

Novel Separation Processes
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Lecture No # 10
Membrane Separation Processes (Contd.)

Good morning everyone. So, we will be looking into the modeling of unstart batch experiments. Now, unstart batch membrane by separation processes; and in the last class we have elaborated what is the importance of this unstart batch experiments, and we will why we **will** looking at two model them, and how we can generate a large number of data in a short period of time. So, what we did we again we wrote down the concentration species balance equation within the mass transfer boundary layer, and in this case the concentration is a function of time as well as space in the y direction. So, we go reduce the equations into and analyze the equations and found out that a similarity solution can be useful in this case because the governing equation is a parabolic partial differential equation, and one of the boundary condition is residing at infinity.

So, in this case we can have a similarity solution like the cross flow cell cross flow experiment, and we can identify made the equation non dimensional and we formulated the boundary conditions as well and made them non dimensional. Then what we did we did an order of magnitude analysis at the edge of the concentration boundary layer and found a similarity parameter in this particular case. And in this class we will be getting the so, what is the next procedure, next procedure is replacing the partial derivative in terms of similarity derivative of similarity parameter. And we expect the partial differential equation will now reduce to the governing ordinary differential equation and the we look into the solution of the ordinary resultant ordinary differential equation.

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Gov. eqn.

$$\frac{\partial C^*}{\partial \tau} - Pe_w \frac{\partial C^*}{\partial y^*} = \frac{\partial^2 C^*}{\partial y^{*2}}$$

at $\tau = 0$, $C^* = 1$. ✓

at $y^* = 0$, $\frac{\partial C^*}{\partial y^*} + Pe_w C^* R_r = 0$ ✓

at $y^* = \infty$, $C^* = 1$. ✓

$$\eta = \frac{y^*}{\delta^*} = \frac{y^*}{\sqrt{\tau}}$$

$$\frac{\partial C^*}{\partial \tau} = \frac{dC^*}{d\eta} \cdot \frac{\partial \eta}{\partial \tau}$$

$$= \frac{dC^*}{d\eta} \cdot \frac{y^* (-\frac{1}{2})}{\sqrt{\tau}}$$

$$= -\frac{\eta}{2\tau} \frac{dC^*}{d\eta}$$

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So, you should remember the governing equation that we have written in the last class for the transient or the time varying infiltration system is $\frac{\partial C^*}{\partial \tau} = \frac{\partial^2 C^*}{\partial y^{*2}} - Pe_w \frac{\partial C^*}{\partial y^*}$; and we wrote down the initial and the boundary condition at time τ is equal to 0 that is my initial condition; C^* is equal to 1 that is $C = C_0$; at y^* is equal to 0, we have the mixed boundary condition that is the convective flux towards the membrane and the diffusive flux away from the membrane will be equal $\frac{\partial C^*}{\partial y^*} + Pe_w C^* R_r = 0$; and at y^* is equal to infinity, we had C^* is equal to 1.

Now, these boundary condition gives the initial condition; this is the boundary condition lowest to equal to zero; and these are another boundary condition at y^* is equal to infinity, constituting the mass transfer boundary layer is really very small order of magnitude are less compare to the actually dimension of the system. So, it is in fact, on the order of 10 to the minus 6 micron meter. Now, this governing equation is a parabolic partial differential equation, and this boundary condition is infinity. So, we can have a similarity solution, and indeed we can have a similarity solution, we have analyze a situation and found out the similarity parameter in this case as y^* by δ^* , which is nothing but, y^* make it non dimensional; so, it becomes y^* by root over τ . So, this is the similarity parameter we have derived in the last class. So, what is left now, the left is basically to get the derivative of this partial derivative of the original, in terms of

the original coordinate system to be expressed, in terms of this similarity parameter or combine variable parameter, eta. So, instead of two independent variables y star and tau, you will be having a single variable eta and our resultant equation will be boiled down into an ordinary differential equation.

So, let us look into these derivatives. So, $\frac{\partial C^*}{\partial \tau}$ will be nothing but, $\frac{dC^*}{d\eta} \frac{d\eta}{d\tau}$. So, this becomes $\frac{dC^*}{d\eta}$ and $\frac{d\eta}{d\tau}$ will be y^* will be remaining as it is τ to the power minus half right. So, it will be minus half to τ to the power minus three by two. So, it will be nothing but, τ root over τ . Now this becomes y^* by root over τ ; this becomes nothing but eta; so, this becomes eta minus eta by two $\frac{dC^*}{d\eta}$, and you will be having another tau here; so, it will be you will be having eta by 2 tau $\frac{dC^*}{d\eta}$. So, this partial derivative with respect to tau can be expressed in terms of the total derivative with respect to eta. Now, we will look into the other derivatives $\frac{\partial C^*}{\partial y^*}$ and $\frac{\partial^2 C^*}{\partial y^{*2}}$.

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The image shows handwritten mathematical derivations on a blue background. At the top right, there is a small logo for '© GET I.I.T. KGP'. The main derivations are as follows:

$$\frac{\partial C^*}{\partial y^*} = \frac{\partial C^*}{\partial \eta} \cdot \frac{\partial \eta}{\partial y^*}$$

$$= \frac{1}{\sqrt{\tau}} \frac{dC^*}{d\eta} \quad \eta = \frac{y^*}{\sqrt{\tau}}$$

$$\frac{\partial^2 C^*}{\partial y^{*2}} = \frac{\partial}{\partial y^*} \left(\frac{\partial C^*}{\partial y^*} \right)$$

$$= \frac{\partial}{\partial \eta} \left(\frac{\partial C^*}{\partial y^*} \right) \frac{\partial \eta}{\partial y^*}$$

$$= \frac{d}{d\eta} \left(\frac{1}{\sqrt{\tau}} \frac{dC^*}{d\eta} \right) \frac{1}{\sqrt{\tau}}$$

$$= \frac{1}{\tau} \frac{d^2 C^*}{d\eta^2}$$

Below these, there are two more equations:

$$-\frac{\eta}{2\tau} \frac{dC^*}{d\eta} - \frac{P_{ew}}{\sqrt{\tau}} \frac{dC^*}{d\eta} = \frac{1}{\tau} \frac{d^2 C^*}{d\eta^2}$$

$$\Rightarrow -\frac{\eta}{2} \frac{dC^*}{d\eta} - P_{ew} \sqrt{\tau} \frac{dC^*}{d\eta} = \frac{d^2 C^*}{d\eta^2}$$

At the bottom left, there is a logo for 'NPTEL'.

So, $\frac{\partial C^*}{\partial y^*}$ will be $\frac{dC^*}{d\eta} \frac{d\eta}{d y^*}$ and if you remember eta is nothing but, y^* by root over τ . So therefore, $\frac{d\eta}{d y^*}$ will be nothing but, one by root over τ and C^* is now suppose to be a sole function of eta; so, this partial derivative will be replaced by a total derivative $\tau \frac{dC^*}{d\eta}$; and what is $\frac{\partial^2 C^*}{\partial y^{*2}}$, $\frac{\partial^2 C^*}{\partial y^{*2}}$ will be nothing but, $\frac{d}{d\eta} \frac{dC^*}{d\eta} \frac{d\eta}{d y^*}$; now, I replaced $\frac{d\eta}{d y^*}$ by $\frac{d\eta}{d\eta} \frac{d\eta}{d y^*}$

del C star del y star remains same this becomes del eta del y star using the general of the differentiation, and what is this del C star del y star it is already expressed in terms of d C d eta.

So, this becomes d d eta this becomes total derivative, because concentration is now a sole function of eta. So, root over this will be root over tau d C star d eta and del eta del y star will be again one over y root over tau. So, this becomes one over tau d square C star d eta square. So, that way we can express all the partial derivatives in terms of the single variable eta; now, you can substitute in the governing equation this becomes what.

So, you just write it del C star del eta becomes minus eta by 2 tau d C star d eta minus P e w del C star del y star becomes one over root tau d C star d eta one over root over tau d C star d eta; and this becomes d square del square C star del y star square will becomes one over tau d square C star d eta square; so, multiply both side by tau and let us see what you get, this becomes minus eta by 2 d C star d eta minus P e w root over tau d C star d eta is equal to d square C star d eta square. Now, we can combine this two terms and this ordinary differential is second order ordinary differential equation can be written in a neat form.

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$$\frac{d^2 C^*}{d \eta^2} = - \left(\frac{n}{2} + P e w \sqrt{\tau} \right) \frac{d C^*}{d \eta}$$

$$P e w \propto \frac{1}{\delta^*}$$

$$\delta^* \propto \sqrt{\tau}$$

$$P e w \propto \frac{1}{\sqrt{\tau}}$$

$$P e w \sqrt{\tau} = \underline{A} = \text{constant.}$$

$$\frac{d^2 C^*}{d \eta^2} = - \left(\frac{n}{2} + A \right) \frac{d C^*}{d \eta}$$

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And that becomes this d square C star d eta square is equal to minus eta by 2 plus P e w root over tau d C star d eta. So, the governing equation becomes a second ordinary differential equation now this if you remember the permeate flux that is the permeate

flux that you are getting out of the membrane will be pretty small and these will be and the thickness of the concentration boundary will be also it is very thin compare to the any physical dimension of the system. So therefore, we assume that the y component velocity in the concentration boundary layer will be same as that is valid at the wall; although, it is an approximation by it will be a valid approximation there; that means, v within the mass transfer boundary layer, within mass transfer boundary layer will be nothing but, approximated as minus J .

So therefore, and we have seen that this value of the permit flux or the solvent flux will be inversely proportional to the thickness of this δ , less be the thickness of the concentration boundary layer more less resistance it will be offering to against the solvent flux; so, you can expect a higher permeate flux at the beginning of the you know at lower period of term now in this case δ will be function of τ time; so, when the thickness becomes more the permeate flux will be small; so, which is inversely proportional to that; so therefore, non dimensional permeate flux will be inversely proportional to δ^* , and δ^* is directly proportional to $\sqrt{\tau}$, that you have already seen because, the definition if you look into a definition of η ; η is nothing but, y^* by δ^* and the δ^* is proportional to t to τ to the power half. So, it will be one by $\sqrt{\tau}$. So, δ^* is directly proportional to $\sqrt{\tau}$ and therefore, Pe_w will be nothing, but it is inversely proportional to $\sqrt{\tau}$; that means Pe_w times $\sqrt{\tau}$ is A constant, and we call this constant as A.

So, in terms of this new constant now we rewrite our governing equation as $d^2 C^* = \eta^2 C^* - A d C^*$; now, we have to solve these equation at the same time, we have to write down the boundary conditions in terms of η as well; now, since this is an these are second order ordinary differential equation with respect to η ; we must be having two boundary conditions on η to solve this governing equation into ordinary differential equation.

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$\text{at } y^* = \infty, C = C_0 \Rightarrow C^* = 1$
 $\text{at } \eta = 0, C^* = 1$ (1)
 $\text{at } \eta = 0, \left[\frac{dC^*}{d\eta} + AC_m^* R_r = 0 \right]$ (2)
 $\text{at } y^* = 0, \frac{\partial C^*}{\partial y^*} + P e_w C_m^* R_r = 0$
 $\frac{1}{\sqrt{\tau}} \frac{dC^*}{d\eta} + P e_w C_m^* R_r = 0$
 $\frac{dC^*}{d\eta} + P e_w \sqrt{\tau} C_m^* R_r = 0$
 $\text{at } \eta = 0, \left[\frac{dC^*}{d\eta} + AC_m^* R_r = 0 \right]$

So, at eta at y star is equal to infinity C was is equal to C naught; so, C star is equal to 1; so therefore, and at eta is equal to infinity, we have C star is equal to one that is one boundary condition expressed in terms of eta, and the other boundary condition was at eta is equal to 0; at eta is equal to 0, you had d C star d eta plus A C m star Rr is equal to 0; how did you get this equation, this equation can be obtained if you look into the governing equation, the original governing equation at eta equal to 0; that means, let us clarify everything, and let us do a couple of steps derivation; so, that will be then things will be clearer at y star is equal to 0, we had the original boundary condition was as del C star del y star plus P e w C m star Rr equal to 0, that was the original boundary condition; now, if we express del C star del y star in terms of eta, we have already derive that one over root over tau d C star d eta, del C star del y star can be written as one over root over tau d C star d eta, we have already looked into a derivation P e w C m star Rr is equal to 0; so, d C star d eta plus P e w root over tau C m star Rr is equal to 0 and this is the constant A that we have already defined earlier.

So, get this expression, which is the boundary condition at eta is equal to 0; y star equal to 0 means eta equal to 0; so, that is how you obtain this boundary condition in term of eta, and these two boundary conditions this is number 1, this is number 2 that defines the you know solution of the governing equation that whatever we have expressed earlier in terms of eta; so, this gives the derivation of the how you obtain these boundary condition

at eta is equal to 0; now, we will be in a position to get the solution of concentration profile.

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The image shows a handwritten derivation on a blue background. The equations are as follows:

$$\frac{d^2 C^*}{d\eta^2} = - \left(\frac{\eta}{2} + A \right) \frac{dC^*}{d\eta}$$

$$\frac{dC^*}{d\eta} = Z$$

$$\frac{dZ}{d\eta} = - \left(\frac{\eta}{2} + A \right) Z$$

$$\Rightarrow \frac{dZ}{Z} = - \left(\frac{\eta}{2} + A \right) d\eta$$

After 1st integration,

$$\frac{dC^*}{d\eta} = Z = K_1 \exp \left(-\frac{\eta^2}{4} - A\eta \right)$$

$$C^*(\eta) = \int_0^\eta \exp \left(-\frac{\eta^2}{4} - A\eta \right) d\eta + K_2$$

Logos for NPTEL and CET I.I.T. KGP are visible in the bottom left and top right corners of the slide respectively.

So, let us write down the governing equation one more time discuss is $C^* d\eta^2$ minus η by 2 plus $A dC^* d\eta$; now, you consider $dC^* d\eta$ is another variable Z ; so, this becomes $dZ d\eta$ minus η by 2 plus A and Z ; so, this becomes dZ by Z is equal to minus η by 2 plus $A d\eta$; so, the one first integration will give you 1 and Z ; so, in terms of the after first integration, what you get is, you get Z is equal to k_1 exponential minus η square by 4 minus $A\eta$. So, the k_1 is the first constant of integration.

So, what is Z , Z is nothing, but $dC^* d\eta$; so, one more integration will result you the expression of C^* as a function of η , what is the expression of C^* , C^* will be k_1 exponential minus η square by 4 minus $A\eta d\eta$ plus k_2 , and what is this integration, **integration** is basically over the variable η ; so, the minimum value of η is η equal to 0 and so, it this integration basically runs from 0 to η implicitly; now you have the two boundary conditions that we have just discussed in the last slide and we using these two boundary conditions, one can evaluate these two constants of integration.

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$\text{at } \eta = \infty, C^* = 1$
 $C^* = K_1 \int_0^{\infty} \exp\left(-\frac{\eta^2}{4} - A\eta\right) d\eta + K_2$
 $C^* = K_1 I + K_2 \dots (1) \checkmark$
 $\text{at } \eta = 0, \frac{dC^*}{d\eta} + A C^* R_r = 0$
 $Z|_{\eta=0} + A C^* R_r = 0$
 $Z = K_1 \exp\left(-\frac{\eta^2}{4} - A\eta\right)$
 $Z|_{\eta=0} = K_1$
 $C^* = K_1 \int_0^{\eta} \exp\left(-\frac{\eta^2}{4} - A\eta\right) d\eta + K_2$
 $C^*|_{\eta=0} = K_2$
 $K_1 + A K_2 R_r = 0 \dots (2) \checkmark$

Now, let us look into the boundary conditions, and evaluate the constants of integration k_1 and k_2 ; the first constant of first boundary condition was at η is equal to infinity; C^* is equal to 1; that means one you just put this condition in your solution, one becomes $k_1 \int_0^{\infty} \exp\left(-\frac{\eta^2}{4} - A\eta\right) d\eta + k_2$; now, again if you know the value of A then these integral is known to you, because it is a definite integral, and once you know the value of A , because that is a constant it must be assuming some value; and how to evaluate the value that we will see later on; once you know the value of A then you can this definite integral can be evaluated by using any you know Trapezoidal rule or Simpson's rule putting the upper limit as you know some higher numbers may be 10 15 something; and how to decide 10 15, I think we have discussed earlier just put any number and you were getting a result that just change the number, just increase the number arbitrarily let us say from 10 to 12 and change, and look into the result if the result does not change in the decimal third order to fourth decimal place; that means 10 is good enough as infinity for this particular problem.

So, let us indicate this integral as I ; so, this becomes $k_1 I + k_2$, this is equation number 1 then we put the other boundary condition, the other condition was at η is equal to 0 you had $\frac{dC^*}{d\eta} + A C^* R_r = 0$, and what is $\frac{dC^*}{d\eta}$, $\frac{dC^*}{d\eta}$ is nothing but, Z the variable Z you have defined Z evaluated at η is equal to 0 plus $A C^* R_r = 0$; now, let us look into what is Z at η is equal to 0, if you look into the expression of Z , the expression of Z was $k_1 \exp\left(-\frac{\eta^2}{4} - A\eta\right)$

square by 4 minus A eta; so therefore, Z evaluated at eta equal to is nothing but k 1, this exponential term becomes one; so, Z at eta is equal to 0 is nothing but k 1, and what is C m star, if you look into the expression of C star, that was k 1 I k 1 0 to eta exponential minus eta square by 4 minus A eta d eta plus k 2; and what is C m star, C m star is nothing but, C star evaluated at y star equal to 0 and y star equal to 0 is nothing, but eta equal to 0; so, C m star is nothing but C star evaluated at eta equal to 0, and once it is eta equal to 0; that means, this upper limit becomes 0; so, this integral will be then from 0 to 0; so, whole integral will vanish; so, this is nothing but the other constant k 2.

So, once you get that then you can write the boundary condition of at eta equal to 0 at d C star d eta at eta equal to 0 is nothing but k 1; so, this will be k 1 plus A C m star will be nothing but k 2 Rr equal to 0; so, you will be getting the second equation from the second boundary conditions; now, there are two equations and two unknowns k 1 and k 2 to be evaluated; once these two constants evaluated, then you will be getting the full concentration profile within the mass transfer boundary layer; now, in the next slide we will be doing the derivation of k 1 and k 2 from this two equations.

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$$(2) \Rightarrow K_1 = -A K_2 R_r I$$

$$K_1 I + K_2 = 1$$

$$-A K_2 R_r I + K_2 = 1$$

$$\Rightarrow K_2 = \frac{1}{1 - A R_r I}$$

$$K_1 = -\frac{A R_r}{1 - A R_r I}$$

Permeate flux $P_w = ?$ ✓
 Permeate conc. $C_p^* = ?$
 $C_p^* = C_n^* (1 - R_r)$ ✓

So, from equation two we will be getting k 1 is nothing but minus A k 2 Rr; so, you just put into the first equation that will be k 1 I plus k 2 is equal to 1; so, k 1 is minus A k 2 Rr I plus k 2 is equal to 1; so therefore, k 2 will be nothing but, one by one minus A Rr I; so, this is the expression of k 2; and what is k 1, k 1 is nothing but minus of A Rr divided

by one minus $A R_r I$; so, this is the expression of k_1 ; so, using these two constants k_1 and k_2 and knowing the value of a_1 one can evaluate the concentration profile within the mass transfer boundary layer; now, but our aim is to basically evaluate what is the permeate flux, what is non dimensional permeate flux, and what is the permeate concentration C_p^* ; and C_p^* can be obtained if you can evaluate the C_m^* the membrane surface concentration through the relationship is C_p^* in nothing but, C_m^* times one minus real retention; if you know the value of real retention that can be you know there are several ways to evaluate the value of real retention of a particular membrane solute system from various independent set of experiments, that we have discussed already earlier.

So, once you know the C_m^* and $P_e w$ from this expression you can calculate the permeate concentration as well; so, you can predict the system performance at every location of time point; so, we have to evaluate the concentration profile at η is equal to 0 that is the C_m^* , what is the expression of C_m^* now, if you look into the derivation earlier C_m^* was nothing but the constant k_2 .

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$$C_m^* = \frac{1}{1 - A R_r I}$$

Darcy's Law:

$$P_e w(\tau) = B \left(1 - \frac{\Delta \pi}{\Delta P}\right)$$

$$J = L_p (\Delta P - \Delta \pi)$$

$$\frac{J A R_r}{D} = \frac{L_p R_r \Delta P}{B} \left(1 - \frac{\Delta \pi}{\Delta P}\right)$$

$$P_e w(\tau) = \frac{B}{D} \left(1 - \frac{\Delta \pi}{\Delta P}\right)$$

$$\pi = a_1 c + a_2 c^2 + a_3 c^3$$

$$\Delta \pi = \pi_m - \pi_p = a_1 (C_m - C_p) + a_2 (C_m^2 - C_p^2) + a_3 (C_m^3 - C_p^3)$$

$$C_p = C_m (1 - R_r)$$

So, write down the C_m^* is nothing but constant k_2 is equal to $1 / (1 - A R_r I)$; and we have the Darcy's law will be valid as the solvent flux equation through the porous medium for every point of time, and if you remember the non dimensional form of Darcy's law will be $P_e w$ into B into $1 - \Delta \pi / \Delta P$; how you obtained

this equation, this equation obtained by the non dimensional of the Darcy's law; if you forget, I can do a couple of steps and you can rebrush your memory that Darcy's law was nothing but J is equal to $\frac{\Delta P}{\Delta x}$, and you just make it non dimensional $J \frac{d}{L}$ by D , in this case d equivalent will be nothing but the radius of the cell because it is the unstart back cell; so, it is a cylindrical cell with their inner radius is R lets say; so, $L \frac{p}{R}$ by D make it non dimensional; so therefore, take ΔP .

So, take ΔP these becomes $\frac{\Delta \pi}{\Delta P}$; so, what is $J \frac{R}{d}$, this is nothing but $P \frac{e}{w}$ non dimensional $P \frac{e}{w}$ this is valid at every τ , and $L \frac{p}{R} \frac{\Delta P}{D}$ is a constant B and $1 - \frac{\Delta \pi}{\Delta P}$; these are non dimensional parameter that takes care of operating pressure diffusivity solid diffuse solid property operating conditions membrane property and of course, the geometry property R .

So, and what is $\Delta \pi$, if you remember π is nothing but, we had an expression π is equal to nothing but a $1 C$ plus a $2 C^2$ plus a $3 C^3$. So, we have the expression of π as a polynomial of concentration; so, what is $\Delta \pi$, $\Delta \pi$ is nothing but the osmotic pressure difference between the membrane surface and in the upstream and the permeate stream. So, $\Delta \pi$ will be nothing but a $1 C_m - C_p$ plus a $2 C_m^2 - C_p^2$ plus a $3 C_m^3 - C_p^3$; so, know you can omit, you can replace C_p in favor of C_m by doing this exercise, C_p is nothing is but C_m into $1 - \text{real retention}$; if you substitute this the $\Delta \pi$ will become a sole function of membrane surface concentration.

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$$\Delta \pi = a_1 C_m R_r + a_2 C_m^2 [1 - (1 - R_r)^2] + a_3 C_m^3 [1 - (1 - R_r)^3]$$

$$= (a_1 c_0 R_r) C_m + a_2 c_0^2 [1 - (1 - R_r)^2] C_m^2 + a_3 c_0^3 [1 - (1 - R_r)^3] C_m^3$$

$$= b_1 C_m + b_2 C_m^2 + b_3 C_m^3$$

$$b_1 = a_1 c_0 R_r; \quad b_2 = a_2 c_0^2 [1 - (1 - R_r)^2]$$

$$b_3 = a_3 c_0^3 [1 - (1 - R_r)^3]$$

$$\frac{d\Delta \pi}{dC_m} = \left(\frac{b_1}{dC_m}\right) C_m + \left(\frac{b_2}{dC_m}\right) C_m^2 + \left(\frac{b_3}{dC_m}\right) C_m^3$$

So, let us do that and see what we get, what we will be getting is that delta pi will be nothing but, a 1 C m Rr plus a 2 C m square, this becomes 1 minus 1 minus Rr square, this becomes a 3 C m cube 1 minus 1 minus Rr whole cube; and now, make C m non dimensional in favor of the bulk concentration; so, this becomes a 1 C 0 Rr times C m star, this becomes a 2 C 0 square 1 minus 1 minus Rr square times C m star square plus a 3 C 0 cube 1 minus 1 minus Rr whole cube C m star cube, and this becomes b 1 times C m star plus b 2 times C m star square plus b three times C m star cube, and your b 1 will be nothing but a 1 C 0 Rr, your b 2 will be nothing but a 2 C naught square 1 minus 1 minus Rr whole square, and b 3 will be nothing but a 3 C 0 cube 1 minus 1 minus Rr whole cube.

So, these are $\left(\frac{b_i}{dC_m}\right)$ of to express delta pi in terms of C m star only, and 1 by delta pi by delta P will be nothing but, b 1 by delta P plus delta P C m star plus b 2 by delta P C m star square plus b 3 by delta P times C m star cube, and now this becomes a non dimensional number, this becomes a non dimensional number, this becomes a non dimensional number, and all the variables are non dimensional. So, whole things becomes non dimensional, and in a sense your osmotic pressure difference can be expressed in terms of C m star only and that is the way how the derivation goes. So, what you have now, let us write down the governing equations whatever you are the relevant equations in one place that will be better otherwise the whole thing becomes cluttered.

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$$C_m^* = \frac{1}{1 - A R r I} \quad \text{--- (1)}$$

$$A = P e w \sqrt{\tau} \quad \text{--- (2)}$$

$$P e w(\tau) = B \left(1 - \frac{\Delta P}{\Delta P}\right) \quad \text{--- (3)}$$

Iterative Algorithm

at a particular op. time τ

- (i) guess value of C_m^*
- (ii) $P e w$ from Eq. (3)
- (iii) Evaluate A from Eq. (2)
- (iv) Evaluate I , $I = \int_0^{\infty} \exp\left(-\frac{r^2}{q} - A \eta\right) d\eta$
- (v) Calculate C_m^* from Eq. (1)
- (vi) Iterate until $|C_m^{*n} - C_m^{*(n-1)}| < \epsilon$

So, the relevant equations are first equation is C_m^* as one minus $A R r I$ and what is A becomes A is nothing but, if you remember A is nothing but $P e w$ root over τ ; and what is the Darcy's law, Darcy's law is nothing but $P e w$, at any point of time becomes B times 1 minus Δp by ΔP and this is a function, sole function of C_m^* .

So, that's it. Now, using an appropriate algorithm one can do an iterative calculation, and can get the permeate flux and permeate concentration at every time point, and you'll be getting a profile as a function of time. So, we use the following iterative algorithm, what is the iterative algorithm, the at a particular operating time, operating time let us say 5 second 10 second or whatever particular time τ you do the following, step number 1: guess the value of C_m^* , once you guess the value of C_m^* we can calculate $P e w$ from equation number 3, what is equation 3 from here, this is let us say 3 2 1.

So, you know the value of C_m^* , because you have already guessed it; and this is a sole function of C_m^* ; ΔP is the operating pressure operating condition that is known to because you are setting it; B is a parameter that is all the values are known to you membrane radius, cell radius and so for. So, you can get the value of $P e w$, once you know the value of $P e w$ from step number 3 is from equation 2 you know the time point, you know the value τ , you know the $P e w$; so, you can get the value of A . So, evaluate A from equation 2; **once you do that then you can you are in a position** once you know the value of A you are in a position to evaluate the definite integral I ; evaluate I from its

definition, I is equal to 0 infinity exponential minus eta square by four minus A eta d eta. So, you know the value of A substitute there; so, everything is known eta is a dummy variable; so, just evaluate the definite integral I; so, once you know the value of I then you can go back to equation number 1, because I is now known to you, A is now known to you, Rr I already determined. So, you can calculate C m star from equation 1, and check whether this calculated value is close to that guess value or not.

So, next one is just iterate, on the absolute value of C m star guess minus C m star calculated is less than epsilon or not; if not then you guess a new value of C m star likewise this loop has to be iterated for a particular time point and what you do next **next** you have an increment of time tau that is a tau plus delta and re iterate this equation this **this** steps, but in the second step in the second step you should have a judicial guess of C m star the judicial guess is basically the converged value of C m star at the earlier time point then it will be the **the the** iteration of the convergence will be faster for the next time point likewise what in the likewise it will be continued it will be continued and what is the final outcome of this exercise final outcome is that you will be getting various a profile of C m star as a function of tau you will **will** be getting a profile of Pe w as a function of tau.

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Predict system performance as a function of time.

Mass Transfer Coefficient
(K)

Definition of mass transfer Coeff

$$K (C_m - C_o) = -D \left(\frac{\partial C}{\partial y} \right)_{y=0}$$

$\frac{\partial C}{\partial y} \rightarrow \frac{dC^*}{d\eta}$

NPTEL

So, out of it you'll be getting Pe w as a function of tau and this is this is Pe w and this is C m star as a function of tau once you know the C m star is a function of tau you can get a

value of C_{Pstar} as a function of τ through the relationship C_{Pstar} is equal to $C_{m star}$ into one minus real retention. So, the this way one can get the one can predict the **performance** system performance as a function of time in an unstart batch cell experiment and why you are conducting the unstart batch cell experiment that I think you have a discussed that issue earlier and **that** that is the way.

Now, one maybe being a chemical engineering maybe interested what is expression of a mass transfer coefficient in this case now **now** this mass transfer coefficient k is a vital parameter as far as the design of any a chemical engineering system is concerned where the mass transfer is involved. So, we start with the definition of mass transfer coefficient k times C_m ; that means, C at y equal to 0 minus C_{naught} is equal to minus $d C d y$ del C del y at y is equal to 0 that is the definition of the mass transfer coefficient.

Now, what do we do I think we have done this exercise earlier in the case of rectangular geometry we have to replace $del C del y$ in terms of $d C star d eta$ in terms of non dimensional similarity parameter eta and before **before** that you have to make these equation non dimensional. So, we **we** first we make then make the equation non dimensional then we put the value of the governing **equation** the **the the** derivative in terms of eta .

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Handwritten derivation for the Sherwood number (Sh):

$$k(C_m^* - C_0) = -D \left(\frac{\partial C}{\partial y} \right)_{y=0}$$

$$k(C_m^* - 1) = -D \left(\frac{\partial C^*}{\partial y} \right)_{y=0}$$

$$y^* = y/R$$

$$k(C_m^* - 1) = -\frac{D}{R} \left(\frac{\partial C^*}{\partial y^*} \right)_{y^*=0}$$

$$\Rightarrow \frac{kR}{D} = \frac{-\left(\frac{\partial C^*}{\partial y^*} \right)_{y^*=0}}{(C_m^* - 1)}$$

$$Sh = \frac{-\frac{\partial C^*}{\partial y^*} |_{y^*=0}}{C_m^* - 1}$$

$$Sh = \frac{-k_1/\sqrt{\pi}}{(C_m^* - 1)} = \frac{-k_1/\sqrt{\pi}}{k_2 - 1}$$

So, $k C_m star$ minus C_0 is equal to minus $d C star del y$ del $C del y$ this there is no star here at y is equal to 0 now you divide both **both** side by C_0 . So, this become C_m

star C_m star minus 1 is equal to minus $d \frac{dC}{C} / 0$. So, it become $\frac{dC}{dy}$ at y is equal to 1.

Make y star as y by R . So, y is nothing, but r times y star y star. So, $k C_m$ star minus one $d \frac{dC}{dy}$ star now this becomes y star is equal to 0 and what is this kR by D is nothing, but minus $\frac{dC}{dy}$ star at y star is equal to 0 divided by C_m star minus one and this quantity kR by D is nothing, but the sherwood number now there is a neat expression of sherwood number in terms of non dimensional parameter and if you now put the expression of $\frac{dC}{dy}$ star in terms of η and this becomes sherwood is equal to minus 1 over root over $\tau d C$ star $d\eta$ at η is equal to zero divided by C_m star minus 1.

Now, we have done this exercise earlier just few minutes back that $d C$ star $d\eta$ $d\eta$ equal to 0 and C_m star in terms of the integration constants and this becomes minus k_1 by root over τ divided by k_2 minus 1 just just put the value of $d C$ star $d\eta$ eta equal 0 we have done that probably few minutes back and C_m star is nothing, but k_2 . So, it becomes k_2 minus 1 now you put the expression of k_1 and k_2 because you have evaluated the k_1 and k_2 in terms of ARr and I and substitute there and see what you get what you get what you will be getting is like that sherwood is nothing, but one over I times root over τ .

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$$Sh = \frac{1}{I \sqrt{\tau}}$$

$$Sh = Sh(\underline{\tau}, \underline{Pe_w})$$

We have discussed
Osmotic Pressure Controlled
Filtration.

- 1) 1-Dim modeling.
- 2) Shortcomings of 1Dim model
- 3) 2 Dimensional mode \rightarrow Rect. batch.

NPTEL

So, this is the expression you will get at the end of it now in this expression. So, so in a sense Sherwood number will be a function of two things that is observed here one is time another is **if** if you remember in the integration of I there is a value there is a constant a and a is nothing, but Pe_w times root over τ . So, this will be function of Pe_w as well. So, the porosity the **the** permeate flux this permeate flux is will **occur** will occur because of the porous membrane present at the wall. So, porosity over of wall will be taking by this parameter Pe_w and the time variation is taken care **care** by the non dimensional quantity τ .

So, therefore, Sherwood will be and if you if you look into **Sherwood** Sherwood will be a function of both permission **permission** as well as the time of operation and as it is evident from this expression the Sherwood is inversely proportional to time; that means, as time progresses the mass transfer coefficient the Sherwood number will go on decreasing; that means, mass transfer coefficient will go on **will** will **will** go on decreasing because mass transfer coefficient is directly proportional to Sherwood number there is a non dimensional version of mass transfer coefficient.

So, therefore, as time progresses the concentration boundary layer develops on the membrane surface it offers physically it means it offers more resistance against the solvent flux and **and and** the permeate **flux** and the mass transfer amount of mass transfer will be reduced since the amount of mass transfer will be reduced by **by** the mass transfer coefficient and the permeate flux will be connected by the boundary condition at y equal to 0 if you remember. So, therefore, the permeate flux will also decrease as time progresses.

So, these maybe of interest to the chemical engineering engineer **student people** a chemical engineers to design efficiently a system. So, that goes the modeling part on calculation **part/part os** part of the osmotic pressure controlled unstart batch cell.

So, let us try to summarize whatever we have done till now we have discussed osmotic pressure controlled filtration. So, we have discussed one dimensional one dimensional modeling we have discussed the short coming of one dimensional models then we discussed the **the the** two dimensional model in both geometries a a rectangular cross flow cell there is a steady state process and in a batch cell there is an unsteady state

process. So, both steady state and transient operations we have seen the modeling how to model the system.

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Solutes of low MW

$$\pi = \frac{RTC}{M}$$

Salt NaCl \rightarrow 58.5
Sucrose \rightarrow 342
dextrose \rightarrow 92

RO, NF, UF of protein
Solute MW $<$ 70,000

Osmotic Pr. model.

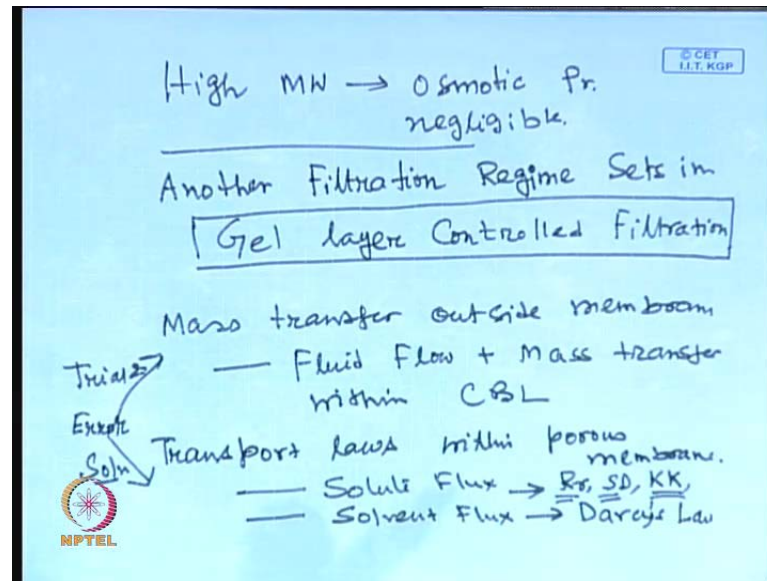
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Now, let us discuss when this osmotic pressure controlled filtration becomes very important the osmotic pressure controlled filtration becomes very important when you are talking about filtration of solutes which will be having very high osmotic pressure solutes of low molecular weight solutes will not be having the osmotic pressure the solvent the solution will be having the osmotic pressure, but the solvent must be containing the solutes of low molecular because if you look into the definition of osmotic pressure the first coefficient is RTC by M . So, if you go for lower molecular solutes lower molecular typically you can you can correlate with the size that you **you** can talk about the smaller sized particles if you have solutes of smaller molecular less the osmotic pressure becomes dominant.

For example salt NaCl its is fifty eight point five sucrose three hundred forty two dextrose 92 I guess 92 or 42 or something like that.

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So, you will be having if you if you have the osmotic pressure in the range of around the molecular in the range of 1000 the osmotic pressure can be very important for the charged neutral solutes for the **for the** charged proteins which you will be having a high **osmotic** molecular for example, 70000,80000 but because of the charged **charged** effect and other interactions the osmotic pressure can be very high.

So, if you are talking of untrust solutes for example, polysaccharides dextran or things like that are very high molecular rate the osmotic pressure the denominator becomes very large in that case osmotic pressure will be negligible for most of the cases of reverse osmosis nano filtration and ultra filtration of proteins and solutes having molecular rate less than 20000 the osmotic pressure module becomes very important it can be even **even** for the putting which will be having a molecular at seventy thousand the osmotic pressure is very **very very** important because of the chartered effect and the other **other** effects.

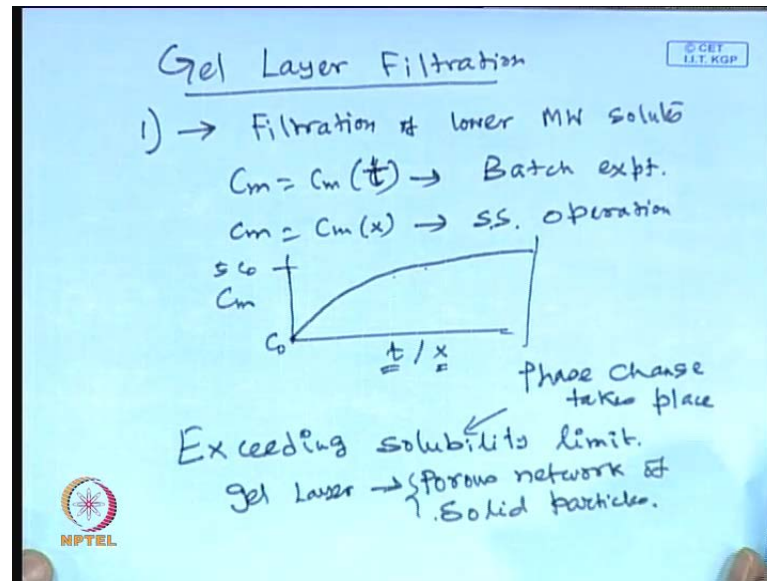
So, you will we talk about the osmotic pressure model in those cases. So, if you remember starting from micro filtration to a nano **filtration** to reverse osmosis most of the operating conditions or the membrane separations they remain under the osmotic pressure controlled regime now for the solutes which have a very high molecular weight the osmotic pressure becomes negligible another filtration regime will set in that is known as the gel filtration or gel layer controlled filtration now **jis/this** gel filtration is

what is gel filtration gel filtration is gel filtration can occur now before going into the gel layer filtration we can we can discuss the various you know aspects of osmotic pressure model that we have already discussed that we have seen the various osmotic pressure models for an and we have identified that the mass transfer outside the mass outside the membrane; that means, within the mass transfer boundary layer that has to be coupled or hooked up by the mass transfer or the **the** solute transport or solvent transport within the porous membrane and mass transfer outside boundary layer sorry outside membrane is solved by fluid flow plus mass transfer within concentration boundary layer that has to be coupled by the transport laws within porous membrane and these transport laws are typically the solute flux equation and solvent flux equation.

Now, this solvent flux is typically and generally and for all the cases is calculated by or **or** quantified by darcy's law on the other hand the solute flux is given by either a definition of real retention or by the **the** solution diffusion model or kevin kachinski model modified models either that definition of real retention or a solution diffusion model that we have already discussed and modified version of the solution diffusion model that is the quedan kochalski model these are the various models that you are talked about the solute flux. So, by **by by** getting some **definetion** the appropriate definition of the solute flux and solvent flux these two has to have up to be coupled and you can get the system performance prediction and the **the** typical feature is you'll be getting a trial and error solution the solution is always trial and error.

Now, if we let us go to the gel layer control filtration because it is another regime in the whole filtration system it is more rampant in case of micro filtration and in **in** at the upper **upper** range of alter filtration cases where the forces are much bigger.

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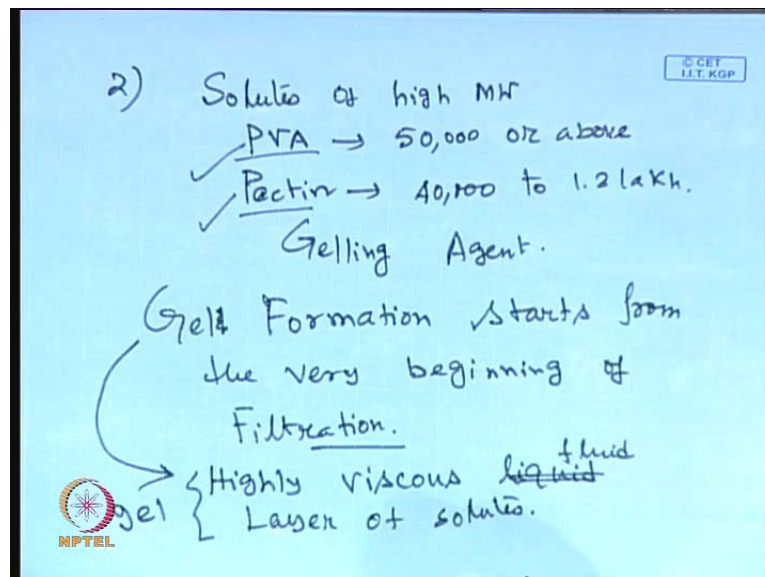
now there are two ways when gel layer filtration can occur in the case in case number one if you talk about any osmotic pressure controlled model basically you **you** are talking about filtration of lower molecular weight solutes. So, initially; that means, if you're talking about a batch cell operation or transient operation initially the filtration is osmotic pressure controlled if you talk about a steady state case for a for **for** initial length of the channel the ultra filtration is osmotic pressure controlled now as time progresses or the length of you know channel it goes on if you go **down** as you go downstream of the channel the membrane surface concentration increases; that means, C_m is a function of t or t or time of operation in case of batch experiment and C_m is a function of X in case of steady state operation whereas, X is the length of the module now and in both the cases you have seen that the minimum value of membrane surface concentration is C_0 and it will be always increasing beyond C_0 and it keeps on increasing.

Now, if you plot C_m versus t or x the curve will be typically something like this starting from C_0 to let us say five times C_0 or whatever now what happens if you talking about such a system which is basically an osmotic pressure controlled initially the membrane surface concentration increases with time for a batch scale batch **batch** system or as a **function** increases with respect to X in a steady state cross flow system and it may happen at the particular temperature it exceeds the solubility limit if it exits the solvability limit at that point onwards the whole system the **the** solute particles will be deposited as a solute particles over the membrane surface because the phase

transformation will occur. So, there may be a phase change takes place; that means, you are exceeding the solubility limit in that case the formation of solute solute particle will **will will** be will lead to formation of a porous network of solute particles over the membrane surface that is known as the gel layer.

So, that leads to the formation of the gel layer, but since it is porous you can still expect some amount of flux solvent flux out of it some amount of g out of it, but that will be typically lower and in this case we would like to avoid that gel layer formation the second way of.

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So, that is **that is that is** one way of formation of gel layer and most common one is that first some solutes it is not osmotic pressure from the very beginning in that case we are talking about solutes of very high molecular rate if you're talking about filtration of solutes of very high molecular rate for example, polyvinyl alcohol which will be having a molecular rate around fifty thousand or above pectin **pectin** is a polysaccharides it is basically a you know it will be having it is a it is a polydispersed material; that means, it'll be having a molecular rate distribution let us say from forty thousand to 1 point 2 lakh it is a it is a polydisperse.

Now, things like the solutes like polyvinyl alcohol pectin they are well known gelling agent pectin is a polyvinyl alcohol is a polymer pectin is basically a polysaccharide that is available in any fruit juice and you would like to like to remove pectin from a citrus

fruit juice because pectin forms a an or complex with protein and it destroys the fruit juice. So, one has to remove the pectin from the fruit juice for a clarified juice which will be having a high shelf life.

So, maybe alter filtration can be used for it. So, if you are talking about the solutes like polyvinyl alcohol pectin or things like that which are well known gelling agents the gel formation occurs from the very beginning of the filtration gel formation starts from the very beginning of filtration. So, there is no nothing like osmotic pressure controller and then the regime changes over to gel layer the gel formation starts from the very beginning of a filtration and you cannot avoid that what you can do you can you can. So, I said in the other earlier mechanism we said that we can you can avoid the gel layer formation.

So, you can set your operating convenience such that that formation can be avoided, but in this case gel layer formation cannot be avoided because, but **but** the operation has to be done because of the juice has to be clarified . So, therefore, one has to be very careful and set the tuning parameters of the operating parameters such that the gel layer formation becomes minimum and what is a definition of a gel layer it is not a solute you know layer it becomes it is a basically highly viscous liquid layer I should not say liquid it is a highly viscous fluid layer it flows stills it flows highly viscous fluid layer of solutes again it is porous it forms over the membrane surface and it is gel we **we** term it a gel in this particular case. So, in that we I stop here in the next class what we will do we will look into the modeling and the calculation of the gel layer control filtration thank you.