is equation number 9. So, if we use  $a$  if we use  $c_a$  is equal to 0 at  $y$  is equal to 0 and  $c_a$  is equal to  $c_{a, i}$  at  $z$  is equal to 0 and  $c_a$  is equal to 0 at  $z$  is equal to infinite. So, if we integrate equation 9 then it can be this is equation 9 if we integrate then  $c_a$  over  $c_{a, i}$  this is equal to  $\frac{f c}{z^4} \frac{D_A b y}{\nu}$  this is equation number 10.

$\int \frac{1}{\sqrt{y}} dy = 2\sqrt{y}$

$\frac{C_A}{C_{A_i}} = \text{erfc} \left( \frac{y}{\sqrt{4D_{AB}t}} \right)$

Complementary error function of  $y$

$\text{erfc } y = 1 - \text{erf } y$

$z=0$  at any position  $y$

$N_A = -D_{AB} \frac{\partial C_A}{\partial y} \Big|_{z=0}$

$= C_{A_i} \frac{D_{AB} \sqrt{y_{max}}}{\tau y}$

$\text{--- Eq 11}$

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Now, where this particular is a complementary error, error function of  $y$  it can be written as  $\text{Erfc } y$  is equal to 1 minus  $\text{Erf } y$  the standard tabulated functions. So, the local molar flux at the surface  $z$  is equal to 0 at any position  $y$  from the top of the entrance can be given as  $n_A$  is equal to minus  $D_{AB} \frac{\partial C_A}{\partial y} \Big|_{z=0}$  and this can be  $C_{A_i} \frac{D_{AB} \sqrt{y_{max}}}{\tau y}$  this is equation number 11.

total mol of A transferred per second to the liquid over the entire length  $y=0$  to  $y=L$

$N_A(L+1) = (1) \int_0^L N_A \Big|_{z=0} dy = (1) \int_0^L C_{A_i} \left( \frac{D_{AB} \sqrt{y_{max}}}{\pi} \right)^{1/2} \frac{1}{y^{1/2}} dy$

$= (L) C_{A_i} \left( \frac{4D_{AB} \sqrt{y_{max}}}{\pi L} \right)^{1/2}$

$\text{--- Eq 12}$

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So, the total mole of A transferred total mole of A transferred per second to the liquid over the entire length  $y$  is equal to 0 to  $y$  is equal to  $L$  where the vertical surface is unit with which can be calculated like  $n_A L$  into 1 1 0 to  $n_A z$  is equal to  $L C_{A_i} \frac{D_{AB} \sqrt{y_{max}}}{\tau y}$  which we are targeting to the power half into 1 over  $y$  to the power half  $dy$  which is equal to  $L C_{A_i} \frac{D_{AB} \sqrt{y_{max}}}{\tau y}$  this can be written as equation number 12.

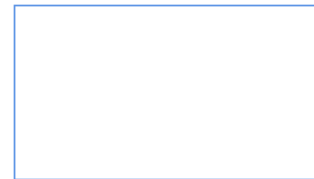


## Problem-1

**Question:** The absorption of pure carbon dioxide is carried out at 1 atmospheric pressure and at 25 degree centigrade by using water film flowing down a vertical wall of 1 meter long. The water is essentially CO<sub>2</sub> free initially. The average velocity of the liquid is 0.2 meter per second. The solubility of CO<sub>2</sub> in water at 25 °C and at 1 atmosphere is  $C_{A,i} = 0.0336$  kmol/m<sup>3</sup>

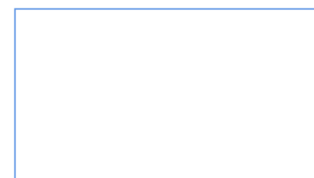
Calculate the film thickness and the rate of absorption of carbon dioxide ?

Use the following properties, DAB:  $2 \times 10^{-9}$  m<sup>2</sup>/s, solution density  $\rho = 997$  kg/m<sup>3</sup> and viscosity  $\mu = 8.95 \times 10^{-4}$  kg/m.s at 25 °C.



Now, let us take up a problem the absorption of pure carbon dioxide is carried out at 1 atmosphere pressure and at 25 degree centigrade by using water film flowing down a vertical wall of 1 meter long and the water is essentially CO<sub>2</sub> free at the outset. The average velocity of the liquid is 0.2 meters per second and the solubility of carbon dioxide in water at 25 degree Celsius and 1 atmosphere is that is  $C_{A,i}$  initially is equal to 0.336-kilo mole per meter cube. So, you need to calculate the film thickness and the rate of absorption of carbon dioxide. Now, some statistical information is given to you like that  $D_{AB}$  is equal to 2 into 10 to the power minus 9-meter square per second and the solution density is equal to  $\rho$  is equal to 997 kilograms per meter cube and the viscosity which is foremost important to solve this particular problem is given as 8.95 into 10 to the power minus 4 kilograms per meter second at 25 degree Celsius.

$$\begin{aligned}
 V_{avg} &= 0.2 \text{ m/s} & \rho &= 997 \text{ kg/m}^3 & \mu &= 8.95 \times 10^{-4} \text{ kg/m.s} & g &= 9.81 \text{ m/s}^2 \\
 \delta &= \left( \frac{34 V_{avg} \mu}{\rho g} \right)^{1/2} = \left( \frac{3 \times 0.2 \times (8.95 \times 10^{-4})}{997 \times 9.81} \right)^{1/2} \\
 \delta &= 2.34 \times 10^{-4} \text{ m} \\
 C_{A,i} &= 0.0336 \text{ kmol/m}^3 & D_{AB} &= 2 \times 10^{-9} \text{ m}^2/\text{s} & L &= 1 \text{ m} \\
 N_A &= C_{A,i} \left( \frac{4 D_{AB} V_{avg}}{\delta L} \right)^{1/2} \\
 N_A &= 0.0336 \left( \frac{4 \times (2 \times 10^{-9}) \times 0.2}{(2.34 \times 10^{-4}) \times 1} \right)^{1/2} \\
 &= 7.58 \times 10^{-7} \text{ kmol/m}^2 \cdot \text{s} \quad \text{Ans}
 \end{aligned}$$

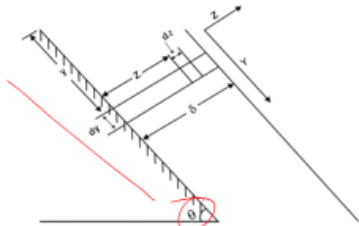


So, let us solve this particular problem. Now, it is given that  $v_y$  average this is equal to 0.2 meters per second,  $\rho$  is equal to 997 kilograms per meter cube,  $\mu$  is equal to  $8.95 \times 10^{-4}$  kilograms per meter second and  $g$  is 9.81 meters per square. So, the  $\delta$  is equal to  $3 v_{\text{average}} \mu$  this is the standard formula that we discussed earlier  $\rho g$  to the power half, and this is  $3 \times 0.2 \times 8.95 \times 10^{-4}$  over  $997 \times 9.81$ , which is equal to one half.

So, the  $\delta$  is equal to  $2.34 \times 10^{-4}$  meters. So, now,  $CAI$  is given to you that is  $0.0336$ -kilo mole per meter cube,  $D_{AB}$  is given to you  $2 \times 10^{-9}$  meter square per second length is given 1 meter. So,  $N_A$  is equal to  $CAI \times 4 D_{AB} v_y \text{ maximum } \tau L$  to the power half. So, if we substitute all these values, then it becomes  $N_A$  is equal to  $0.0336 \times 4 \times 2 \times 10^{-9}$  into  $0.2 \tau_1$  to the power half and this is coming out to be  $7.58 \times 10^{-7}$  kilo mole per meter square second and this is our required answer.

### Laminar falling film in an inclined surface

- In any liquid **flowing down** a surface a velocity profile is established with **the velocity increasing from 0** at the surface itself to a maximum where it is **contact with the surrounding atmosphere**.



**Flow of liquid over a surface**

- **The velocity distribution** may be obtained in a manner similarly used in **connection with the pipe flow** but that **the driving force is due to gravity rather than the pressure gradient**.

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Duggal R. E., "Mass Transfer Operations", 2nd Ed., 1999, McGraw Hill.

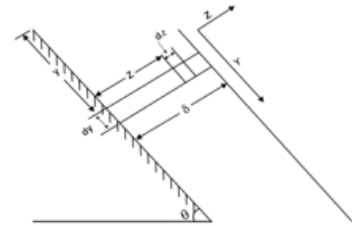
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Now, let us talk about the laminar falling film in an inclined surface. Sometimes this inclined surface plays a very vital role in this aspect.

Now, this is our inclined surface at angle  $\theta$ . So, for any liquid flowing down a surface, a velocity profile which is usually established with the velocity increasing, let us say from 0 at the surface itself to a maximum where it is in contact with the surrounding atmosphere. So, the velocity distribution may be obtained in a manner similar what we used in connection with the pipe flow, but the driving force is attributed to the gravity rather than the pressure gradient. So, for the flow of liquid of the depth of say,  $\delta$  down a plane surface with the width  $W$ , which is inclined at an angle of  $\theta$  to the horizontal. So, a force balance in the  $y$  direction parallel to the surface may be written in this particular aspect.

## Laminar falling film in an inclined surface

- If the drag force of the atmosphere is negligible, the retarding force for laminar flow is attributable to the viscous drag in the liquid at a distance  $y$  from the surface =  $\mu \cdot dV_y/dz \cdot w \cdot dy$



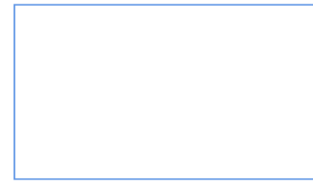
Where,

$V_y$  is the velocity of the fluid at that position.

- Thus, at equilibrium

$$(\delta - z)w \cdot dy \cdot \rho \cdot g \cdot \sin \theta = \mu \frac{dV_y}{dz} w \cdot dy$$

Flow of liquid over a surface



$$(\delta - z)w \cdot dy \cdot \rho \cdot g \cdot \sin \theta = \mu \frac{dV_y}{dz} w \cdot dy$$

Now, in an element of length  $dy$ , the gravitational force acting in the part of the liquid which is at the distance greater than  $Z$  from the surface, this can be written as  $\delta - Z$  into  $W \cdot dy \cdot \rho \cdot g \cdot \sin \theta$ . Now, if we talk about the drag force, so if the drag force of the atmosphere is negligible, the retarding force for laminar flow is attributable to the viscous drag in the liquid at a distance  $y$  from the surface and that is equal to  $\mu \cdot dy \cdot dV_y/dz \cdot W \cdot dy$ . Now, where  $V_y$  this is the velocity of the fluid at the position and therefore the equilibrium thus the equilibrium we can write  $\delta - Z$  into  $W \cdot dy \cdot \rho \cdot g \cdot \sin \theta$  is equal to  $\mu \cdot dV_y/dz \cdot W \cdot dy$ . Now since there will be normally no slip between the liquid and a surface where  $V_y$  is equal to 0 when  $Z$  is equal to 0. So, we can write the integrated integration like from 0 to  $V_y$   $dV_y$  is equal to  $\rho \cdot g \cdot \sin \theta / \mu$  and 0 to  $Z$   $\delta - Z$  into  $dz$  which can be modified to  $V_y$  is equal to  $\rho \cdot g \cdot \sin \theta / \mu$  into  $\delta Z - \frac{1}{2} Z^2$ .

$$\int_0^{V_y} dV_y = \frac{\rho \cdot g \cdot \sin \theta}{\mu} \int_0^Z (\delta - z) dz$$

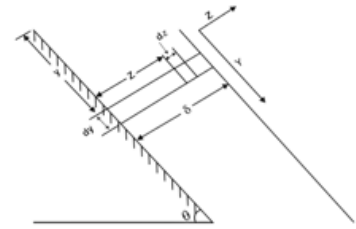
$$V_y = \frac{\rho \cdot g \cdot \sin \theta}{\mu} \left( \delta Z - \frac{1}{2} Z^2 \right)$$

## Laminar falling film in an inclined surface

- Since there will be normally no slip between the liquid and the surface, then  $V_y = 0$  when  $z = 0$ :

$$\int_0^{V_y} dV_y = \frac{\rho \cdot g \cdot \sin \theta}{\mu} \int_0^{\delta} (\delta - z) dz$$

$$V_y = \frac{\rho \cdot g \cdot \sin \theta}{\mu} \left( \delta z - \frac{1}{2} z^2 \right)$$



Flow of liquid over a surface

## Laminar falling film in an inclined surface

- The mass rate of flow that is  $\dot{m}$  of liquid down the surface can be calculated as:

$$\dot{M} = \int_0^{\delta} \frac{\rho \cdot g \cdot \sin \theta}{\mu} w \left( \delta z - \frac{1}{2} z^2 \right) \rho dz$$

$$\dot{M} = \frac{\rho^2 g \cdot \sin \theta}{\mu} w \left( \frac{\delta^3}{2} - \frac{\delta^3}{6} \right)$$

$$\dot{M} = \frac{\rho^2 g \cdot \sin \theta \cdot w \cdot \delta^3}{3 \cdot \mu}$$

$$\dot{M} = \int_0^{\delta} \frac{\rho \cdot g \cdot \sin \theta}{\mu} w \left( \delta z - \frac{1}{2} z^2 \right) \rho dz$$

$$\dot{M} = \frac{\rho^2 g \cdot \sin \theta}{\mu} w \left( \frac{\delta^3}{2} - \frac{\delta^3}{6} \right)$$

$$\dot{M} = \frac{\rho^2 g \cdot \sin \theta \cdot w \cdot \delta^3}{3 \cdot \mu}$$

## Laminar falling film in an inclined surface

- The average velocity of fluid is:

$$V_{y,avg} = \frac{\dot{M}}{\rho \cdot w \cdot \delta} = \frac{\rho \cdot g \cdot \sin \theta \delta^2}{3 \cdot \mu}$$

- For a **vertical surface**,  $\sin \theta = 1$

$$V_{y,avg} = \frac{\rho \cdot g \cdot \delta^2}{3 \cdot \mu}$$

- The **maximum velocity** which occurs at **the free surface** is given by:

$$V_y = \frac{\rho \cdot g \cdot \sin \theta \delta^2}{2 \cdot \mu}$$

And this is **1.5 times the mean velocity** of the liquid

$$V_{y,avg} = \frac{\dot{M}}{\rho \cdot w \cdot \delta} = \frac{\rho \cdot g \cdot \sin \theta \delta^2}{3 \cdot \mu}$$

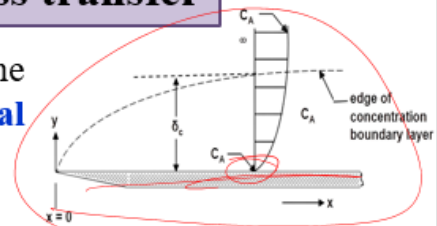
$$V_{y,avg} = \frac{\rho \cdot g \cdot \delta^2}{3 \cdot \mu}$$

$$V_y = \frac{\rho \cdot g \cdot \sin \theta \delta^2}{2 \cdot \mu}$$

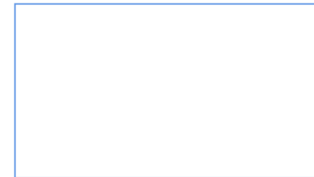
So, the mass rate of the flow, that is the  $\dot{m}$  of the liquid or mass of the liquid down to the surface, can be calculated as  $\dot{m}$  is equal to  $0$  to  $\delta$   $\rho \cdot g \cdot \sin \theta$  over  $\mu$  into  $W \cdot \delta \cdot Z$  minus  $1$  over  $2$   $Z$  square into  $\rho \cdot Z$  which can be after simplification it can be written as  $\dot{m}$  is equal to  $\rho$  square  $g \cdot \sin \theta$  over  $\mu$  into  $W$  into  $\frac{1}{3} \delta^3$  minus  $\frac{1}{2} \delta^3$  over  $Z$  and which can be further modified to this final shape. Now, if we talk about the laminar falling film in an inclined surface, so average velocity of the fluid can be represented as  $v_y$  average is equal to  $\dot{m}$  over  $\rho \cdot W \cdot \delta$  is equal to  $\rho \cdot g \cdot \sin \theta \cdot \delta^2$  over  $3 \cdot \mu$  and for vertical surface  $\sin \theta$  is equal to  $1$ . So, the equation can be modified as  $v$  is equal to  $\rho \cdot g \cdot \delta^2$  over  $3 \cdot \mu$  the maximum velocity that occurs at the free surface is given by  $v_y$  is equal to  $\rho \cdot g \cdot \sin \theta \cdot \delta^2$  over  $2 \cdot \mu$  and this is almost  $1.5$  times the mean velocity of the fluid. Now, let us talk about the boundary layer theory in mass transfer.

## Boundary layer theory in mass transfer

- An exact solution can be obtained for the hydrodynamic boundary layer for isothermal laminar flow past a plate.
- An extension of the Blasius solutions can be extended to derive an expression for convective heat transfer.
- In the analogous manner we can use the Blasius solutions for convective mass transfer as well for the same geometry and laminar flow.



Laminar flow of fluid passes a flat plate and concentration boundary layer



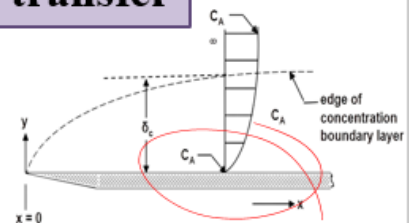
The exact solution can be obtained for the hydrodynamic boundary layer for isothermal laminar flow usually past a plate. Now, this is the typical laminar flow of the fluid that passes a flat plate, this is the flat plate and the concentration boundary layer. So, at this level this is  $C_A$ . Now, an extension to the Blasius solution can be extended to derive an expression for the convective heat transfer. Now, in the analog manner, we can use the Blasius solution for convective mass transfer as well as the same geometry and laminar flow.

## Boundary layer theory in mass transfer

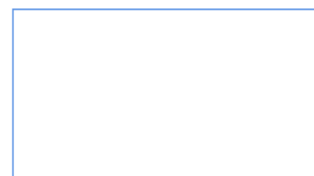
Here,

- $C_{A\infty}$  = is the concentration of A in the fluid approaching the plate.
- $C_{AS}$  = is the concentration of A in the fluid adjacent to the surface
- We start with the differential mass balance and simplifying it for steady state process where,

$$\frac{\partial C_A}{\partial t} = 0$$



Laminar flow of fluid passes a flat plate and concentration boundary layer



$$\frac{\partial C_A}{\partial t} = 0$$

$$V_x \frac{\partial C_A}{\partial x} + V_y \frac{\partial C_A}{\partial y} = D_{AB} \frac{\partial^2 C_A}{\partial y^2}$$

$$V_x \frac{\partial V_x}{\partial x} + V_y \frac{\partial V_x}{\partial y} = \frac{\mu}{\rho} \frac{\partial^2 V_x}{\partial y^2}$$

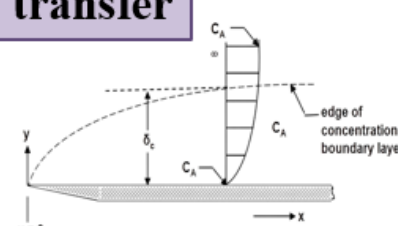
Now, here  $C_A$  infinite is the concentration of A in the fluid approaching the plate. Now,  $C_{AS}$  is the concentration of A in the fluid adjacent to the surface. So, in other words, we can say like this. So, if we start with the differential mass balance and simplify it for steady-state process where  $\text{Del } C$  over  $\text{Del } t$  is equal to 0. Now, flow only direction flow in the direction of x and y.

### Boundary layer theory in mass transfer



- **Flow only in the x and y directions**, so  $\mathbf{V}_z = 0$  and neglecting diffusion in the x and z directions to give:
 

$$V_x \frac{\partial C_A}{\partial x} + V_y \frac{\partial C_A}{\partial y} = D_{AB} \frac{\partial^2 C_A}{\partial y^2} \quad \dots(13)$$
- **The momentum boundary layer** is very similar, So:
 

$$V_x \frac{\partial V_x}{\partial x} + V_y \frac{\partial V_x}{\partial y} = \frac{\mu}{\rho} \frac{\partial^2 V_x}{\partial y^2} \quad \dots (14)$$



Laminar flow of fluid passes a flat plate and concentration boundary layer

Joseph R. E., "Mass Transfer Operations", 3rd Ed., 1990, McGraw Hill

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So, if so,  $v_z$  is equal to 0 and neglecting diffusion in x and z directions. So, it can be given as  $v_x$  into  $\text{Del } C$  over  $\text{Del } x$  plus  $v_y$  into  $\text{Del } C_A$  over  $\text{Del } y$  is equal to  $D_{AB} \text{Del }^2 C_A$  over  $\text{Del } y^2$  and this is can be represented as equation number 13. So, the momentum boundary layer is very similar. Therefore, this equation can be modified as  $v_x$  to  $\text{Del } v_x$  over  $\text{Del } x$  plus  $v_y$  is equal to  $\text{Del } v_x$  over  $\text{Del } y$  plus is equal to  $\frac{\mu}{\rho} \text{Del }^2 v_x$  over  $\text{Del } y^2$  and this can be represented as equation number 14.

The thermal boundary layer is also similar

$$V_x \frac{\partial T}{\partial x} + V_y \frac{\partial T}{\partial y} = \frac{K}{\rho \cdot C_p} \frac{\partial^2 T}{\partial y^2}$$

Dimensionless concentration boundary conditions are;

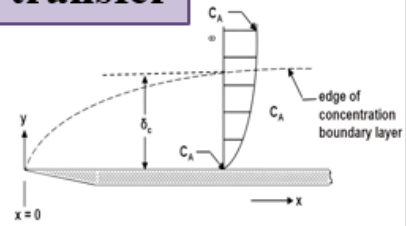
$$\frac{V_x}{V_\infty} = \frac{T - T_S}{T_\infty - T_S} = \frac{C_A - C_{AS}}{C_{A\infty} - C_{AS}} = 0, \text{ at } y=0.$$

$$\frac{V_x}{V_\infty} = \frac{T - T_S}{T_\infty - T_S} = \frac{C_A - C_{AS}}{C_{A\infty} - C_{AS}} = 1, \text{ at } y \text{ is at infinity}$$

## Boundary layer theory in mass transfer

- The thermal boundary layer is also similar:

$$V_x \frac{\partial T}{\partial x} + V_y \frac{\partial T}{\partial y} = \frac{K}{\rho C_P} \frac{\partial^2 T}{\partial y^2} \quad \dots (15)$$

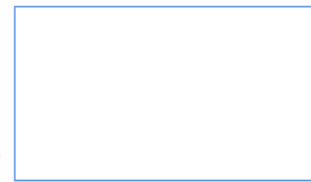


Laminar flow of fluid passes a flat plate and concentration boundary layer

- The dimensionless concentration boundary conditions are

$$\frac{V_x}{V_\infty} = \frac{T - T_S}{T_\infty - T_S} = \frac{C_A - C_{AS}}{C_{A\infty} - C_{AS}} = 0 \quad \text{at } y = 0 \quad \dots (16)$$

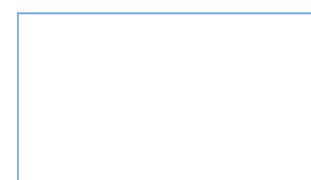
$$\frac{V_x}{V_\infty} = \frac{T - T_S}{T_\infty - T_S} = \frac{C_A - C_{AS}}{C_{A\infty} - C_{AS}} = 1 \quad \text{at } y = \infty \quad \dots (17)$$



So, the thermal boundary layer if we talk about is also similar. So, this can be represented as  $v_x$  into  $\frac{\partial t}{\partial x}$  plus  $v_y$  is equal to  $\frac{\partial t}{\partial y}$  is equal to  $\frac{k}{\rho C_P} \frac{\partial^2 t}{\partial y^2}$ , and this is equation number 15. And if we adopt the dimensionless correct concentration boundary conditions, then at  $y$  is equal to 0, the  $\frac{v_x}{v_\infty}$  is equal to  $\frac{t - T_S}{t_\infty - T_S}$  is equal to  $\frac{C_A - C_{AS}}{C_{A\infty} - C_{AS}}$  is equal to 0. So, this can be represented as equation number 16. And on simplification, if we go for  $y$  is equal to infinite, this equation can be further modified, which is represented here, that is  $\frac{v_x}{v_\infty}$  is equal to  $\frac{t - T_S}{t_\infty - T_S}$  is equal to  $\frac{C_A - C_{AS}}{C_{A\infty} - C_{AS}}$  is equal to 1. So, this is equation number 17.

## Boundary layer theory in mass transfer

$\frac{\mu}{\rho} = N_A = 1.0$   
 $\frac{\mu}{D_{AB}} = N_{Sc} = 1.0$   
 $N_{Re} x = x \frac{\rho V_\infty}{\mu}$   
 $\frac{\sqrt{x}}{V_\infty} = \frac{C_A - C_{AS}}{C_{A\infty} - C_{AS}} = (C_{A\infty} - C_{AS}) \left( \frac{0.332}{x} N_{Re} x \right)$   
 $\left( \frac{\partial C_A}{\partial y} \right)_{y=0}$   
 The velocity gradient  $\left( \frac{\partial v_x}{\partial y} \right)_{y=0} = 0.332 \frac{V_\infty}{x} N_{Re} x$   
 (18)  
 (19)  
 (20)



Now, the Blasius equation is applied to convective heat transfer in that case  $\frac{\mu}{\rho \alpha}$  is equal to  $N_A$ , which can be represented as 1. Now, we use the same type of solution for the



laminar convective mass transfer, where  $\frac{\mu}{\rho D_{AB}}$  is equal to  $N_{Sc}$  that is Schmidt number  $1$  minus  $0$ , and the velocity gradient, the velocity gradient is derived in the fluid mechanics and that is  $\frac{d v_x}{d y}$  at  $y = 0$  that is  $0.332 \frac{v_{\infty}}{x} \sqrt{Re_x}$ . So, where  $Re_x$  that is Reynolds number at  $x$ ,  $v_{\infty}$  is  $\rho v_{\infty}$ , this is equation number 18. Now, if we recall equation number 17 here, then  $\frac{v_x}{v_{\infty}}$  is equal to  $\frac{C_A - C_{A\infty}}{C_{A0} - C_{A\infty}}$ , this is equation number 19.

Fick's eq  
 $N_{AB} = k'_c (C_{A\infty} - C_{A0}) = -D_{AB} \left( \frac{\partial C_A}{\partial y} \right)_{y=0}$  --- (21)

$\frac{k'_c x}{D_{AB}} = N_{Sh,x} = 0.332 \sqrt{Re_x} N_{Sc}$  --- (20)

The relationship is restricted to gases with  $N_{Sc} = 1.0$

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Now, if you differentiate equation 19 and combine the result with equation 18, then it can be written as  $\frac{d C_A}{d y}$  at  $y = 0$  is equal to  $\frac{C_{A0} - C_{A\infty}}{x} \sqrt{Re_x}$ , this is equation number 20. So, the convective mass transfer equation can be written as and also, if we relate with Fick's law, Fick's equation, then  $N_{AB}$  is equal to  $k'_c (C_{A\infty} - C_{A0}) = -D_{AB} \frac{d C_A}{d y}$  at  $y = 0$ , this is equation number 21. So, from equations 20 and 21, we get  $\frac{k'_c x}{D_{AB}}$  is equal to  $0.332 \sqrt{Re_x} N_{Sc}$  and this is equation number 22. So, the relationship, this relationship is restricted to gas, gases with a Schmidt number is equal to  $1.0$ .

Relation between thickness and concentration boundary layer

$$\frac{\delta}{\delta_c} = N_{Sc}^{\frac{1}{3}}$$

Equation for local convective mass transfer coefficient

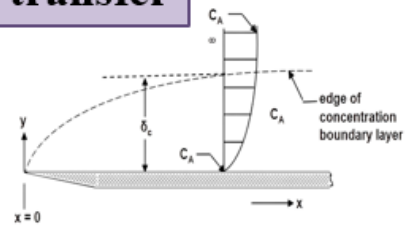
$$\frac{k'_c x}{D_{AB}} = N_{Sh,x} = 0.332 N_{Re,x}^{\frac{1}{2}} N_{Sc}^{\frac{1}{3}}$$

## Boundary layer theory in mass transfer

- The relations between the thickness  $\delta$  of the hydrodynamic and the  $\delta_c$  of the concentration boundary layer, where the  $N_{Sc} \neq 1.0$  is

$$\frac{\delta}{\delta_c} = N_{Sc}^{\frac{1}{3}}$$

..... (23)

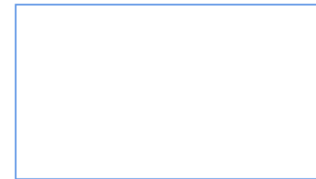


**Laminar flow of fluid passes a flat plate and concentration boundary layer**

- As a result, the equation for local convective mass transfer coefficient is

$$\frac{k'_C x}{D_{AB}} = N_{Sh,x} = 0.332 N_{Re,x}^{\frac{1}{2}} N_{Sc}^{\frac{1}{3}}$$

..... (24)



Now, the relationship between the thickness  $\delta$  of the hydrodynamic and  $\delta_c$  of the concentration boundary layer where Schmidt number is not equal to 1 is  $\delta_c / \delta = N_{Sc}^{-1/3}$  and that is equation number 23. Now, as a result, the equation for local convective mass transfer coefficient is given as  $k'_C x / D_{AB} = N_{Sh,x} = 0.332 N_{Re,x}^{1/2} N_{Sc}^{1/3}$  and this can be represented as equation number 24. So, we can obtain the mass transfer coefficient  $k'_C$  from  $x = 0$  to  $L$  for a plate with  $B$  by integrating this particular equation. So, this can be the  $k'_C$  is equal to  $B / B \int_0^L k'_C dx$  and this is equation number 25. So, we can represent the integration like  $k'_C$  over  $B$  is equal to  $0.664 N_{Re,L}^{1/2} N_{Sc}^{1/3}$ .

Mean mass transfer coefficient

$$\overline{k'_C} = \frac{1}{L} \int_0^L k'_C dx$$

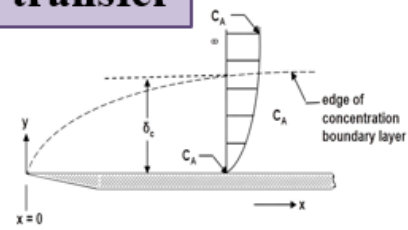
Or

$$\overline{k'_C} \cdot L / D_{AB} = N_{Sh,L} = 0.664 N_{Re,L}^{\frac{1}{2}} N_{Sc}^{\frac{1}{3}}$$

## Boundary layer theory in mass transfer

- We can obtain the **mean mass transfer coefficient**  $\bar{k}_C'$  from  $x = 0$  to  $x = L$  for a plate of width  $b$  by integrating as follows:

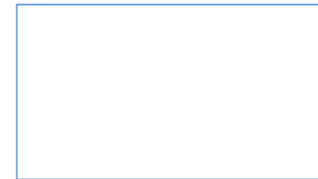
$$\bar{k}_C' = \frac{b}{b.L} \int_0^L k_C' dx \quad \dots (25)$$



**Laminar flow of fluid passes a flat plate and concentration boundary layer**

- The result is

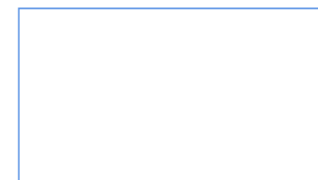
$$\frac{\bar{k}_C' L}{D_{AB}} = N_{Sh,L} = 0.664 N_{Re,L}^{\frac{1}{2}} N_{Sc}^{\frac{1}{3}} \quad \dots (26)$$



Now, let us take another problem, a large volume of pure water at 25 degrees Celsius is flowing parallel to a flat surface of solid benzoic acid where  $L$  is 0.244 meter in the direction of the flow and the water velocity is 0.061 meter per second, and the solubility of the benzoic acid in water is given that is 0.2948 kilo mole per meter cube and the diffusivity of the benzoic acid is  $1.245 \times 10^{-9}$  meter square per second.

### Problem-2

**Question:** A large volume of pure water at 25 °C is flowing parallel to a flat surface of solid benzoic acid, where  $L$  is 0.244 m in the direction of flow. The water velocity is 0.061 m/s. The solubility of benzoic acid in water is 0.02948 kmol/m<sup>3</sup>. The diffusivity of benzoic acid is  $1.245 \times 10^{-9}$  m<sup>2</sup>/s. Calculate the mass transfer coefficient  $\bar{k}_C'$  and the flux  $N_A$ . Given that  $\mu = 8.71 \times 10^{-4}$  kg/m.s and  $\rho = 996$  kg/m<sup>3</sup>.



$$N_{Sc} = \frac{\mu}{\rho D_{AB}} = \frac{8.71 \times 10^{-4}}{996 (1.245 \times 10^{-9})} = 702$$

$$N_{Re} = \frac{L u \rho}{\mu} = \frac{0.244 (0.061) (996)}{8.71 \times 10^{-4}} = 1.7 \times 10^4$$

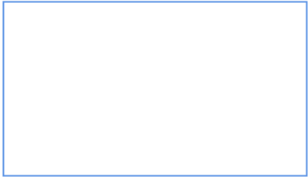
$$\frac{k_c' L}{D_{AB}} = 0.664 N_{Re}^{1/2} N_{Sc}^{1/3}$$

$$k_c' = 3.92 \times 10^{-6} \text{ m/s}$$

You need to calculate the mass transfer coefficient to  $k_c'$  and the flux and given is whatever given that is  $\mu$  is equal to  $8.71 \times 10^{-4}$  kilograms per meter second and density  $\rho$  is 996 kilogram per meter cube. So, let us solve this particular problem. So, the Schmidt number that is  $N_{Sc}$  that is  $\mu$  is equal to  $\rho D_{AB}$  this is  $8.71 \times 10^{-4} / (996 \times 1.245 \times 10^{-9})$  that is equal to 702 and Reynolds number is given by this particular mathematical representation and if you substitute the value, then it can be into  $996 / (8.71 \times 10^{-4})$  and that is  $1.7 \times 10^4$ . So,  $k_c' / D_{AB}$  is equal to  $0.664 \times \text{Reynolds number to the power } 1/2 \times \text{Schmidt number to the power } 1/3$ .

So, if we substitute all the values that we calculated and  $D_{AB}$ , which is given then  $k_c'$  comes out to be  $3.92 \times 10^{-6}$  meters per second. So, if we calculate  $N_A$  we have  $N_A$  is equal to  $k_c' \times (C_{A1} - C_{A2})$  which is equal to  $k_c' C_{A1} - C_{A2}$ . Now, if the solution is very dilute then this can be represented at  $k_c'$  is almost equal to  $k_c$ .

$N_A = \frac{k_c'}{X_{BLM}}$ 
 $(C_{A1} - C_{Av}) = k_c (C_{A1} - C_{Av})$ 
  
 $X_{BLM} = 10$  and  $k_c' \approx k_c$ 
  
 $C_{A1} = 2.948 \times 10^{-4} \text{ kg mol/m}^3$  (solubility)
   
 $C_{Av} = 0$  → large volume of fresh water
   
 $N_A = (3.92 \times 10^{-6}) \times (0.02948 - 0)$ 
  
 $= 1.15 \times 10^{-7} \text{ kg mol/m}^2 \text{ s}$ 
  
 Ans



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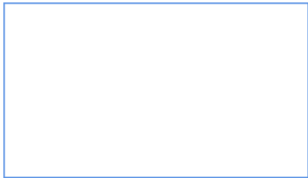
Also, initially it is given the solubility is given 2.948 into 10 to the power minus 4-kilogram mole per meter cube that is solubility and  $C_{A2}$  is equal to 0 for a large volume of fresh water. So, if we substitute all these values it 3.9 N A is equal to 3.92 into 10 to the power minus 6 into 0.02948 minus 0. This comes out to be 1.15 into 10 to the power minus 7-kilogram mole per meter cube. This is our answer.

## Mass transfer coefficient in turbulent flow

- There are **many theories** which attempt to interpret or **explain the behavior of mass transfer coefficient**.

Such as:

- a) **Film theory** → Nernst (1904)
- b) **Penetration theory** → Higbie (1935)
- c) **Surface-renewal theory** → Danckwerts (1951)



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Now, if we talk about the turbulent flow, so there are many theories that attempt to interpret the explaining explain the behaviour of mass transfer coefficient like NIST given the film theory the penetration theory is given by Higbie in 1935 and the surface enable theory is given by Dankwerts in 1951. Now, NIST the film theory the NIST postulates that near the interface, there exists a stagnant film and the stagnant film, which is hypothetical since we really do not know the details of the velocity profile near the interface.

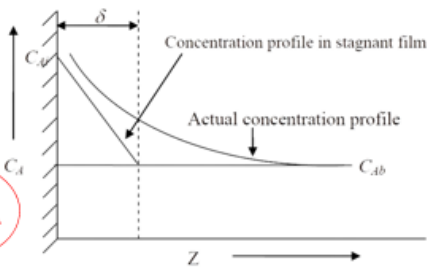
The basic concept is that the resistance to diffusion can be considered equivalent to that in the stagnant film of a certain thickness.

$$N_A = \frac{-D_{AB}dC_A}{dz}$$


**Film theory**


- **Mass transfer occur by molecular diffusion** through the fluid layer at phase boundary that is at solid wall. Beyond this film the **concentration is homogeneous and is  $C_{Ab}$** .
- Mass transfer through the film occurs at **steady state**.
- **Flux is low** and **mass transfer occurs at low concentration**. Hence,
 

$$N_A = \frac{-D_{AB}dC_A}{dz}$$



Concentration profile with stagnant film


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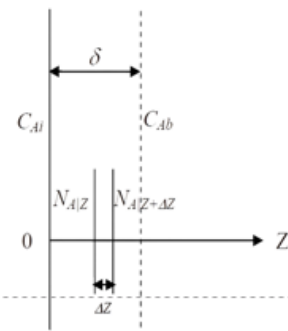
The mass transfer occurs by the molecular diffusion through the fluid layer at a phase boundary that is a solid wall and beyond this particular film the concentration is homogeneous and it is given by  $C_A B$ . So, the mass transfer through the film occurs at a steady state and flux is very low and the mass transfer occurs at low concentration. So therefore, this can be  $N_A$  can be represented like this. Now, a steady state mass balance over the elementary volume thickness if we try to find out at  $\Delta z$ .

$$\lim_{\Delta Z \rightarrow 0} \frac{N_A|_z - N_A|_{z+\Delta Z}}{\Delta Z} = 0$$

$$\frac{dN_A}{dz} = 0 \frac{d}{dz} \left( \frac{-D_{AB}dC_A}{dz} \right) = 0 - D_{AB} \frac{d^2 C_A}{dz^2} = 0 \frac{d^2 C_A}{dz^2} = 0$$

## Film theory

- The steady state mass balance over an elementary volume of thickness  $\Delta Z$ .
- Rate of input of solute at  $Z = N_A|_Z$
- Rate of output of solute at  $Z + \Delta Z = N_A|_{Z+\Delta Z}$



- Rate of accumulation = 0  
= rate of input - rate of output
- Therefore,

At steady state:

$$N_A|_Z - N_A|_{Z+\Delta Z} = 0$$



## Film theory

Consider

$$\lim_{\Delta Z \rightarrow 0} \frac{N_A|_Z - N_A|_{Z+\Delta Z}}{\Delta Z} = 0$$

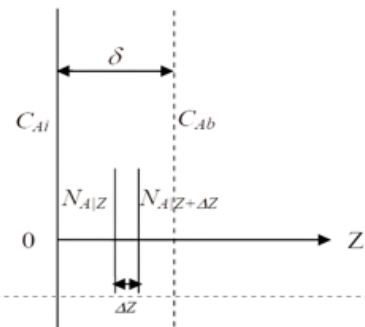
Then,

$$\frac{dN_A}{dz} = 0 \quad \dots (27)$$

$$\frac{d}{dz} \left( -D_{AB} \frac{dC_A}{dz} \right) = 0 \quad \dots (28)$$

$$-D_{AB} \frac{d^2 C_A}{dz^2} = 0 \quad \dots (29)$$

$$\frac{d^2 C_A}{dz^2} = 0 \quad \dots (30)$$



So, the rate of input solute at  $z$  is equal to  $N_A$  at  $z$  and the rate of output can be given as like this. So, if we take the rate of accumulation is equal to 0 that is the rate of input minus the rate of output; therefore, at a steady state, mathematically it can be represented like this. Now, let us consider that if we go for  $z$  tends to 0, then  $N_A$  at  $z$  minus  $N_A$  at  $z + \Delta z$  over  $\Delta z$  is equal to 0. So, upon simplification we can go we can obtain various equations like  $dN_A$  over  $dz$  is equal to 0 and then  $d$  over  $dz$  minus  $dAB$   $dC_A$  over  $dz$  is equal to 0. So, we can obtain various equations which are useful in due course of time.

$$C_A = C_{A,i} - (C_{A,i} - C_{A,b}) \frac{z}{\delta}$$

$$N_A = -D_{AB} \frac{dC_A}{dz} \Big|_{z=0}$$

$$N_A = \frac{D_{AB}(C_{A,i} - C_{Ab})}{\delta}$$

## Film theory

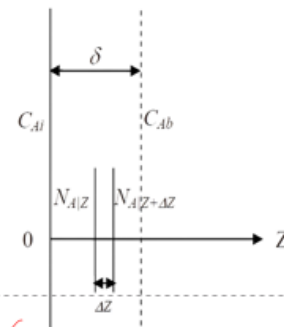
- Integrating equation 3 for the following boundary conditions

$$C_A = C_{A,i} \text{ when } Z = 0$$

$$C_A = C_{Ab} \text{ when } Z = \delta$$

We have now:

$$C_A = C_{A,i} - (C_{A,i} - C_{Ab}) \frac{Z}{\delta}$$



- Hence, according to film theory

**Concentration profile in stagnant film is linear**

- Molar flux through film,  $N_A$ :

Therefore,

$$N_A = -D_{AB} \left. \frac{dC_A}{dz} \right|_{Z=0}$$

$$N_A = \frac{D_{AB}(C_{A,i} - C_{Ab})}{\delta} \dots (31)$$



Now, if we integrate the previous equation, which we developed over like equation number 3, then with the different conditions, a different boundary layer condition that is  $C_A$  is equal to  $C_{A,i}$  when  $z$  is equal to 0, and  $C_A$  is equal to  $C_{A,b}$  where  $z$  is equal to the delta. So, we have now mathematically represented like this. Therefore, according to the film theory, the concentration profile in stagnant film is linear, and the molar flux through  $N_A$  is given by these two equations.

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So, dear friends, at the outset, we discussed the mass transfer operation in various aspects like falling film, thickness all those things, and for your convenience, we have enlisted a couple of references you can utilize as per your requirement. Thank you very much.