

Heat Transfer
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Lecture - 26
Heat and Momentum Transfer Analogy

So, we were discussing about the dimensionless form of the momentum equation, dimensionless form of the energy equation, the boundary conditions again in dimensionless forms; essentially for the case of fluid flow, the no slip velocity. And what would be the condition on what will be the condition of velocity at a point plot from the plate such that at that point the velocity is going to be equal to the local free stream velocity outside of the boundary layer. And similarly, we are also looking at the energy equation in what would be, what would be the form of the boundary conditions?

For example, what is going to T^* that is the dimensionless temperature at any actual location; but on the plate itself? So, T^* at x^* comma 0; that means, y^* equals to 0 would be equal to 0 because of the way we have defined the dimensionless temperature T^* . But T^* were T^* was simply $T - T_s$ by $T_\infty - T_s$. So, on the plate T is equal to T_s ; therefore, T^* would be equal to 0. At a point plot from the plate, the temperature of the fluid would simply be equal to T_∞ and the value of T^* in that case would be equal to 1.

So, we were looking at 2 equations to governing equations for 2 processes; one for heat transfer and the other for momentum transfer. And we saw that the ones that separate, the combination of terms that separate these 2 equations that that differentiates between these 2 equations are the presence of similarity parameters. One is the Reynolds number for the case of momentum transfer and the second is Reynolds times Prandtl number for the case of heat transfer.

So, these are the only difference between Heat transfer and Momentum transfer. So, what we would like to do is we have we have proposed then that if we could keep the Reynolds number to be the same for heat transfer as well as or momentum transfer and if you can choose a hypothetical fluid with a Prandtl number to be equal to 1, then these two equations dimensionless form of these 2 transfer equations are identical.

And if in addition, we assume that the flow is taking place over a flat plate, then the governing equation the boundary conditions of the governing equations are also going to be identical. So, that is the case of dynamic similarity which they tells us that that a for a dynamic similar system, the expression of dependent variable for the case of momentum transfer which is u^* can be replaced by the replaced by the dependent variable of the other equation which is T^* .

And therefore, and analogy, a similarity and equableness between the momentum transfer and heat transfer can be established to obtain expressions, known expressions of one dependent variable from the known expression of another independent another dependent variable. So, will be look at that it will be very clear towards the end of this class, how it is done.

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The image shows handwritten notes on a grid background, organized into three columns: **Gov. EON**, **SIMILARITY PARAMETER**, and **BOUNDARY CONDITIONS**.

Gov. EON

Momentum equation:
$$u^* \frac{\partial u^*}{\partial x^*} + v^* \frac{\partial u^*}{\partial y^*} = -\frac{dP}{dx^*} + \frac{\nu}{L} \frac{\partial^2 u^*}{\partial y^{*2}} = \frac{1}{Re_L}$$

Energy equation:
$$u^* \frac{\partial T^*}{\partial x^*} + v^* \frac{\partial T^*}{\partial y^*} = \frac{\alpha}{L} \frac{\partial^2 T^*}{\partial y^{*2}} = \frac{1}{Re_L Pr} = \frac{\nu}{L} \frac{1}{c_p \mu}$$

SIMILARITY PARAMETER

Momentum: Re_L

Energy: Re_L, Pr

BOUNDARY CONDITIONS

WALL

Momentum: $u^*(x^*, 0) = 0$
 $v^*(x^*, 0) = 0$
NO SLIP.

Energy: $T^*(x^*, 0) = 0$
 $T^* = \frac{T - T_s}{T_\infty - T_s}$

FREESTR.-EAM

Momentum: $u^*(x^*, \infty) = \frac{U_\infty(x^*)}{V} \approx 1$ (FLAT PLATE)

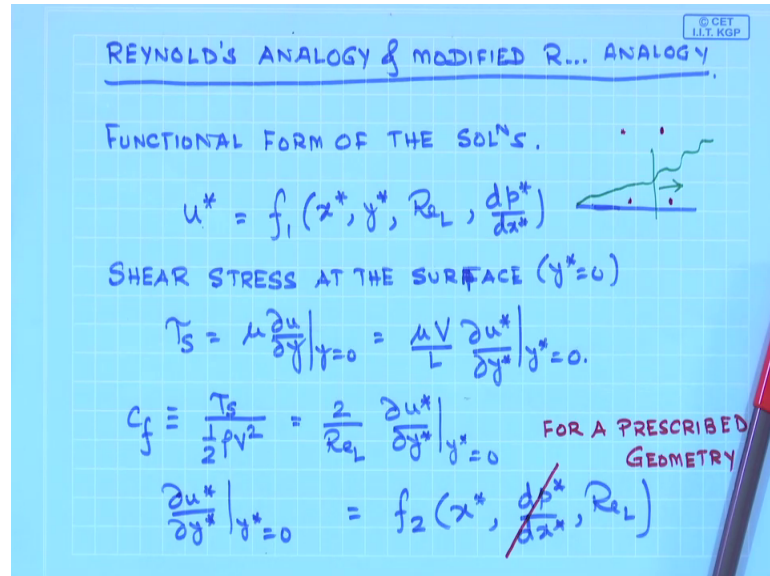
Energy: $T^*(x^*, \infty) = 1$

So, let us look at the this slide which was the last slide on the previous class, where I have identified the governing equations, the similarity parameters, Reynolds number and Reynolds and Prandtl number. This is for momentum; this is for energy and the boundary conditions using no slip and a point plot from the flat plate, what would be the velocity condition? The temperature at y equals to 0 and temperature at y equals to infinity.

So, with this knowledge when we by keeping the Prandtl number to be equal to 1 and keeping the Reynolds number to be the same and assuming that the flow is taking place over a flat plate; everything in this equation in the between these equations and the

boundary conditions are identical, so we have a similar system, dynamically similar system.

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So, what I am going to what we going to do is find out what is Reynolds analogy and modified Reynolds analogy. So, for that I am going to look at the function what could be the functional form of u star. I do not know what would be the exact form of it; but I know that if I could write if I could write the u star if I could write the functional form of u star, it should contain the independent variable x star, independent variable y star, the similarity parameter Reynolds number and the pressure gradient present in the system which is d p star d x star.

So, my functional form of u star is going to be x star y star Reynolds number based on the entire length of the plate and d p star d x star. I do not know how u is going to be connected with x y or Reynolds number, but I know that a functional form like this would exist for the case of flow. Now in terms of engineering interest we would like to find out what is the shear stress at the surface? That means, at this by at by at the surface, I mean at y star to be equal to 0 that is at the surface.

So, which I let us call it as tau s, the shear stress which would be mu times del u del y at y equals to 0. So, this is the shear stress at the surface and if I if I non-dimensionalize it; it is simply going to be mu u by L del u star by del y star at y star equals to 0.

So, that would give me the expression for shear stress and shear stress coefficient, we understand that by definition its τ_s by the dynamic pressure which is half ρV^2 ; V is the approach velocity, ρ is the density. So, that is a definition of C_f . So, the definition of C_f it can be written as Reynolds number for the entire length $\frac{\rho u_{star}}{\mu} \frac{\Delta y}{\rho u_{star}}$ at $y_{star} = 0$. This was simply obtained by putting the value putting the expression of τ_s over here and observing this half ρV^2 in it to obtain to by a real $\frac{\rho u_{star}}{\mu} \frac{\Delta y}{\rho u_{star}}$ at $y_{star} = 0$.

So, this if I would write take write to find what is the, what is the $\frac{\rho u_{star}}{\mu} \frac{\Delta y}{\rho u_{star}}$ functional form of $\frac{\rho u_{star}}{\mu} \frac{\Delta y}{\rho u_{star}}$ at $y_{star} = 0$? So, if you look at the expression the functional form, the hypothetical functional form of u_{star} , I am trying to find out $\frac{\rho u_{star}}{\mu} \frac{\Delta y}{\rho u_{star}}$ at $y_{star} = 0$. Since, I am assigning a specific value of y_{star} to be equal to 0; this must be a function of x_{star} $\frac{d p_{star}}{d x_{star}}$ and Reynolds number based on length. Since, I have specified the value of y_{star} to be equal to 0. So, therefore, the y_{star} does not appear over here.

Now, this is the flow; this is a flat plate over which the flow is taking place and this site is the turbulent flow. Now, if the geometry is prescribed, then you would be able to obtain $\frac{d p_{star}}{d x_{star}}$ separately. So, this for a prescribed geometry, I will I will explain on it in a moment. Remember that what I have told you before is between in inside the boundary layer, the flow is viscous; outside of the boundary layer, the flow is in viscous. So, there is no effect of viscosity in here. Since, you have viscosity present in effect of viscosity present inside the boundary layer, you cannot use known equations which are available to give which are there to provide what is the pressure drop as a function of distance.

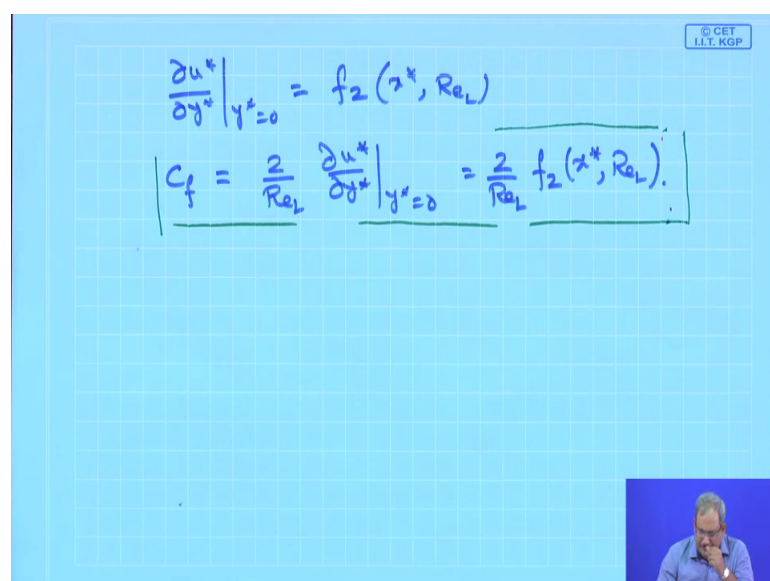
Now, when you when if you if you if someone tells you that what is the equation that provides the pressure drop in a flow? The name that comes to your mind is Bernoulli's equation because Bernoulli's equation would relate the pressure gradient the pressure head, the velocity head and the gravity head. Now, if I assume the plate to a horizontal which is the case in this case. So, therefore, the it is going to be the summation of pressure head and velocity head to be constant. So, if I know this velocity or I can express the change in pressure in terms of change in velocity head that is what Bernoulli's equation is all about. Now, there is catch though; the Bernoulli's equation is strictly valid for in viscid flow for flow where the effect of viscosity is absent.

So, inside the boundary layer, technically I cannot use Bernoulli's equation as the flow is viscous there. So, this solution; but the observation is outside of the boundary layer the flow is inviscid. So, if the geometry is known to me then I would be able to use Bernoulli's equation in the flow domain outside of the boundary layer to obtain an expression for $\frac{dp}{dx}$ or $\frac{dp^*}{dx^*}$ independent of everything.

So, if someone gives me the geometry I should be able to obtain, what is $\frac{dp^*}{dx^*}$ outside of the boundary layer through the use of Bernoulli's equation and since the thickness of the boundary layer is very small, there is no change in pressure with y . That is an assumption which is, a valid assumption considering the small thickness of the boundary layer. So, I use Bernoulli's equation to find out what is $\frac{dp^*}{dx^*}$. So, $\frac{dp^*}{dx^*}$ can be obtained and for a prescribed geometry $\frac{dp^*}{dx^*}$ is a constant; for that reason the term $\frac{\partial u^*}{\partial y^*}$ at $y^* = 0$ which had otherwise contained $\frac{dp^*}{dx^*}$, I can drop that. Since for a given geometry this pressure gradient is known to me after every and is a constant.

So, in terms of functional form of the equation whatever I have written over here can simply be written once again as $\frac{\partial u^*}{\partial y^*}$ at $y^* = 0$ is a function only of x^* and Reynolds number I need not mention $\frac{dp^*}{dx^*}$ for a prescribed geometry.

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$$\frac{\partial u^*}{\partial y^*} \Big|_{y^*=0} = f_2(x^*, Re_L)$$

$$C_f = \frac{2}{Re_L} \frac{\partial u^*}{\partial y^*} \Big|_{y^*=0} = \frac{2}{Re_L} f_2(x^*, Re_L)$$

So, in other words my $\frac{\partial u^*}{\partial y^*}$ at $y^* = 0$ is simply going to be f^2 the function the unknown function x^* times Reynolds number $x^* \text{Re} L$. Now, now if use C_f which I have seen to be equals 2 by $\text{Re} L \frac{\partial u^*}{\partial y^*}$ at $y^* = 0$ which I have seen in here. This is the, this is where I have obtain the expression for C_f . So, my C_f is simply going to be 2 by $\text{Re} L f^2$ of x^* and $\text{Re} L$.

So, my C_f therefore, would be 2 by $\text{Re} L f^2$ the yet to be evaluated functional function of x^* and $\text{Re} L$. So, these are the 2 equations that one needs to take a look at. First of all u is a function of all the independent variables, the operational parameter and the pressure gradient. From there I obtained the shares stress; from the shear stress, I obtained C_f and for $\frac{\partial u^*}{\partial y^*}$ at $y^* = 0$, I obtained the functional form for this special case when the geometry is known to me. So, this should give me the expression for C_f for flow momentum transport inside a boundary layer. Now, let us see what is going to happen to the temperature profile? So, if I look at the temperature in expression over here, over here that we have obtained.

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Gov. Eqn

$$u^* \frac{\partial u^*}{\partial x^*} + v^* \frac{\partial u^*}{\partial y^*} = -\frac{dP^*}{dx^*} \left(\frac{\nu}{VL} \right) \frac{\partial^2 u^*}{\partial y^{*2}} = \frac{1}{\text{Re} L}$$

$$u^* \frac{\partial T^*}{\partial x^*} + v^* \frac{\partial T^*}{\partial y^*} = \frac{\alpha}{VL} \frac{\partial^2 T^*}{\partial y^{*2}} = \frac{1}{\text{Re} L \text{Pr}} = \frac{\nu}{VL} \frac{k}{c_p \mu}$$

SIMILARITY PARAMETER

$\text{Re} L$

$\text{Re} L, \text{Pr}$

BOUNDARY CONDITIONS

WALL

$u^*(x^*, 0) = 0$

$v^*(x^*, 0) = 0$

NO SLIP.

$T^*(x^*, 0) = 0$

$T^* = \frac{T - T_s}{T_\infty - T_s}$

FREESTR-EAM

$u^*(x^*, \infty) = \frac{U_\infty(x^*)}{V} \approx 1$ (FLAT PLATE)

$T^*(x^*, \infty) = 1$

My temperature profile T^* would be function of $u^* x^* v^* y^* \text{Re} L$ number and Prandtl number; but this u^* and v^* are already a function are already known function of x^* and y^* and so on. For example, in this expression itself what we have seen is that u^* is a function of once; you specify $x y \text{Re} L$ and $d p d x$, u^* star is specified.

So, in the governing equation here you do not need to write T star is a function of u star because the moment you write T star is a function of x star y star and Reynolds number, you essentially specify u star. So, by incorporating u star once again in your a functional form that would be simply a repetition.

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$$\frac{\partial u^*}{\partial y^*} \Big|_{y^*=0} = f_2(x^*, Re_L)$$

$$C_f = \frac{2}{Re_L} \frac{\partial u^*}{\partial y^*} \Big|_{y^*=0} = \frac{2}{Re_L} f_2(x^*, Re_L)$$

$$T^* = f_3(x^*, y^*, Re, Pr, \frac{dp^*}{dx^*}) \Rightarrow$$

$$q_s = -k_f \frac{\partial T}{\partial y} \Big|_{y=0} \Rightarrow -k_f \frac{\partial T}{\partial y} \Big|_{y=0} = h(T_s - T_\infty)$$

$$h = -\frac{k_f}{L} \frac{(T_\infty - T_s) \frac{\partial T^*}{\partial y^*} \Big|_{y^*=0}}{(T_s - T_\infty)}$$

So, therefore, based on the knowledge of the of this governing equation, one should be able to write the functional form T star to be equal to a function f 3 and I put it as x star y star Reynolds number Prandtl number in d p star d x star. This d p star d x star I am keeping it just as just as to make it complete.

But we understand that for a prescribed geometry, I can drop this d p star d x star. So, the same way I have done it for the case of shear stress. I am going to write the same thing for the case of surface heat flux which I call it as q s. So, this is a solid plate, you have the profile and have flow taking place; I am trying to find out what is the surface heat flux at y star equal to 0. So, the surface heat flux is k thermal conductivity of the fluid times del T by del y at y equals to 0.

So, that is the Fourier law equivalent. That is a Fourier's law in which can be expressed as minus k f del T del y at y to be equal to 0 divided by T s minus T infinity and this is going to be equal to h. Because my q s, this is the equality of conduction and convection at this point, at the point where the liquid molecules by due to no slip are stuck to the solid.

So, the heat transfer from the immobile liquid molecules to the mobile liquid molecules, there you have the conduction and convection equality. So, this q_s can be expressed in terms of Fourier's law; this q_s can also be expressed in terms of Newton's law which is h times T_s minus T_∞ . So, h times T_s minus T_∞ is also equal to. So, these 2 are simultaneously valid at y equals to 0 and therefore, the expression for h can be obtained in this fashion.

So, when you express it in dimensionless form this becomes h is equal to minus k_f by L T_∞ minus T_s by T_s minus T_∞ times $\frac{\partial T^*}{\partial y^*}$ at y^* equals to 0. So, this gives I am slowly moving towards the dimensionless form of the expression over here.

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$$h = \frac{k_f}{L} \left. \frac{\partial T^*}{\partial y^*} \right|_{y^*=0}$$

$$\frac{hL}{k_f} = Nu = \left. \frac{\partial T^*}{\partial y^*} \right|_{y^*=0}$$

LOCAL $\rightarrow Nu_x = f_4(x^*, Re_L, Pr)$, (FOR A PRESCRIBED GEOMETRY)

AVG. $\rightarrow \bar{Nu} = f_5(Re_L, Pr)$.

REYNOLD'S ANALOGY $dp/dx=0, Pr=1$

EXPR. of U^* & T^* MUST BE IDENTICAL

$f_1 = f_3$

TRUE ALSO FOR FRICTION COEFF. & Nu

$f_2 = f_4$

So, when you do that when you cancel that the numerator and the denominator what you get is h is equal to k_f the thermal conductivity of the solid times $\frac{\partial T^*}{\partial y^*}$; y^* to be equal to 0 or in other words, you can write hL by k_f is equal to $\frac{\partial T^*}{\partial y^*}$ at y^* equals to 0.

So, what is hL by k_f this is nothing but the Nusselt number. So, we it in convection, we always try to find what is h or what is the expression for Nusselt number? So, now, I write a Nusselt number is used for f_1 f_2 and f_3 over here. So, Nusselt number is $\frac{\partial T^*}{\partial y^*}$ that that y^* equal to 0. So, when I say its $\frac{\partial T^*}{\partial y^*}$ at y^*

equal to 0 that function should be a function of x^* Reynolds number and Prandtl number provided the geometry is known to us.

So thus, this Nusselt number expression would be some function f_4 ; I do not know what this f_4 would be? But, some function of x^* Reynolds number sorry x^* Reynolds number and Prandtl number. So, this is obviously, for a prescribe geometry and if would like to find out what is the average value of Nusselt number, length average value of Nusselt number; the moment you do that, the length average value of Nusselt number; then, x^* is obviously drop it should be another function f_5 $Re L$ times Pr .

So, this is the local value of Nusselt number; this is nu_x and this is the. So, this is local value of Nusselt number and this is the average value of Nusselt number and the bar over Nu simply denotes it is the average value which is the function of this in the f . For length average value, it would simply be a function of Reynolds number and Prandtl number.

Now when we when we use the Reynolds condition, Reynolds analogy; what is at dp/dx is 0 and Prandtl number is equal to 1 and if that is the case, then the expression of u^* and T^* must be identical. This is what we were discussing so far. So, the expressions of u^* and T^* must be identical. So, what is expression of T^* and u^* ? So, u^* is f_1 and T^* is f_3 . So, if your if your Prandtl number is equal to 1. So, the equation becomes dynamically similar; dp/dx is in is the dependence of dp/dx is not there.

So, therefore, f_1 must a f_1 must be equal to f_1 and f_3 ; f_1 and f_3 are going to be identical ok. So, f_1 and f_3 are identical. It is also then true that the expression for the friction coefficient which is this f_2 must also be equal to the f_4 which is which is the relation for this case. So, expression for u^* and T^* must be identical would simply give you that f_1 is equal to f_3 ok.

In true also for friction coefficient and Nusselt number; so, if it is true for friction coefficient Nusselt number what you would get is f_2 is equal to f_4 . So, these are collectively known as the Reynolds analogy. The important point here is the major problem that you one would face in the practical application of Reynolds analogy is the requirement that Prandtl number has to be equal to 1.

Where are you going to get a fluid who is Prandtl number is equal to 1 and if it is equal to 1, how are you going to use this analogy for other cases? So, f_3 f_2 is equal to f_4 ; how does how does that help us? f_4 is this, f_4 and f_2 if these 2 are identical; if f_2 and f_4 are identical, I will write these 2 equations once again to show how we can use them use them in this case.

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$$C_f = \frac{2}{Re_L} \left. \frac{\partial u^*}{\partial y^*} \right|_{y^*=0} = f_2(x^*, Re_L).$$

$$Nu = f_4(x^*, Re_L, Pr).$$
 IF f_2 & f_4 ARE EQUIVALENT.

$$\frac{C_f \cdot Re_L}{2} = Nu \quad \text{--- REYNOLD'S ANALOGY (Pr=1)}$$

$$\frac{C_f}{2} = \frac{Nu}{Re_L \cdot Pr} = St \quad (St = \text{STANTON NO.})$$

$$\boxed{\frac{C_f}{2} = St} \quad \text{'R' ANALOGY}$$

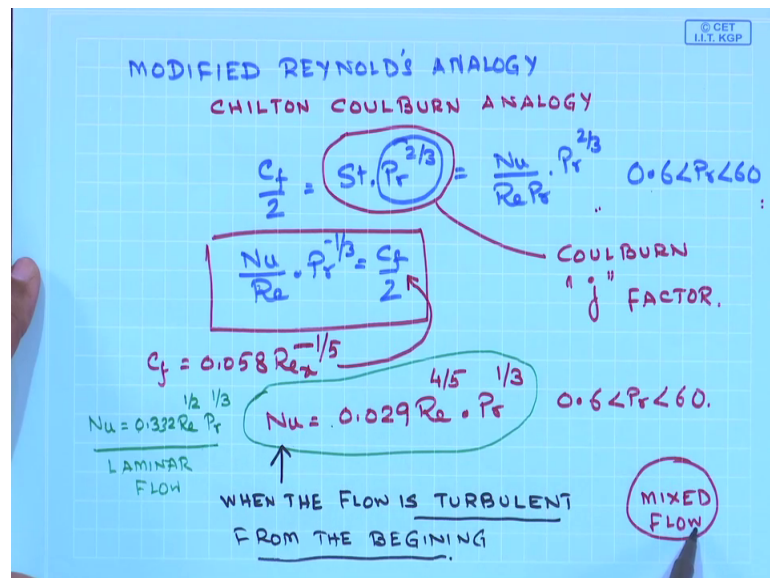
So, C_f is 2 by Re_L del u^* by del y^* at y^* equal to 0 and we understand that del u^* by del y^* at y^* to be equal to 0 is equals f_2 x^* times Reynolds number based on length and for the case of Nusselt number, Nusselt number is simply equal to f_4 x^* Re_L times Prandtl number ok. So, if f_2 and f_4 are equivalent, then what we can say is that C_f , C_f from here times Re_L by 2 . So, C_f times Re_L by 2 is this which is f_2 must be equal to Nusselt number. So, if f_2 and f_4 are equivalent, then C_f times Re_L by 2 must be equal to Nusselt number.

So, this is known as Reynolds Analogy. This is a some extremes in some cases, it is modified in a exactly different way; where it is written that C_f by 2 is equal to Nusselt by Reynolds number based on the length. And since the value of Prandtl number is equal to 1 , there is no harm in adding a Prandtl number in this case. I can do that since Prandtl number in Reynolds analogy is equal to 1 . So, this Nusselt by Reynolds into Prandtl has a special name which is called Stanton Number. So, I can use Stanton number the I can introduce Stanton number. So, this is the value of Prandtl number is equal to 1 .

So, the general form of Reynolds analogy is $C_f/2$ is equal to Stanton number. This is the commonly used form of Reynolds analogy. So, this connects the key engineering parameter of C_f in fluid friction with h on Nusselt number in convective heat transfer. So, I would also like to draw your attention to the previous slide that I was showing Nusselt number is equal to ΔT^* by Δy^* at $y^* = 0$. This again reinforces my statement that the significance of Nusselt number is nothing but the dimensionless temperature gradient at the solid liquid interface.

So, that would be the definition of Nusselt number. The more important 1 is a Nusselt number contains h ; this is an engineering parameter and here I connect Nusselt number with C_f friction coefficient which also is an engineering parameter. So, through the use of this analogy, I connect the heat transfer with momentum transfer; but there is as I understand, there is a problem that is only valid for the case when Prandtl number is equal to 1. So, therefore, in order to extend the validity of Reynolds analogy 2 situations; 2 fluids whose Prandtl number may not be equal to 1; a correction factor is added to this analogy and then, it takes the is called the modified Reynolds analogy.

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And is also known as the Chilton Coulburn Analogy to extend the, to extend the Reynolds analogy. A correction factor is added to this as Stanton into Prandtl to the power 2 by 3. So, this is the correction factor which is added in the Stanton Prandtl to the power 2 by 3 is Nusselt by Reynolds into Prandtl into Prandtl to the power 2 by 3 and

this extends the Prandtl number to a large range of Prandtl number. So, what you get then is Nusselt by Reynolds in $2 \text{ Prandtl}^{-1/3}$ is equal to $C_f^{1/2}$ and this whole thing Stanton into Prandtl to the power $2/3$ this is called the Colburn "j" factor.

So, this is the expression for modified Chilton modified Reynolds analogy or Chilton Colburn analogy and the validity of this is extended in most of the real systems real fluids, they have Prandtl number in the range; except for heavy oils which has Prandtl number more than 60 and the other extreme is liquid metals which has Prandtl number way below 0.6. So, for heavy metals sorry liquid metals and heavy oils, if we exclude these 2 special type of fluids most of the liquids the most of the fluids that you ordinarily use, ordinarily come across would be in this range. And therefore, the Chilton Colburn analogy extends this for a wide range of Prandtl number.

The advantage, what is the advantage? The advantage is as I mentioned C_f expression is already known to us $Re^{-1/2} Pr^{-1/3}$; put it in here and what you get is an expression for Nusselt number as $0.029 Re^{4/5} Pr^{1/3}$. The range of validity between Prandtl number 0.6 and 60. See the beauty of it. This is something which is really interesting. You have got an expression for Nusselt number, you have got an expression for h by simply using an analogy which has solid foundation. So, you the expression for C_f is known to you; you are looking at the governing equations, non-dimensionalizing the governing equations; the similarity parameters clearly obtained out of this exercise.

You look at the dimensionless boundary conditions; see under which condition these 2 equations governing equations become dynamically similar. The moment they become dynamic similar, the solution of one can be used as the solution of the other. So, $\frac{\partial u}{\partial y}$ at $y = 0$ which is connected with C_f can be substituted by $\frac{\partial T}{\partial y}$ at $y = 0$ which is connected with Nusselt number.

So, the gradient of velocity or the gradient of temperature, all in dimensionless form; one related to C_f , the other is related to Nusselt number. The momentum with them dynamical similar, these 2 gradients are identical and what you have then is an expression for C_f and an expression for Nusselt number. The expression for C_f is

already known to you. Therefore, you obtain an expression for the Nusselt number in turbulent flow.

So, without getting into the complicated statistical analysis of AD formation, velocity distribution, unknown velocity distribution, the fluctuations in temperature in velocity; you have a tool now through the use of an analogy and an extended analogy by incorporating Prandtl number corrections, you now have the expression for convective heat transfer coefficient in turbulent flow. That is the beauty of this analysis or this analogy.

So, Reynolds analogy or modified Reynolds analogy which is also known as Chilton Colburn analogy is a powerful tool which lets you find out the expression for h in highly turbulent flow. So, now, I have the complete picture in heat transfer; external heat transfer, flow the heat transfer in external flow simplest possible example flow over a flat plate. I have an expression for h in the early part where the flow is laminar up to a value of Reynolds number 5 into 10 to the power 5. And through the use of analogy, I have an expression for the Nusselt number beyond Reynolds number 5 into 10 to the power 5; that means, when the flow is turbulent.

So, together they give me a complete picture of what would be the heat transfer coefficient in laminar flow and what would be the heat transfer coefficient in turbulent flow? More importantly, this I would show you the next class a corollary of that is the flow is never fully turbulent and the flow can change from laminar to turbulent. So, in most cases any flow has a turbulent portion to sorry a laminar portion to begin with and then, it becomes turbulent.

So, those kind of flows are commonly encountered the known as Mixed flow. The early part its laminar later part it turns turbulent. So, how these relations can be modified to express the average heat transfer coefficient for the case of mixed flow. But that that is there is no new concepts and involved there. What is important is again, I would bring your attention to this equation which simply gives you the Nusselt number for the case of turbulent flow as a function of Reynolds number and as a function of Prandtl number.

I should mention as I was telling you this is when the flow is turbulent from the beginning. So, when the flow is turbulent from the beginning. This expression can be used to obtain to obtain the value of h and so on. But in most of the cases the flow is

laminar to start with and then it turns turbulent those kind of flows are known as Mixed flow.

So, I will give you the expressions or mixed flow based on the expression of Nusselt, Nusselt number in laminar flow and in turbulent flow in the next class. But however, I would once again write the Nusselt number for the case of laminar flow which here just to compare them is $0.3332 Re$ to the power half into Prandtl to the power one-third.

So, this is for laminar flow and this one is for turbulent flow. So, together if you if I combined this and this together what I get is mixed flow. But this is obtain almost completely analytically, this has some approximation built into it; but it gives us the analogy is give us a powerful tool to convert heat transfer data from the momentum transfer obtain an expression for heat transfer and vice versa.

So, will solve her quite a few problems on this to clarify the concepts and to show you how this analogy can be effectively employed in problem solving.