

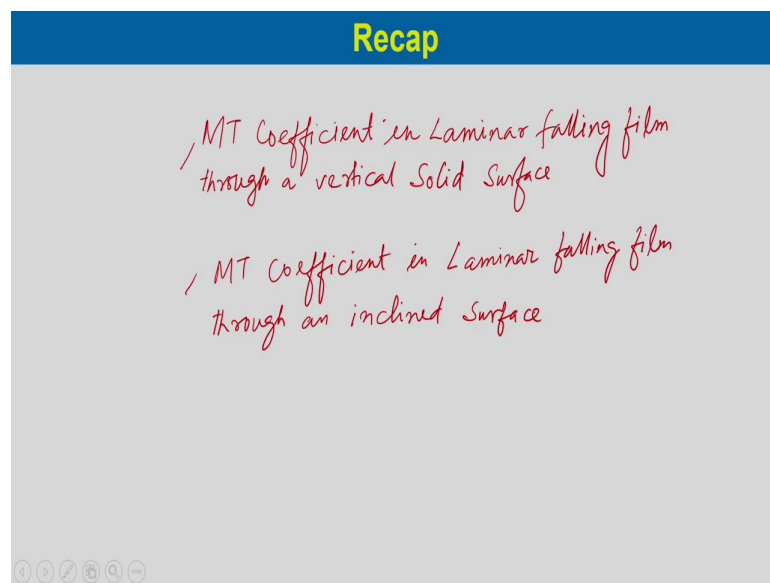
Mass Transfer Coefficients
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Lecture – 12

Boundary Layer Theory and mass transfer coefficients in turbulent flow

Welcome to the 4th lecture of module 2 on mass transfer operation. In module 2 we are discussing the Mass Transfer Coefficient. So, before going to this lecture let us have small recap on our earlier lecture. In our last lecture we discussed, no mass transfer coefficient in laminar flow. In the laminar flow conditions we have considered two cases; one case is mass transfer coefficient in laminar falling film through a vertical solid surface.

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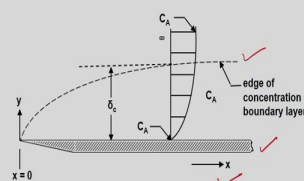


The second case we have considered mass transfer coefficient in laminar flow condition falling film through an inclined surface.

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Boundary Layer Theory In Mass Transfer

- An exact solution can be obtained for the hydrodynamic boundary layer for isothermal laminar flow past a plate
- An extension of the Blasius solution can be extended to derive an expression for convective heat transfer
- In an analogous manner we use the Blasius solution for convective mass transfer for the same geometry and laminar flow.



The diagram shows a flat plate along the x-axis starting at x=0. A velocity boundary layer is shown as a shaded region above the plate, with a vertical arrow labeled δ_v indicating its thickness. A concentration boundary layer is shown as a shaded region above the plate, with a vertical arrow labeled δ_c indicating its thickness. The concentration profile is shown as a curve starting at C_A at the plate and approaching a free-stream concentration $C_A = \infty$. The edge of the concentration boundary layer is marked with a red checkmark. The x and y axes are shown, with x=0 at the leading edge of the plate.

Fig 1: Laminar flow of a fluid passed a flat plate and concentration boundary layer

In this lecture, we will consider boundary layer theory in mass transfer and then we will consider mass transfer coefficient in turbulent flow. If we considered a case the flow through a solid surface in a horizontal surface and then we see that the when the fluid pass the solid surface there is a gradient of the velocity and also the gradient of concentration over the surface, near to the surface and the we call a layer which is formed where the concentration variation exist is called the boundary layer.

So, let us see the boundary layer, how it is formed. We have a solid surface over here and when the fluid flow horizontally through the x direction where x is equal to 0 at the start of the velocity profile or the boundary layer and then we see that; no there is a concentration gradient over the surface and this is the edge of concentration boundary layer and then the thickness of this boundary layer is δ_c .

So, an exact solution can be obtained for the hydrodynamic boundary layer for isothermal laminar flow past a plate, if it is a plate. An extension of the Blasius solutions can be extended to derive an expression for convective heat transfer. In the analogous manner we can use the Blasius solutions for convective mass transfer as well for the same geometry and laminar flow.

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Boundary Layer Theory In Mass Transfer

- Here:
 - $C_{A\infty}$ = the concentration of A in the fluid approaching the plate
 - C_{As} = the concentration of A in the fluid adjacent to the surface.
- We start with the differential mass balance and simplifying it for steady state where
$$\frac{\partial C_A}{\partial t} = 0 \quad \checkmark$$

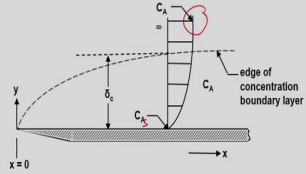


Fig 1: Laminar flow of a fluid passed a flat plate and concentration boundary layer

Now in this case to define the system $C_{A\infty}$ is the concentration of A in the fluid approaching the plate. So, concentration is $C_{A\infty}$ at this location and this is the concentration variation and C_{As} which is at the surface it is C_{As} concentration of A in the fluid adjacent to the surface so, which is over here.

We start with the differential mass balance and simplifying it for steady state process. If we consider the differential mass balance and if we considered the steady state, for steady state we considered $\frac{\partial C_A}{\partial t}$ is equal to 0. So, $\frac{\partial C_A}{\partial t}$ this is equal to 0. Now for the flow which is we considered only in x and y directions.

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Boundary Layer Theory In Mass Transfer

- Flow only in the x and y directions, so $v_z = 0$, and neglecting diffusion in the x and z directions to give

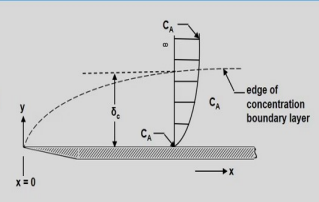
$$v_x \frac{\partial C_A}{\partial x} + v_y \frac{\partial C_A}{\partial y} = D_{AB} \frac{\partial^2 C_A}{\partial y^2} \quad 1$$


Fig 1: Laminar flow of a fluid passed a flat plate and concentration boundary layer

- The momentum boundary layer is very similar:

$$v_x \frac{\partial v_x}{\partial x} + v_y \frac{\partial v_x}{\partial y} = \frac{\mu}{\rho} \frac{\partial^2 v_x}{\partial y^2} \quad 2$$

Then the velocity component that is v_z would be 0 in the z direction and if we neglect the diffusion which is happening towards the x and z direction. So, since the fluid is flowing through the flat surface in the x direction, the convective mass transfer is much prominent than the diffusion. Hence, we can neglect the diffusion part in the x and z direction.

So, we will be know arising out of this know equations which is $v_x \frac{\partial C_A}{\partial x} + v_y \frac{\partial C_A}{\partial y} = D_{AB} \frac{\partial^2 C_A}{\partial y^2}$. So, this we can derive from the fixed second law equations without no chemical reactions and understudy state condition. The momentum layer is also very similar. It will be $v_x \frac{\partial v_x}{\partial x} + v_y \frac{\partial v_x}{\partial y} = \frac{\mu}{\rho} \frac{\partial^2 v_x}{\partial y^2}$. Here μ is the viscosity and ρ is the density.

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Boundary Layer Theory In Mass Transfer

- The thermal boundary layer is also similar:

$$v_x \frac{\partial T}{\partial x} + v_y \frac{\partial T}{\partial y} = \frac{K}{\rho C_p} \frac{\partial^2 T}{\partial y^2}$$
3
- The dimensionless concentration boundary conditions are

$$\frac{v_x}{v_\infty} = \frac{T - T_s}{T_\infty - T_s} = \frac{C_A - C_{As}}{C_{A\infty} - C_{As}} = 0 \quad \text{at } y = 0$$
4

$$\frac{v_x}{v_\infty} = \frac{T - T_s}{T_\infty - T_s} = \frac{C_A - C_{As}}{C_{A\infty} - C_{As}} = 1 \quad \text{at } y = \infty$$
5

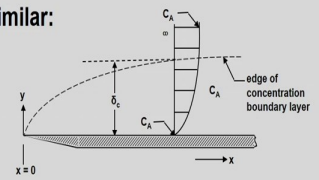


Fig 1: Laminar flow of a fluid passed a flat plate and concentration boundary layer

Now, thermal boundary layer also would be similar which is $v_x \frac{\partial T}{\partial x} + v_y \frac{\partial T}{\partial y}$ is equal to $\frac{k}{\rho C_p} \frac{\partial^2 T}{\partial y^2}$. Now, if we write in terms of the dimensionless concentration in that case we can write $\frac{v_x}{v_\infty}$ by $\frac{v_x}{v_\infty}$. This v_∞ is the known free stream velocity over here is equal to $\frac{T - T_s}{T_\infty - T_s}$ which would be equal to $\frac{C_A - C_{As}}{C_{A\infty} - C_{As}}$ would be equal to 0 at $y = 0$. So, under this conditions, we can write this boundary condition in a dimensionless form. Now similarly we can write at $y = \infty$ at a very long distance from the surface $\frac{v_x}{v_\infty}$ would be equal to $\frac{T - T_s}{T_\infty - T_s}$ which would be equal to $\frac{C_A - C_{As}}{C_{A\infty} - C_{As}}$ is equal to 1.

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Boundary Layer Theory In Mass Transfer

- The Blasius solution is applied to convective heat transfer when

$$\left(\frac{\mu}{\rho}\right) / \alpha = N_{Pr} = 1.0$$
- We use the same type of solution for laminar convective mass transfer when

$$\left(\frac{\mu}{\rho}\right) / D_{AB} = N_{Sc} = 1.0$$
- The velocity gradient was derived in fluid mechanics:

$$\left(\frac{\partial v_x}{\partial y}\right)_{y=0} = 0.332 \frac{V_\infty}{x} N_{Re,x}^{1/2}$$

6 **Where,** $N_{Re,x} = \frac{x v_\infty \rho}{\mu}$ ✓

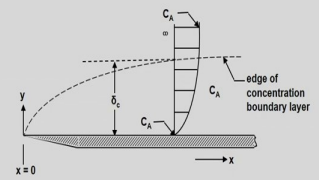


Fig 1: Laminar flow of a fluid passed a flat plate and concentration boundary layer

Now, the Blasius solution is applied to convective heat transfer when μ by ρ by α that is the Prandtl number which is equal to 1. So, when the Prandtl number equal to 1, we can apply the Blasius solutions for convective heat transfer. Similarly in case of mass transfer we use for the laminar convective mass transfer when the Schmidt number equal to 1. If the Schmidt number equal to 1, that is μ by ρ D_{AB} is equal to 1, if that is 1 then we can apply the Blasius solutions.

So, the velocity gradient was derived in fluid mechanics, you have already derived that equations the velocity gradient that is $\frac{\partial v_x}{\partial y}$ at y equal to 0 is equal to $0.332 \frac{v_\infty}{x} N_{Re,x}^{1/2}$. So, $N_{Re,x}$ is Reynolds number. So, Reynolds number to the power half. So, this $N_{Re,x}$ is $x v_\infty \rho$ by μ , which is Reynolds number.

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Boundary Layer Theory In Mass Transfer

- From equation (5)

$$\frac{v_x}{v_\infty} = \frac{C_A - C_{As}}{C_{A\infty} - C_{As}} \quad (7)$$
- Differentiating equation (7) and combining the results with equation (6)

$$\left(\frac{\partial C_A}{\partial y}\right)_{y=0} = (C_{A\infty} - C_{As}) \left(\frac{0.332}{x} N_{Re,x}^{1/2}\right) \quad (8)$$

$\frac{v_x}{v_\infty} = \frac{T - T_s}{T_\infty - T_s} = \frac{C_A - C_{As}}{C_{A\infty} - C_{As}} = 1 \text{ at } y = \infty \quad (5)$

$\left(\frac{\partial C_A}{\partial y}\right)_{y=0} = 0.332 \frac{v_\infty}{x} N_{Re,x}^{1/2} \quad (6)$

Fig 1: Laminar flow of a fluid passed a flat plate and concentration boundary layer

Now, from this equation know 5 which the boundary conditions which we have written in terms of the dimensionless concentration or dimensionless velocity or dimensionless temperature at y equal to infinity, we can write v x by v infinity would be equal to C A minus C As divided by C A infinity minus C As. Now if we differentiate equation 7, this equations if we differentiate and combined with the results of equation 6. So, we can write del C A del y at y equal to 0 would be equal to C A infinity minus C As into 0.332 by x Reynolds number to the power half. So, I think this is very clear for you to understand by differentiating this equation and then substituting this equation 6 in that you will be able to obtain this relation.

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Boundary Layer Theory In Mass Transfer

- The convective mass transfer equation can be written as follows and also related with Fick's equations for dilute solutions:

$$N_{A,y} = k_c (C_{A,s} - C_{A,\infty}) = -D_{AB} \left(\frac{\partial C_A}{\partial y} \right)_{y=0} \quad (9)$$

$$\left(\frac{\partial C_A}{\partial y} \right)_{y=0} = (C_{A,\infty} - C_{A,s}) \left(\frac{0.332}{x} N_{Re,x}^{1/2} \right) \quad (8)$$

Fig 1: Laminar flow of a fluid passed a flat plate and concentration boundary layer

- From eq.(8) and (9) we obtain:

$$\frac{k_c \cdot x}{D_{AB}} = N_{Sh,x} = 0.332 N_{Re,x}^{1/2} \quad (10)$$

- The relationship is restricted to gases with a Schmidt number $N_{Sc} \approx 1.0$.

$$k_c (C_{A,s} - C_{A,\infty}) = -D_{AB} (C_{A,\infty} - C_{A,s}) \frac{0.332}{x} N_{Re,x}^{1/2}$$

$$\Rightarrow \frac{k_c \cdot x}{D_{AB}} = 0.332 N_{Re,x}^{1/2}$$

Now, the convective mass transfer equation can be written as follows and also related with the fix equation for dilute solutions. So, for convective mass transfer, we write the flux N_A at any position y would be equal to $k_c (C_{A,s} - C_{A,\infty})$. So, this is the convective mass flux. Now, if we for a very dilute solution we can relate with the diffusive flux which is $-D_{AB} \frac{\partial C_A}{\partial y}$ at y equal to 0.

Now, we know equation 8 which we have derived which is $\frac{\partial C_A}{\partial y}$ at y equal to 0 is equal to $(C_{A,\infty} - C_{A,s}) \frac{0.332}{x} N_{Re,x}^{1/2}$. So, if we substitute this know $\frac{\partial C_A}{\partial y}$ at y equal to 0 over here so, we will obtain $k_c (C_{A,s} - C_{A,\infty}) = -D_{AB} (C_{A,\infty} - C_{A,s}) \frac{0.332}{x} N_{Re,x}^{1/2}$. So, $k_c \cdot x$ by D_{AB} would be equal to Sherwood number, which is equal to $0.332 N_{Re,x}^{1/2}$.

So, substituting this equations over here like we can write $k_c (C_{A,s} - C_{A,\infty})$ would be equal to $-D_{AB} (C_{A,\infty} - C_{A,s}) \frac{0.332}{x} N_{Re,x}^{1/2}$. Now, this if we rearrange this will canceled out and we will get $k_c \cdot x$ divided by D_{AB} would be equal to $0.332 N_{Re,x}^{1/2}$. So, this is nothing but say sherwood number with respect to x . So, we can obtain this equation 10. So, this relation is valid when this schmidt number is 1; so, which we have said before.

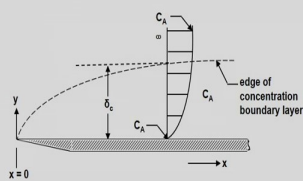
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Boundary Layer Theory In Mass Transfer

- The relation between the thickness δ of the hydrodynamic and δ_c of the concentration boundary layers where the $N_{Sc} \neq 1.0$ is

$$\frac{\delta}{\delta_c} = N_{Sc}^{1/3} \quad \text{11}$$

- As a result, the equation for local convective mass transfer coefficient is

$$\frac{k_c' x}{D_{AB}} = N_{Sh,x} = 0.332 N_{Re,x}^{1/2} N_{Sc}^{1/3} \quad \text{12}$$


The diagram shows a flat plate at the bottom with a coordinate system where x is the distance along the plate and y is the vertical distance. A hydrodynamic boundary layer is shown as a shaded region starting from the leading edge at x=0. A concentration boundary layer is shown as a region above the hydrodynamic one, with its thickness labeled as δ_c . The concentration profile is shown as a curve starting from C_A at the plate and reaching a free stream value $C_A = \infty$. The edge of the concentration boundary layer is indicated by a dashed line.

Fig 1: Laminar flow of a fluid passed a flat plate and concentration boundary layer

Now, the relations between the thickness δ of the hydrodynamic and the δ_c of the concentration boundary layer. So, the hydrodynamic boundary layer thickness is δ and the concentration boundary layer thickness is δ_c . When Schmidt number is not equal to 1, if Schmidt number is not equal to 1, then there is a relations between the hydrodynamic boundary layer thickness and the concentration boundary layer thickness and which is δ / δ_c would be equal to Schmidt number to the power one-third. The equation for local convective mass transfer coefficient that is $k_c' x$ by D_{AB} which is Sherwood number x with respect to x would be equal $0.332 \text{Re}_x^{1/2} \text{Sc}^{1/3}$. So, this is question 12.

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Boundary Layer Theory In Mass Transfer

- We can obtain the mean mass transfer coefficient $\overline{k_c}$ from $x = 0$ to $x = L$ for a plate of width b by integrating as follows:

$$\overline{k_c} = \frac{b}{bL} \int_0^L k_c' dx \quad \text{13}$$
- The result is $\frac{\overline{k_c} L}{D_{AB}} = N_{Sh,L} = 0.664 N_{Re,L}^{1/2} N_{Sc}^{1/3}$ 14
- This is similar to the heat transfer equation for a flat plate.

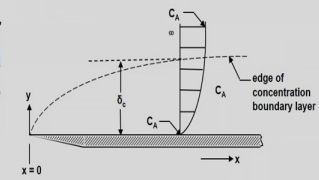


Fig 1: Laminar flow of a fluid passed a flat plate and concentration boundary layer

We can obtain the mean mass transfer coefficient $\overline{k_c}$ from x is equal to 0 to x is equal to L for a plate of width b by integrating as follows. So, we can integrate $\overline{k_c}$ would be equal to $\frac{b}{bL} \int_0^L k_c' dx$. Now if we substitute $\overline{k_c}$ which we have derived before $\overline{k_c}$ from here then, we would be able to obtain $\overline{k_c} L$ divided by D_{AB} is equal to Sherwood number N with respect to Sherwood number would be equal to 0.664 Reynolds number to the power half into Schmidt number to the power one-third. So, the derivation after substituting $\overline{k_c}$ from the equation 13 to 14 is left to you as a part of your homework, you can do it very easily. So, this is similar to the heat transfer equations for a flat plate.

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Example 1

A large volume of pure water at 25 °C is flowing parallel to a flat plate of solid benzoic acid, where $L = 0.244$ m in the direction of flow. The water velocity is 0.061 m/s. The solubility of benzoic acid in water is 0.02948 kmol/m³. The diffusivity of benzoic acid is 1.245×10^{-9} m²/s. Calculate the mass transfer coefficient $k_{L,avg}$ and the flux N_A . Given that $\mu = 8.71 \times 10^{-4}$ kg/m s and $\rho = 996$ kg/m³

Now, let us take an example, a large volume of pure water at 25 degree centigrade is flowing parallel to a flat surface of solid benzoic acid, where L is 0.244 meter in the direction of flow. The water velocity is 0.061 meter per second. The solubility of benzoic acid in water is 0.02948 kilo mole per meter cube. The diffusivity of benzoic acid is 1.245 into 10 to the power minus 9 meter square per second. We need to calculate the average mass transfer coefficient $k_{L,avg}$ and the flux N_A and the viscosity and the density are given.

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Example 1: Solution

A large volume of pure water at 25 °C is flowing parallel to a flat plate of solid benzoic acid, where $L = 0.244$ m in the direction of flow. The water velocity is 0.061 m/s. The solubility of benzoic acid in water is 0.02948 kmol/m³. The diffusivity of benzoic acid is 1.245×10^{-9} m²/s. Calculate the mass transfer coefficient $k_{L,avg}$ and the flux N_A . Given that $\mu = 8.71 \times 10^{-4}$ kg/m s and $\rho = 996$ kg/m³

Solution:

- The Schmidt number is

$$N_{Sc} = \frac{\mu}{\rho D_{AB}} = \frac{8.71 \times 10^{-4}}{996(1.245 \times 10^{-9})} = 702$$

- The Reynold's number is

$$N_{Re,L} = \frac{L u \rho}{\mu} = \frac{0.244(0.061)(996)}{8.71 \times 10^{-4}} = 1.7 \times 10^4$$

So, now we need to calculate the Schmidt number. So, Schmidt number is μ by ρD_{AB} , μ is given 8.71×10^{-4} kg per meter second, ρ is 996 and diffusivity is given. So, you can calculate the Schmidt number. Similarly, we can calculate the Reynolds number which is $L u \rho$ by μ L is given 0.244 meter and then we know the velocity 0.061 meter per second. So, velocity is given L is known and ρ is given 996 kg per meter cube and also the viscosity. So, we can calculate know the Reynolds number 1.7×10^4 , which is in the turbulent region.

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Example 1: Solution

A large volume of pure water at 25 °C is flowing parallel to a flat plate of solid benzoic acid, where $L = 0.244$ m in the direction of flow. The water velocity is 0.061 m/s. The solubility of benzoic acid in water is 0.02948 kmol/m³. The diffusivity of benzoic acid is 1.245×10^{-9} m²/s. Calculate the mass transfer coefficient $k_{L,avg}$ and the flux N_A . Given that $\mu = 8.71 \times 10^{-4}$ kg/m s and $\rho = 996$ kg/m³

Solution: (Cont.)

$$\frac{\bar{k}_c L}{D_{AB}} = 0.664 N_{Re,L}^{1/2} N_{Sc}^{1/3}$$

$$\therefore \bar{k}_c = 3.92 \times 10^{-6} \text{ m/s}$$

- In this case, diffusion is for A through non-diffusing B,

$$N_A = \frac{\bar{k}_c}{x_{BM}} (C_{A1} - C_{A2}) = k_c (C_{A1} - C_{A2})$$

Handwritten notes:
 $N_{Re,L} = 1.7 \times 10^4$
 $N_{Sc} = 702$
 $D_{AB} = 1.245 \times 10^{-9}$
 $L = 0.244 \text{ m}$

Now, if we just use the know value of N Reynolds number L which is know 1.7×10^4 to the power 4 and then Schmidt number which is 702, substituting this we would be and then the diffusion coefficient is known to us D_{AB} is given 1.245×10^{-9} meter square per second and then the L which is given 0.244 meter. So, if we substitute we will get k_c dash average would be 3.92×10^{-6} meter square per second.

In this case, diffusion of A through non diffusing B we can write the flux, N_A is k_c dash average by x_{BM} C_{A1} minus C_{A2} would be equal to k_c C_{A1} minus C_{A2} .

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Example 1: Solution

A large volume of pure water at 25 °C is flowing parallel to a flat plate of solid benzoic acid, where $L = 0.244$ m in the direction of flow. The water velocity is 0.061 m/s. The solubility of benzoic acid in water is 0.02948 kmol/m³. The diffusivity of benzoic acid is 1.245×10^{-9} m²/s. Calculate the mass transfer coefficient $k_{L,avg}$ and the flux N_A . Given that $\mu = 8.71 \times 10^{-4}$ kg/m s and $\rho = 996$ kg/m³

Solution: (Cont.)

- Since the solution is very dilute $x_{BM} \cong 1.0$ ✓
and $\bar{k}_c \cong k_c$ ✓
- Also, $C_{A1} = 2.948 \times 10^{-2}$ kg mol/m³ (solubility) ✓
- and $C_{A2} = 0$ (large volume of fresh water).

$$\therefore N_A = (3.92 \times 10^{-6}) \times (0.02948 - 0) = 1.156 \times 10^{-7} \text{ kg mol/m}^2\text{s}$$

Now, if we considered very dilute solutions x_{BM} would be approximately equal to 1 and \bar{k}_c approximately equal to k_c . So, the concentration C_{A1} which is 2.948×10^{-2} kg mole or kmol per meter cube which is given a solubility over here. So, C_{A1} is known and since the large volume of fresh air it is said a large volume of pure water once it said pure water. So, C_{A2} would be equal to 0. So, then we can write the flux would be equal to k_c into the concentration gradient $0.02948 - 0$. So, which would be equal to 1.156×10^{-7} kg mole per meter square second.

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Mass Transfer Coefficient in Turbulent Flow

- There are many theories which attempt to interpret or explain the behavior of mass transfer coefficients.

Such as > **Film theory** → Nernst (1904)

> **Penetration theory** → Higbie (1935)

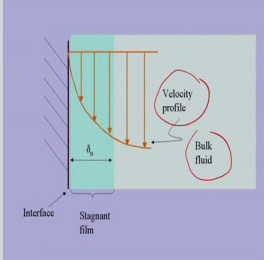
> **Surface-renewal theory** → Danckwerts (1951)

Now, we will consider mass transfer coefficient in turbulent flow. There are many theories which attempt to interpret or explain the behavior of mass transfer coefficient. So, we will try to cover few theories which are applicable in case of the turbulent flow. One of them is film theory which is proposed by Nernst in 1904. The second theory turbulent flow turbulent mass transfer coefficient is penetration theory, which is proposed by Higbie in 1935. So, it is popularly known as Higbie's theory and then third one is surface renewal theory which is proposed by Danckwerts and it is in 1951. So, this is you know evolution of this theory film theory, penetration theory and surface renewal theory we will try to explain them sequentially.

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Film Theory

- Nernst postulated that near the interface there exists a **stagnant film**.
- This stagnant film is **hypothetical** since we really don't know the details of the velocity profile near the interface.
- **Basic concept** – the resistance to diffusion can be considered equivalent to that in stagnant film of a certain thickness



Now, you can see there is an interface over here and there is a stagnant film over here which is given as δ_n and you could see the fluid bulk and the velocity profile look like this. The Nernst postulated that near the interface there exist or stagnant film. So, this film is considered as stagnant film at a gas liquid interface towards the liquid site. The stagnant film which is hypothetical this is considered as hypothetical since we really do not know the details of the velocity profile near the interface. So, this is considered as hypothetical.

The basic concept the resistance to diffusion can be considered equivalent to that in the stagnant film of a certain thickness. So, most of the resistance which create for mass

transfer is due to the stagnant film. So, most of the resistance is considered in this stagnant film and has a certain thickness of δ considered over here.

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Film Theory

- Mass transfer occurs by molecular diffusion through a fluid layer at phase boundary (solid wall). Beyond this film, concentration is homogeneous and is C_{Ab} .
- Mass transfer through the film occurs at steady state.
- Flux is low and mass transfer occurs at low concentration.

Figure. Concentration profile with stagnant film

Hence,

$$N_A = \frac{-D_{AB} dC_A}{dZ}$$

Now, you can see the concentration profile with the stagnant film how it is varying, this is δ the thickness and the concentration is varying C_A at the interface the concentration is C_{Ai} and at the bulk the concentration is C_{Ab} . The stagnant film concentration you can see this is a linear drop in concentration, but you can see the actual profile look like this, this is the actual profile, we will discuss know this things over here.

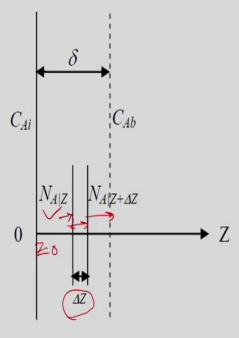
Mass transfer occur by molecular diffusion through the fluid layer at phase boundary that is at solid wall. Beyond this film the concentration is homogeneous and is C_{Ab} . So, at the face boundary know the molecular diffusion takes place and beyond at a certain distance δ concentration is homogeneous that is considered as C_{Ab} . Mass transfer through stagnant film which occurs in this film is basically steady state. So, flux in this film is very low and mass transfer occurs at very low concentrations. Hence, we can write N_A would be equal to minus $D_{AB} dC_A / dZ$.

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Film Theory

- Steady state mass balance is done over an elementary volume of thickness ΔZ
- Rate of input of solute at $Z = N_A|_Z$
- Rate of output of solute at $Z + \Delta Z = N_A|_{Z + \Delta Z}$
- Rate of accumulation = 0
= rate of input - rate of output

Therefore, At steady state:

$$N_A|_Z - N_A|_{Z + \Delta Z} = 0$$


Now, if we do the steady state mass balance over an elementary volume of thickness ΔZ . So, considered this is a elemental volume of ΔZ this one thickness and flux at Z , Z is varying from the surface interface towards this side. So, at Z is equal to 0 over here at Z is equal to Z the flux is considered $N_A|_Z$ that is the rate of input of solute at Z . So, over here at this location and the rate which is out over here is $N_A|_{Z + \Delta Z}$. So, rate of output of solute at Z plus ΔZ would be equal to $N_A|_{Z + \Delta Z}$. The accumulation is considered here is 0, the rate of accumulation is 0. So, which is equal to rate of input minus rate of output. So, from this we can write at steady state $N_A|_Z - N_A|_{Z + \Delta Z} = 0$.

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Film Theory

1 $\lim_{\Delta Z \rightarrow 0} \frac{N_A|_Z - N_A|_{Z+\Delta Z}}{\Delta Z} = 0$

2 $\Rightarrow \frac{dN_A}{dZ} = 0$

3 $\Rightarrow \frac{d}{dZ}(-D_{AB} \frac{dC_A}{dZ}) = 0$

4 $\Rightarrow -D_{AB} \frac{d^2 C_A}{dZ^2} = 0$

5 $\Rightarrow \frac{d^2 C_A}{dZ^2} = 0$

So, if we considered very small thickness δ and taking limit $\delta \rightarrow 0$ we can write this equation N_A at Z minus N_A at $Z + \delta$ by δ would be equal to 0. So, we can write $\frac{dN_A}{dZ}$ would be equal to 0. Now N_A if we substitute this N_A from the Fick's law which is $-D_{AB} \frac{dC_A}{dZ}$ which is equal to 0. So, we can write $-D_{AB} \frac{d^2 C_A}{dZ^2}$ would be equal to zero since, $D_{AB} \neq 0$.

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Film Theory

- Integrating Equation (4) for the following boundary conditions:
 - $C_A = C_{Ai}$ when $Z=0$
 - $C_A = C_{Ab}$ when $Z=\delta$
 We have now:

$$C_A = C_{Ai} - (C_{Ai} - C_{Ab}) \frac{Z}{\delta}$$
- Hence, according to film theory, **concentration profile in stagnant film is linear**
- Molar flux through film, $N_A = -D_{AB} \frac{dC_A}{dZ} \Big|_{Z=0}$

$$N_A = \frac{D_{AB} (C_{Ai} - C_{Ab})}{\delta}$$

4 $-D_{AB} \frac{d^2 C_A}{dZ^2} = 0$

Now, if we integrate this equation 4 that is $-D_{AB} \frac{d^2 C_A}{dZ^2} = 0$ with the boundary conditions that is at C_A would be C_{Ai} at Z is equal to 0 at the interface it will

be C_{Ai} that is at Z is equal to 0 and C_{Ab} when Z is equal to δ . So, that is after the thickness it will remain C_{Ab} the bulk concentration. So, we have now if we just integrate these equations we can obtain C_A would be equal to C_{Ai} minus C_{Ab} into Z by δ . So, the integration with the boundary conditions given over here is left to you as a homework. So, you can just know integrate this equation and we will get this concentration profile.

Now, according to film theory we can write the concentration profile in stagnant film is linear. As per film theory the concentration profile over here would be linear and molar flux through the film N_A we can write N_A would be equal to minus $D_{AB} \frac{dC_A}{dz}$ at Z is equal to 0. So, if we differentiate this equation 6 and then substitute in this relation we will get N_A would be equal to $D_{AB} \frac{C_{Ai} - C_{Ab}}{\delta}$.

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Film Theory

Now comparing eq. (7) with definition of mass transfer coefficient

$$k_L = \frac{D_{AB}}{\delta}$$
8

$$N_A = k_L (C_{Ai} - C_{Ab})$$

$$N_A = \frac{D_{AB} (C_{Ai} - C_{Ab})}{\delta}$$
7

- In this film transport is governed essentially by **molecular diffusion**.
- Therefore, Fick's law describes flux through the film.

$$J = -D \frac{\partial C}{\partial X} \quad \left(\text{typical units } \frac{\text{mg}}{\text{cm}^2 \cdot \text{sec}} \right)$$

Now, if we compare equation 7 with definition of mass transfer coefficient. So, this is equation 7 and if we compare with the mass transfer coefficient equations that is know mass transfer coefficient is k_L and the flux equation we can write N_A is equal to k_L into C_{Ai} minus C_{Ab} . So, if we compare these two relation we can calculate k_L would be D_{AB} by δ . So, in this film transport is governed essentially by molecular diffusion. Therefore, Fick's law describes flux through the film and we can write J is equal to minus $D \frac{dC}{dX}$ and typically you need for this is milligram per meter square second, it is mass flux.

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Film Theory

- If the thickness of the stagnant film is given by δ_n then the gradient can be approximated by:
$$\frac{\partial C}{\partial X} \approx \frac{C_b - C_i}{\delta_n}$$

C_b and C_i are concentrations in the bulk and at the interface, respectively.
- At steady-state if there are no reactions in the stagnant film there will be no accumulation in the film (Assume that $D = \text{constant}$).
 - Therefore the gradient must be linear and the approximation is appropriate.
 - And:
$$J = -D \frac{(C_b - C_i)}{\delta_n}$$

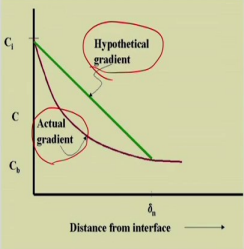
If the thickness of the stagnant film is given by δ_n then the gradient can be approximated by $\frac{dC}{dX}$ which is equal to $\frac{C_b - C_i}{\delta_n}$. And C_b and C_i are concentration in the bulk and at the interface. So, we can add any components. So, we can write this equation at steady state if there is no reactions in the stagnant film there will be no accumulation in the film that is D is constant and therefore, the gradient must be linear and then the approximation is appropriate which we have considered. So, we can write J would be equal to minus $D \frac{C_b - C_i}{\delta_n}$.

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Film Theory

- Calculation of C_i is done by assuming that equilibrium (Henry's Law) is attained instantly at the interface.
- This assumes that the other phase doesn't have a "film". So for the moment
$$C_i = \frac{C_g}{H_C}$$

(if the film side is liquid and the opposite side is the gas phase).
- A problem with the model is that the effective diffusion coefficient is seldom constant since some turbulence does enter the film area.
- So the concentration profile in the film looks more like:



So, calculation of C_i is done by assuming that the equilibrium that is the Henry's law is attained instantly at the interface. So, when we know if we assume the equilibrium which will be attained at the interface very instantaneously then we can apply the Henry's law at the interface to calculate the interfacial concentration. So, C_i we would be able to calculate by Henry's law if the equilibrium is achieved at the interface instantaneously. This assumes that the other phase does not have a film. So, for the moment C_i would be equal to C_g by H_c Henry's law H_c is Henry's law constant and C_g is the gas phase concentration of that component. If the film side is liquid and the opposite side is the gas phase. So, C_i the interfacial concentration we can calculate C_g by H_c . So, this is essentially considered that the gas there is no resistance in the gas phase. A problem with the model is that, the effective diffusion coefficient is seldom constant since some turbulence does occur in the film area.

So, the film theory which states that in the film the concentration profile is linear, but actually it is not linear because the effective diffusion coefficient is not constant because it varies due to the turbulence in the film, which actually exist. So, the concentration profile in the film will look like actual gradient will look like this because of the turbulence which exist, but the hypothetical gradient as per the film theory is linear.

Thank you for hearing this lecture and we will continue with the mass transfer coefficient in turbulent flow in the next lecture.