

Thermodynamics of Fluid Phase Equilibria
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Lecture – 12
Thermodynamic Calculus – 3

In the last lecture we were trying to cover up. We looked at multivariate calculus relevant for a thermodynamic such as in this course we will continue this exercise.

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The image shows a whiteboard with the following handwritten equations:

$$f = f_1(x, y) \qquad f = f(x, y, z)$$

$$y = f_2(x, z)$$

$$f = f_1(x, f_2(x, z)) = f_3(x, z)$$

$$\left. \frac{\partial f}{\partial x} \right|_y = \frac{\partial f_1}{\partial x}(x, y)$$

$$\left. \frac{\partial f}{\partial x} \right|_z = \frac{\partial f_3}{\partial x}(x, z)$$

$$\left. \frac{\partial f}{\partial x} \right|_y \neq \left. \frac{\partial f}{\partial x} \right|_z$$

So, let us assume functions f is a function of x, y , let us say f_1 , where y is dependent on x and z ok. So, in other word f is function of f, x, f_2, x, z or some function f_3, x, z ok.

So, the reason we are considering such a complicated function, because typically thermodynamic functions are not just f is equal to x and y , where x and y are independent. So, most of the time we have more complicated functions. Now in this case if you want to find out partial derivatives such as $\frac{\partial f}{\partial x}$, it has no meaning until you specify what is fixed for example, $\frac{\partial f}{\partial x}$ at constant y is $\frac{\partial f_1}{\partial x}(x, y)$ or $\frac{\partial f}{\partial x}$ at constant z is $\frac{\partial f_3}{\partial x}(x, z)$ because this is s and z we are going to consider f_3 exit. So, clearly in this case $\frac{\partial f}{\partial x}|_y$ is not equal to $\frac{\partial f}{\partial x}|_z$. So, what is the relation between this 2 ok?

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The image shows a whiteboard with handwritten mathematical equations. At the top, it states $\left. \frac{\partial f}{\partial x} \right|_y \neq \left. \frac{\partial f}{\partial x} \right|_z$. Below this, the differential df is expressed as $df = \left. \frac{\partial f}{\partial x} \right|_z dx + \left. \frac{\partial f}{\partial z} \right|_x dz$. This is followed by an equation: $\Rightarrow \left. \frac{\partial f}{\partial x} \right|_y = \left. \frac{\partial f}{\partial x} \right|_z \frac{\partial x}{\partial x} \Big|_y + \left. \frac{\partial f}{\partial z} \right|_x \frac{\partial z}{\partial x} \Big|_y$. The final result is boxed: $\left. \frac{\partial f}{\partial x} \right|_y = \left. \frac{\partial f}{\partial x} \right|_z + \left. \frac{\partial f}{\partial z} \right|_x \frac{\partial z}{\partial x} \Big|_y$. The whiteboard interface includes a toolbar at the top and a page number '16' at the bottom.

So, we are going to consider the differential of f can be written as $\left. \frac{\partial f}{\partial x} \right|_z dx + \left. \frac{\partial f}{\partial z} \right|_x dz$ at constant z multiplied by dx plus $\left. \frac{\partial f}{\partial z} \right|_x dz$ at constant x ok. Now this can be further differentiated such that $\left. \frac{\partial f}{\partial x} \right|_y$ is $\left. \frac{\partial f}{\partial x} \right|_z$ and this is written as $\left. \frac{\partial f}{\partial x} \right|_z + \left. \frac{\partial f}{\partial z} \right|_x \frac{\partial z}{\partial x} \Big|_y$ plus ok. So, this is of course, going to be 1 and this you can write $\left. \frac{\partial f}{\partial x} \right|_z$ at constant z , plus $\left. \frac{\partial f}{\partial z} \right|_x \frac{\partial z}{\partial x} \Big|_y$ at constant x ok. So, this is a relation between these 2 partial derivatives ok.

So, we can further continue particularly now considering f as a function of more variables, something which is what we are going to encounter in thermodynamics.

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$$\text{If } f = f(x_1, x_2, \dots, x_n)$$
$$df = \sum_{i=1}^n \left. \frac{\partial f}{\partial x_i} \right|_{x_1, x_2, \dots, x_{i-1}, x_{i+1}, \dots, x_n} dx_i$$

$$f(x_1, x_2, \dots, x_n) - f(x_1^0, x_2^0, \dots, x_n^0)$$
$$= \sum_{i=1}^n \int_{\Gamma} \left. \frac{\partial f}{\partial x_i} \right|_{x_j \neq x_i} dx_i$$

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So, let us assume that f is a function of x_1, x_2 and so forth. In that case df can be written as summation of partial derivatives of f with respect to x_i , where i is equal to 1 to n and this derivative is at keeping $x_1, x_2, x_{i-1}, x_{i+1}$ till x_n constant. So, this is something which we can write the differentiation or differential amount of f in terms of the summation of the partial derivatives.

We can integrate this also all right. So, if you integrate this from x_1^0 's to x_1 , then we can get something like this minus f of x_1^0, x_2^0 this is nothing, but summation of i equal to 1 to n , integral across the path we want to take $\frac{\partial f}{\partial x_i} dx_i$, x_0 is not equal to $x_i dx_i$. Now remember that the left hand side is independent of the path, but this right hand side the expressions will depend on which path you want to take. So, in order to illustrate we will take some example.

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$$f(x_1, x_2, \dots, x_n) - f(x_0, x_2, \dots, x_n)$$

$$= \sum_{i=1}^n \int_{\Gamma} \left. \frac{\partial f}{\partial x_i} \right|_{x_j \neq x_i} dx_i$$

composed of straight line segments

$$\Delta f = f(x_1, y_1) - f(x_0, y_0)$$

$$= \int_{\Gamma} \left. \frac{\partial f}{\partial x} \right|_y dx + \int_{\Gamma} \left. \frac{\partial f}{\partial y} \right|_x dy$$

So, let us consider plot y against x where we have x_0, y_0 and somewhere here is x_1, y_1 . So, the easiest path we can consider is basically composed of straight line segments ok. So, what we can do is, we can consider straight lines such as this ok. This is one path composed of 2 line segments, let me say this is Γ_a and there is another one from here till here. So, in this case we are keeping x constant in this case we are keeping y constant. This would be another side path composed of 2 segments and let me say this is Γ_b .

Ok and. So, what we are interested in is Δf ok. So, Δf is in this case since its a function of 2 variables. So, this is going to be $f(x_1, y_1)$ minus $f(x_0, y_0)$ ok. So, based on this definition which we have put it here, this Δf you can be written as $\frac{\partial f}{\partial x} dx + \frac{\partial f}{\partial y} dy$ ok.

So, now since we can consider 2 different path we are going to derive this expression for 2 different path, and you can clearly see that the expressions are going to be different ok. So, let me consider this a all right.

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The image shows a whiteboard with handwritten mathematical notes. At the top, there is an equation:
$$= \int_{\Gamma} \frac{\partial f}{\partial x} dy + \int_{\Gamma} \frac{\partial f}{\partial y} dx$$
 Below this, two paths are defined:
$$\Gamma_a \equiv (x_0, y_0) \rightarrow (x_1, y_0) \rightarrow (x_1, y_1)$$

$$\Gamma_b \equiv (x_0, y_0) \rightarrow (x_0, y_1) \rightarrow (x_1, y_1)$$
 The main derivation is for Γ_a :
$$\Delta f = \int_{(x_0, y_0)}^{(x_1, y_0)} \frac{\partial f}{\partial x} \Big|_{y=y_0} dx + \int_{(x_1, y_0)}^{(x_1, y_1)} \frac{\partial f}{\partial y} \Big|_{x=x_1} dy$$
 The first integral is annotated with a horizontal line segment from (x_0, y_0) to (x_1, y_0) and the second with a vertical line segment from (x_1, y_0) to (x_1, y_1) . There are some additional scribbles and a small '18/68' at the bottom right of the whiteboard.

So, this will be x_0, y_0 to x_1, y_0 and finally, x_1, y_0 to x_1, y_1 and this is going to be x_0, y_0 to x_0, y_1 to x_1, y_1 . So, just to make this clear we this is going to be x_0, y_0 and this is going to be x_1, y_0 , and this is going to be x_1, y_1 ok. So, let me just further write it now my expression now this is going to be for Γ_a a path $\frac{\partial f}{\partial x}$ in this case y_0 is constant. So, this is nothing, but y equal to y_0 dx plus. So, this is part 1 and similarly you can also integrate x_0, y_0 to x_1, y_1 sorry y_0 , $\frac{\partial f}{\partial y}$, x device. So, this is going from x_0, y_0 to x_1, y_0 plus. So, this is again plus here, plus you have another path another line segment from x_1, y_0 to x_1, y_1 .

So, its going to be x_1, y_0 to x_1, y_1 $\frac{\partial f}{\partial y}$. So, this is basically again y is equal to y dx , since x_1 is fixed and y is variable which means this term is going to be 0 considering that x is fixed. And the second term is going to be again x_1, y_0 to x_1, y_1 $\frac{\partial f}{\partial y}$ dy where x is nothing, but x_1 ok. So, what I did was I actually you have 2 parts or 2 partial derivatives. So, this is basically the partial derivative integration over first line segment, this is the derivation over the line second line segment. In the first line segment y is equal to y is constant. So, the first derivative is this one and the second derivative this one since x is constant this should be also 0 assimilate this for the second side line segment.

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$\Gamma_b \equiv (x_0, y_0) \rightarrow (x_1, y_1)$

$\Gamma_a:$

$$\Delta f = \int_{(x_0, y_0)}^{(x_1, y_0)} \frac{\partial f}{\partial x} \Big|_{y=y_0} dx + \int_{(x_1, y_0)}^{(x_1, y_1)} \frac{\partial f}{\partial y} \Big|_{x=x_1} dy$$

$$+ \int_{(x_0, y_0)}^{(x_1, y_0)} \frac{\partial f}{\partial y} \Big|_{x=x_0} dy + \int_{(x_1, y_0)}^{(x_1, y_1)} \frac{\partial f}{\partial x} \Big|_{y=y_0} dx$$

$$\Delta f = \int_{x_0}^{x_1} \frac{\partial f}{\partial x} \Big|_{y=y_0} dx + \int_{y_0}^{y_1} \frac{\partial f}{\partial y} \Big|_{x=x_1} dy$$

$$\Gamma_b \Delta f = \int_{x_0}^{x_1} \frac{\partial f}{\partial x} \Big|_{y=y_1} dx + \int_{y_0}^{y_1} \frac{\partial f}{\partial y} \Big|_{x=x_0} dy$$

So, when you simplify this expression delta f turns out to be x_0, x_1 since y is constant here, I can simply write $\frac{\partial f}{\partial x}$ by $\frac{\partial f}{\partial x}$, y is equal to y_0 dx, plus this one which is going to be since x_1 is constant I can simply write y_0 to y_1 , $\frac{\partial f}{\partial y}$ by $\frac{\partial f}{\partial y}$, x is equal to x_1 dy ok. So, this is one expression, now if you had used Γ_b you could show that $\frac{\partial f}{\partial x}$ is x_0, x_1 $\frac{\partial f}{\partial x}$ by $\frac{\partial f}{\partial x}$, y is equal to y_1 dx plus y_0, y_1 $\frac{\partial f}{\partial y}$ by $\frac{\partial f}{\partial y}$, x is equal to x_0 dy ok.

So, the constants are different ok. So, this is a simple way of finding out changes in the functions using paths, which are broken into different line segments ok. So, we are going to continue this exercise now, and now we are going to look at Euler theorem for homogeneous function ok.

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Euler's theorem for homogeneous fn.

Thermo fn & derivative - extensive
or intensive

$$f(\lambda x) = \lambda f(x)$$

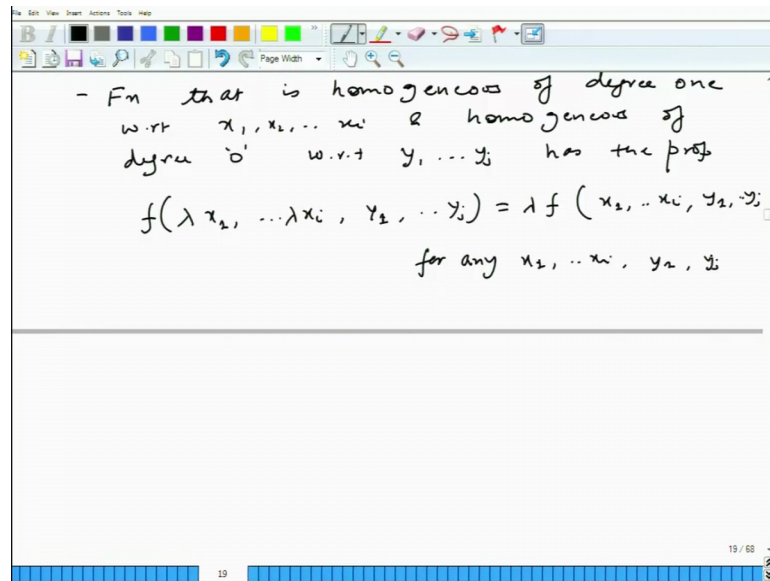
are called homogeneous
of degree one

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Now this is the important component of classical thermodynamics. So, let us make a comment here, the thermodynamic function and its derivative are either extensive or intensive. So, this is something which we observe that all these functions are either extensive or intensive that we have already defined. Now mathematically a function which satisfies a relation such as this, for all values of x are called homogeneous function of degree 1 ok.

So, this is a typical extensive property of which we know from thermodynamics ok. So, the intensive property would be a homogenous function of degree 0 ok.

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So; that means, the function that is homogeneous of degree 1 with respect to variable let us say x_1, x_2, x_i and homogeneous of degree 0 with respect to variables y_1 till y_j has the following property, has the property that you can write the following as this ok. So, what I did basically nothing, but the lambda which is multiplied here in this variables, have been taken out I multiplied the whole function by a lambda that is what basically nothing, but Euler theorem ok.

So, this is true for any x_1, x_i, y_1, y_j ok. Now thermodynamic function a homogeneous function of degree 1 with respect to extensive variable and degree 0 with respect to intensive variable. So, this Euler theorem basically provides a link between its function is derivative, and we are going to show that we are going to prove that ok.

So, let us consider again this function, which we know f and just take it its derivative.

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$$\frac{\partial f(\lambda x_1 \dots \lambda x_i, y_1 \dots y_j)}{\partial \lambda} = \frac{\partial \lambda f(x_1, \dots, y_j)}{\partial \lambda} = f(x_1, \dots, y_j)$$

$$\frac{\partial f(\lambda x_1 \dots)}{\partial (\lambda x_1)} \frac{\partial (\lambda x_1)}{\partial \lambda} + \dots + \frac{\partial f(\lambda x_1 \dots \lambda x_i \dots)}{\partial (\lambda x_i)} \frac{\partial (\lambda x_i)}{\partial \lambda}$$

$$= x_1 \frac{\partial f(x_1, \dots, y_1 \dots y_j)}{\partial x_1} + \dots + x_i \frac{\partial f(\dots)}{\partial x_i}$$

Del f lambda x 1 till lambda x i, and y 1 till x j y until y j with respect to lambda. So, we are taking a derivative of f with respect to lambda ok. Knowing this a Euler theorem, we can write this expression as del lambda f x 1 till y 1 and so forth ok.

Now, this is with respect to lambda, thus this is nothing, but f of x 1, x i comma y 1 till y j all right. So, this derivative of f with respect to lambda is nothing, but simple function f. Now you can also consider also rewrite this expression as follows this is nothing, but based on the differential information this is nothing, but del f lambda x 1 del lambda x 1 del lambda x 1 del lambda plus and so on until you have this del f lambda x 1 lambda x i del lambda x i del lambda x i, divided by del lambda. So, this is nothing, but x i ok.

So, this turns out to be nothing, but x 1 ok. So, what we have is and similarly that 1 is for x i, now this is just a multiplication of lambda lambda can be can come out here and similarly f can be written as f of this can be written as lambda times f and lambda lambda gets cancel. So, this is nothing, but equivalent of saying x 1 del f, x 1 x i y 1, y j divided by del x 1 plus x i del f again there is del x i and this is nothing, but f from this term. So, this gives us basically Euler theorem for homogeneous function in this form ok.

So, this is basically because x 1 to x i are homogeneous function of degree 1. Now having derived this expression we can also apply it to thermodynamic functions. So, let me just try that for internal energy. So, if you consider.

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$$f = x_1 \frac{\partial f(x_1, \dots, x_i, y_1, \dots, y_n)}{\partial x_1} + \dots + x_i \frac{\partial f(\dots)}{\partial x_i}$$

Apply to $U(S, V, N_1, N_2, \dots, N_m)$

$$U = S \left. \frac{\partial U}{\partial S} \right|_{V, N_i} + V \left. \left(\frac{\partial U}{\partial V} \right)_{S, N_i} + \sum_{i=1}^m N_i \left. \left(\frac{\partial U}{\partial N_i} \right)_{S, V, N_{j \neq i}} \right.$$

So, the objective now is to apply to let us say u as a function of u, v and 1 and 2 and so forth ok. Now these are all our extensive variable. So, this is all homogeneous function of degree 1 ok. So, with respect to all the variable. So, in that case you can be written as u is nothing, but f can be written as x_1 is $s \frac{\partial f}{\partial s}$ keeping all of the variable constant v plus $v \frac{\partial u}{\partial v}$ by $\frac{\partial u}{\partial v}$ keeping all variable constant plus summation $N_i \frac{\partial u}{\partial N_i}$ by $\frac{\partial u}{\partial N_i}$ keeping all available all other variable constant.

Ok. So, now, oh sorry this f should have been u . So, we want to write it here u ok.

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$$f = x_1 \frac{\partial f(x_1, \dots, x_i, y_1, \dots, y_n)}{\partial x_1} + \dots + x_i \frac{\partial f(\dots)}{\partial x_i}$$

Apply to $U(S, V, N_1, N_2, \dots, N_m)$

$$U = S \left. \frac{\partial U}{\partial S} \right|_{V, N_i} + V \left. \left(\frac{\partial U}{\partial V} \right)_{S, N_i} + \sum_{i=1}^m N_i \left. \left(\frac{\partial U}{\partial N_i} \right)_{S, V, N_{j \neq i}} \right.$$

$$U = TS - PV + \sum_{i=1}^m \mu_i N_i$$

Euler Integrated FE

Now, what is the partial derivative of u with respect to s is nothing, but T , what is a partial derivative with respect to U , partial derivative with respect to V ; this is nothing, but if you look at the first law of thermodynamics, you can derive that this is nothing, but minus p and this is nothing, but μ_i .

So, we have now relation U is equal to TS minus PV plus summation N_i or $\mu_i N_i$. So, we have now got an equation which is Euler integrated fundamental equation ok. So, this is based on the very simple mathematics and this expression are going to be extremely useful to derive something called gives you an equation which you are going to talk in the next lecture. So, with that I am going to stop and we will see you in the next lecture.