

Introduction to Interfacial Waves
Prof. Ratul Dasgupta
Department of Chemical Engineering
Indian Institute of Technology, Bombay

Lecture - 23
Duffing equation (contd.)

We were looking at the solution to the Duffing equation using the method of multiple scales, we had written down the equation and order epsilon.

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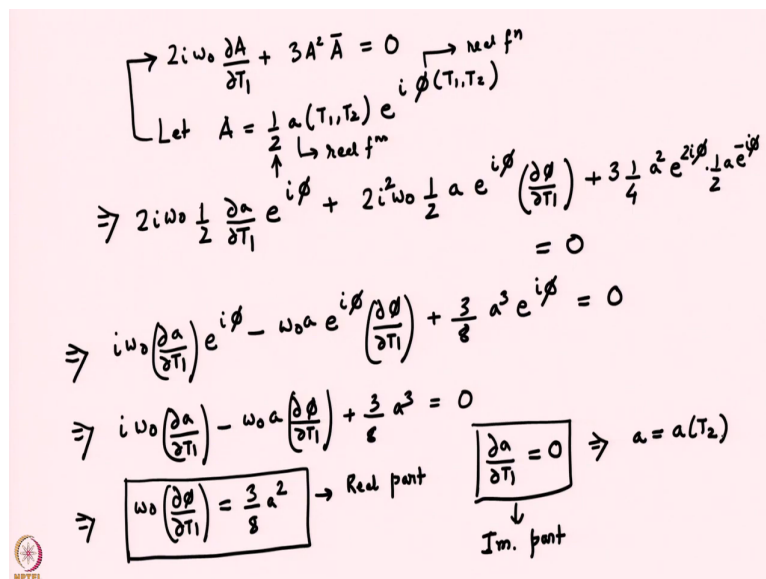
$$\begin{aligned}
 \underline{O(1)}: \quad u_0 &= A(T_1, T_2) e^{i\omega_0 T_0} + \text{c.c.} \quad \leftarrow \\
 \underline{O(\epsilon)}: \quad (D_0^2 + \omega_0^2) u_1 &= -2D_0 D_1 u_0 - u_0^3 \\
 &= \boxed{-2i\omega_0 \frac{\partial A}{\partial T_1} e^{i\omega_0 T_0}} - (A e^{i\omega_0 T_0} + \text{c.c.})^3 + \text{c.c.} \\
 &\quad \downarrow \text{MISTAKE CORRECTED} \\
 \Rightarrow &= -2i\omega_0 \frac{\partial A}{\partial T_1} e^{i\omega_0 T_0} - A^3 e^{3i\omega_0 T_0} - 3A^2 \bar{A} e^{i\omega_0 T_0} + \text{c.c.} \\
 &= \underbrace{\left(-2i\omega_0 \frac{\partial A}{\partial T_1} + 3A^2 \bar{A} \right)}_{=0} e^{i\omega_0 T_0} - A^3 e^{3i\omega_0 T_0} + \text{c.c.} \\
 &\quad \text{Eliminate resonant forcing}
 \end{aligned}$$

Note that there was a small error in writing down the first term. I have indicated that in a red box here. We had also got terms from this second term which was to the power 3 and I had shown you that how I get this term and that term and then, there are complex conjugates of those which we do not write down.

So, we have this line which is the first term and then, the next two terms come from the cubic term and then, plus c.c. which indicates that complex conjugate of all the three terms. It is useful to put in 1 bracket all terms which have e to the power $i\omega_0 T_0$. So, that I have done here and then, e to the power $i3i\omega_0 T_0$ is there and then, there is a complex conjugate part.

Now, the reason why we are looking at terms of the which are proportional to e to the power $i\omega_0 T_0$ is because like earlier, they are resonant forcing terms, they basically these are terms which will oscillate at the same frequency as the natural frequency of the oscillator. The it, this is a solution to the homogeneous equation $d^2 u/dt^2 + \omega_0^2 u = 0$. So, we have to set this term equal to 0 to eliminate resonant forcing and this will give us an equation for capital A as a function of T1. Let us solve that equation.

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$$\begin{aligned}
 & \rightarrow 2i\omega_0 \frac{\partial A}{\partial T_1} + 3A^2 \bar{A} = 0 \quad \text{Let } A = \frac{1}{2} a(T_1, T_2) e^{i\phi(T_1, T_2)} \\
 & \Rightarrow 2i\omega_0 \frac{1}{2} \frac{\partial a}{\partial T_1} e^{i\phi} + 2i^2 \omega_0 \frac{1}{2} a e^{i\phi} \left(\frac{\partial \phi}{\partial T_1} \right) + 3 \frac{1}{4} a^2 e^{2i\phi} \frac{1}{2} a e^{-i\phi} = 0 \\
 & \Rightarrow i\omega_0 \left(\frac{\partial a}{\partial T_1} \right) e^{i\phi} - \omega_0 a e^{i\phi} \left(\frac{\partial \phi}{\partial T_1} \right) + \frac{3}{8} a^3 e^{i\phi} = 0 \\
 & \Rightarrow i\omega_0 \left(\frac{\partial a}{\partial T_1} \right) - \omega_0 a \left(\frac{\partial \phi}{\partial T_1} \right) + \frac{3}{8} a^3 = 0 \\
 & \Rightarrow \boxed{\omega_0 \left(\frac{\partial \phi}{\partial T_1} \right) = \frac{3}{8} a^2} \rightarrow \text{Real part} \quad \boxed{\frac{\partial a}{\partial T_1} = 0} \rightarrow \text{Im. part} \Rightarrow a = a(T_2)
 \end{aligned}$$

So, the equation is twice $i\omega_0$ to 0. For convenience, we choose A , A is a complex function. So, A is equal to half small a which is a function of T_1 and T_2 into e to the power some ϕ , some ϕ which is also a function of T_1 and T_2 . So, now, if I go back and replace it in this equation, then we obtain twice $i\omega_0$. This half is just for convenience because there is a factor of 2 in the equation.

So, if you put a half, things will get cancelled out. So, I have half $\frac{\partial a}{\partial T_1} e$ to the power $i\phi$ and then, we will have twice $i\omega_0$ and then, I have to take the derivative of e to the power $i\phi$ which is just i . So, that will give me i square and then, I have a factor of half $A e$ to the power $i\phi$ into $\frac{\partial \phi}{\partial T_1}$.

This is the derivative of the first term in the equation which governs A . I have one more term which is 3 and A square makes it 1 by 4 small a square e to the power twice $i\phi$ into A bar; A bar is just the complex conjugate of capital A . So, half small a into e to the power minus $i\phi$ is equal to 0. So, we get this equation which is $i\omega_0 \frac{\partial a}{\partial T_1} e$ to the power $i\phi$ minus, minus because there is an i square and then, this is $\omega_0 a e$ to the power $i\phi \frac{\partial \phi}{\partial T_1} + \frac{3}{8} a^2 e$ to the power $2i\phi$ and you can see that this is there is an a cube here and e to the power $i\phi$ is equal to 0.

So, we can simplify this by getting rid of e to the power $i\phi$ because e to the power $i\phi$ is in general not 0. So, this simplifies to. Recall that small a and small ϕ are actually real functions because we are writing them in complex notation. So, it is a times e to the power $i\phi$. So, a is a real function and ϕ is also a real function. So, this equation has an imaginary term here, you can see that the first term in the equation is has a i .

So, this equation can in general be split into a real part and an imaginary part and both of them have to be separately equated to 0. So, we will have. So, if I the real part is just the second and the third terms, so though that gives me $\omega_0 \frac{\partial \phi}{\partial T_1} a$ is equal to $\frac{3}{8} a^2$. I have cancelled out one a . This is the real part.

The imaginary part is just the first term and we have to say because ω_0 is not 0. So, $\frac{\partial \phi}{\partial T_1}$ is equal to 0, this is the imaginary part and this implies that a is not a function of T_1 . So, a at most is a function of T_2 . We will later take a to be a constant because we are not really going to solve the problem up to order ϵ^2 .

So, if there is any variation of a , you will find it only at order only at long enough times of the order ϵ^2 into t . So, of the time small t is of the order 1 by ϵ^2 . So, we are later going to take a as a constant, but let us work on the real part.

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Handwritten mathematical derivation on a pink background:

$$\frac{\partial \phi}{\partial T_1} = \frac{3}{8\omega_0} a^2$$

$$\Rightarrow \phi = \frac{3}{8\omega_0} a^2 T_1 + \phi_0(T_2) \leftarrow$$

$$A = \frac{1}{2} a(T_1, T_2) e^{i\phi(T_1, T_2)}$$

$$= \frac{1}{2} a(T_2) e^{i\left\{\frac{3\epsilon^2}{8\omega_0} T_1 + \phi_0(T_2)\right\}}$$

$$(D_0^2 + \omega_0^2) u_1 = -A^3 e^{3i\omega_0 T_0} + \text{c.c.}$$

P.I. : $\alpha e^{3i\omega_0 T_0}$ (Substitute)

$\alpha = \frac{A^3}{8\omega_0^2}$

\therefore general soln is

$$u_1(T_0, T_1, T_2) = B(T_1, T_2) e^{i\omega_0 T_0} + \frac{A^3}{8\omega_0^2} e^{3i\omega_0 T_0} + \text{c.c.}$$

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So, the real part gives us $\frac{\partial \phi}{\partial T_1}$ is equal to $\frac{3}{8\omega_0} a^2$ and we have just seen that a is not a function of T_1 . So, while doing this integration, you can treat a

as a constant because a is at most a function of T_1 . So, this integration, we are doing with respect to T_1 .

So, ϕ just becomes $3 \text{ by } 8 \omega_0 a^2 T_1$ plus some constant ϕ_0 which could potentially be a function of T_2 . But once again, we will treat ϕ_0 as a constant because as I have said before, we are not going to solve the problem up to T_2 or up to order ϵ^2 . But we will need to go write down equations up to order ϵ^2 in order to determine the solution up to order ϵ completely.

So, we thus, obtain a , we recall that we had written capital A as small a which was a function of which we thought was a function of T_1, T_2 into e to the power $i \phi$ which is also a function of T_1, T_2 . Now, we have determined that small a is not a function of T_1 . So, I will just write it as a function of T_2 and later, take it to be a constant at order ϵ .

So, this into and I am going to write down whatever we found for ϕ here. So, this would be $3 a^2 \text{ by } 8 \omega_0 \text{ naught } T_1$ plus some function ϕ_0 of this. So, we now know how a depends on T_1 . So, now, the solution at this order, so we are looking at the problem at order ϵ and so, the solution at this order looks like $D_0^2 \text{ plus } \omega_0^2 \text{ into } u_1$.

Recall that we had two terms, we had a part which depended on e to the power $i \omega_0 \text{ naught } T_0$ and another part which depended on e to the power $3 i \omega_0 \text{ naught } T_0$ you can see that the first part has been set equal to 0. So, that is not there anymore. So, the equation at this order is just equal to the second term plus its complex conjugate.

If we write that, then it is just minus $A^3 e$ to the power $3 i \omega_0 \text{ naught } T_0$ plus complex conjugate. The term which is proportional to e to the power $i \omega_0 \text{ naught } T_0$, it has been eliminated by setting its coefficient equal to 0 and that has told us what is the expression for capital A ok.

So, now, we have to determine. So, you can see that this has the structure, this although this is an a partial differential equation, D_0 is actually $\partial^2 / \partial T_0^2$; but we

can solve it like an ordinary differential equation and we have to determine the particular integral. The particular integral in this case has to be proportional to $e^{3i\omega_0 t}$.

So, we set it equal to $\alpha e^{3i\omega_0 t}$ and if we substitute it into this equation, then we are supposed to determine the value of α . So, if you substitute it here, you will find. So, substitute and you will find that α is equal to a cube by $8\omega_0^2$. Just one can substitute and find this.

Therefore, the general solution at order ϵ is $u_1(t)$, $u_2(t)$ is some the complementary function some constant which potentially depends on t , $u_3(t)$ into $e^{i\omega_0 t}$ plus the particular integral which is A cube α times $e^{3i\omega_0 t}$; α , we have determined to be A cube by $8\omega_0^2$.

So, this is just $\alpha e^{i\omega_0 t}$ plus complex conjugate. This is the structure of the general solution at this order. Now, we cannot stop at this order, the reason being that we do not know B as a function of t . If we knew B as a function of t , the problem would be over at order ϵ ; we would not have to proceed to the next order.

Recall that I have told you that we are not going to solve the problem correctly all the way up to order ϵ^2 . So, we do not need to worry about the t^2 dependence of B , of B . But we do need to worry about the t dependence of B ok. So, for that, let us proceed to the next order to see if we can find an equation which tells us what is B as a function of t .

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$$\begin{aligned}
 \text{O}(\epsilon^4): (D_0^2 + \omega_0^2) u_2 &= \textcircled{1} + \textcircled{2} + \textcircled{3} + \textcircled{4} \\
 \textcircled{1}: -2D_0 D_1 u_1 &= -2D_0 D_1 \left[B e^{i\omega_0 T_0} + \frac{A^3}{8\omega_0^3} e^{3i\omega_0 T_0} + \text{c.c.} \right] \\
 &= -2D_1 \left[i\omega_0 B e^{i\omega_0 T_0} + \frac{3i\omega_0 A^3}{8\omega_0^3} e^{3i\omega_0 T_0} + \text{c.c.} \right] \\
 &= -2i\omega_0 (D_1 B) e^{i\omega_0 T_0} - \frac{6i}{8\omega_0} 3A^2 (D_1 A) e^{3i\omega_0 T_0} + \text{c.c.} \\
 &\quad D_1 A = -\frac{3A^2 \bar{A}}{2i\omega_0} \left[\text{From } 2i\omega_0 \frac{\partial A}{\partial T_1} + 3A^2 \bar{A} = 0 \right] \\
 &\quad \quad \quad \downarrow D_1 A \\
 \textcircled{1} &= -2i\omega_0 (D_1 B) e^{i\omega_0 T_0} + \frac{27}{8\omega_0^2} A^4 \bar{A} e^{3i\omega_0 T_0}
 \end{aligned}$$

$D_1 A^3$
 $= \frac{\partial}{\partial T_1} (A^3)$
 $= 3A^2 (D_1 A)$

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$$\begin{aligned}
 & \frac{d^2 u}{dt^2} + \omega_0^2 u + \varepsilon u^3 = 0 \\
 & T_0 \equiv t, \quad T_1 \equiv \varepsilon t, \quad T_2 \equiv \varepsilon^2 t \\
 & \frac{d}{dt} = (D_0 + \varepsilon D_1 + \varepsilon^2 D_2 + \dots) \\
 & \frac{d^2}{dt^2} = (D_0^2 + 2\varepsilon D_0 D_1 + \varepsilon^2 D_1^2 + 2\varepsilon^2 D_0 D_2) \\
 & u = u_0(T_0, T_1, T_2) + \varepsilon u_1(T_0, T_1, T_2) + \varepsilon^2 u_2(T_0, T_1, T_2) + \dots
 \end{aligned}$$

$$\begin{aligned}
 O(1) : & (D_0^2 + \omega_0^2) u_0 = 0 \\
 O(\varepsilon) : & (D_0^2 + \omega_0^2) u_1 = -2 D_0 D_1 u_0 - u_0^3 \\
 O(\varepsilon^2) : & (D_0^2 + \omega_0^2) u_2 = -2 D_0 D_1 u_1 - 2 D_0 D_2 u_0 - D_1^2 u_0 - 3 u_0^2 u_1
 \end{aligned}$$

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So, let us do that. So, this is the equation at order epsilon square. I have already written this equation here, there are four terms on the right. So, I will label them 1, 2, 3, 4. Some amount of algebra needs to be done in evaluating these terms. The algebra tends to be slightly lengthy.

I am going to work out some of these terms and I am going to leave the rest of the terms to you as an exercise. The exercise is quite straightforward and I will tell you how to do it. I leave it to you to show that the expressions are indeed what I will write down here.

So, just let me rewrite the equation once again. So, at order epsilon square $D_0^2 + \omega_0^2 u_2$ is equal to 1 plus 2 plus 3 plus 4, 4 terms. What is 1? 1 is minus 2 $D_0 D_1 u_1$. I will work out this term. So, minus 2 $D_0 D_1 u_1$ yes and we have just found that u_1 ,

I have written it in the last slide is just B plus the particular integral that we just found plus complex conjugate.

So, I have to differentiate this expression. So, this is minus $2 D_1$. Next the differentiation with respect to D_0 , B does not participate in it because B is just a function of T_1 and T_2 . So, this will just become $i \omega_0 B$ and from the second term, similarly A is again a function of T_1 and T_2 that will also not participate. So, I will just get thrice $i \omega_0$ A^3 .

Now, I am going to differentiate with respect to T_1 or I am going to operate the D_1 on B . So, you can see that this is going to become twice $i \omega_0 D_1$ operating on B which is basically $\frac{\partial B}{\partial T_1}$ into e to the power $i \omega_0 T_0$ minus $6 i$. If I cancel out an ω_0 , then I get ω_0 only in the denominator and then, I have D_1 of A^3 that can be simplified D_1 is just $\frac{\partial}{\partial T_1}$, so D_1 of A^3 . This is what I have to do.

This is just $\frac{\partial}{\partial T_1}$ of A^3 which is $3 A^2 D_1$ of A . So, this becomes $3 A^2 \frac{\partial A}{\partial T_1}$ into e to the power $i \omega_0 T_0$ plus c.c. Now, this is the place, where I am going to introduce some simplification which will help the algebra. We have an expression for D_1 of A from before.

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$$\begin{aligned}
 & 2i\omega_0 \frac{\partial A}{\partial T_1} + 3A^2 \bar{A} = 0 \quad \left\{ \begin{array}{l} \text{real fn} \\ \phi(T_1, T_2) \end{array} \right. \\
 & \left\{ \begin{array}{l} \text{Let } A = \frac{1}{2} a(T_1, T_2) e^{i\phi} \\ \quad \quad \quad \uparrow \text{real fn} \end{array} \right. \\
 & \Rightarrow 2i\omega_0 \frac{1}{2} \frac{\partial a}{\partial T_1} e^{i\phi} + 2i^2 \omega_0 \frac{1}{2} a e^{i\phi} \left(\frac{\partial \phi}{\partial T_1} \right) + 3 \frac{1}{4} a^2 e^{2i\phi} \cdot \frac{1}{2} a e^{-i\phi} = 0 \\
 & \Rightarrow i\omega_0 \left(\frac{\partial a}{\partial T_1} \right) e^{i\phi} - \omega_0 a e^{i\phi} \left(\frac{\partial \phi}{\partial T_1} \right) + \frac{3}{8} a^3 e^{i\phi} = 0 \\
 & \Rightarrow i\omega_0 \left(\frac{\partial a}{\partial T_1} \right) - \omega_0 a \left(\frac{\partial \phi}{\partial T_1} \right) + \frac{3}{8} a^3 = 0 \\
 & \Rightarrow \boxed{\omega_0 \left(\frac{\partial \phi}{\partial T_1} \right) = \frac{3}{8} a^2} \rightarrow \text{Real part} \quad \boxed{\frac{\partial a}{\partial T_1} = 0} \Rightarrow a = a(T_2) \\
 & \quad \quad \quad \downarrow \\
 & \quad \quad \quad \text{Im. part}
 \end{aligned}$$

Recall that in the previous order, we had an equation which look like this. This was obtained by setting a resonant forcing term on the right hand side at that order to 0. Note that this is just D^{-1} of A . ΔA by ΔT^{-1} by definition is D^{-1} of A . So, I am just going to take this expression. Next replace D^{-1} of A from this and substitute it in the equation that I have.

This will just make it easier for me to do the algebra. So, from that equation that I just showed you, we obtained. This is obtained from the equation $\frac{d}{dt} \left(\frac{1}{2} m \dot{r}^2 \right) = 0$ by $\frac{d}{dt} \left(\frac{1}{2} m \dot{r}^2 \right) = 0$. This was the equation which allowed us to determine A , capital A .

So, this is just D_1 of A . So, I am just going to replace the expression for D_1 of A . From this into there. If I do that and simplify a little bit, then you can easily show that you get this expression twice $i\omega_0 D_1$ of B e to the power $i\omega_0 T_0$ that term remains

untouched. And then, if you replace D_1 of A with that term, you get 27 by $8\omega_0^2$ square, we get A^4 and then, we get $A\bar{e}$ to the power $3i\omega_0 T_0$. So, this is my expression for term 1. So, I will put term 1 here and put a bracket around it ok.

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$$\begin{aligned}
 \textcircled{2} &: -2D_0D_2u_0 \\
 &= -2i\omega_0(D_2A)e^{i\omega_0T_0} + \text{c.c.} \\
 \textcircled{3} &: -D_1^2u_0 \quad \boxed{D_1\bar{A} = \overline{D_1A}} \\
 &= \frac{9}{4} \frac{A^3\bar{A}^2}{\omega_0^2} e^{i\omega_0T_0} + \text{c.c.} \\
 \textcircled{4} &: -3u_0^2u_1 \\
 &= -3\left(Ae^{i\omega_0T_0} + \bar{A}e^{-i\omega_0T_0}\right)^2 \\
 &\quad \times \left(\underbrace{Be^{i\omega_0T_0}}_{\uparrow} + \underbrace{\frac{A^3}{8\omega_0^2}e^{3i\omega_0T_0}}_{\uparrow} + \underbrace{\bar{B}e^{-i\omega_0T_0}}_{\downarrow} + \underbrace{\frac{\bar{A}^3}{8\omega_0^2}e^{-3i\omega_0T_0}}_{\downarrow}\right) \\
 &= -3\left(\underbrace{A^2e^{2i\omega_0T_0}}_{\uparrow\uparrow} + \underbrace{\bar{A}^2e^{-2i\omega_0T_0}}_{\uparrow} + 2A\bar{A}\right) \times \left(\right)
 \end{aligned}$$

Now, let us go to term 2, term 2 is a simple term. I will leave that to work it out yourself and I will give you the final expression. It is just two lines of algebra and if you do that algebra, you will find just use the expression that we have for u_1 , u_1 we already have in my previous slides.

So, u_1 is there at the top of the slide at order 1, you have to just substitute u_1 . So, this is just twice $i\omega_0 u_1$ that is because of the derivative with respect to D_1 into D_2 into A , we can replace this D_2 into A later e to the power $i\omega_0 T_1$ plus complex conjugate. This is easy to do.

So, this is my second term on the right hand side. Let us work on the third term. The third term is minus $D_1^2 u_0$. This term is again easy to do. Again, like the first term, you will have to replace the expression for D_1 of A later on and somewhere you will have to use the fact that D_1 of A bar is equal to D_1 of A bar ok. So, if you have to if you use this, then you should be able to get this term ok.

So, this term actually turns out to be equal to after some a few lines of algebra, it just turns out to be equal to $9/4 A^3 A^* \omega_0^2 e^{i\omega_0 T_0}$ plus complex conjugate. There is some algebra here which I encourage you to try yourself, it is not difficult. One just has to be careful and do those lines of algebra to get this. This is my expression for the third term.

Then fourth term and this term, I will show it how to how do we do this? Because there is a product involved and one has to be a bit careful; otherwise, there is a possibility of doing mistakes. So, the fourth term is $u_0^2 u_1$. Now, let us work this term out. So, here there is a product of u_0^2 into u_1 and so, I have to be careful while doing this.

So, then, we will have u_0 by definition is $A e^{i\omega_0 T_0}$ and now, instead of writing a c.c., I am going to actually write down the complex conjugate term. So, the complex conjugate is just A^* into this and there is a square here multiplied by u_1 , u_1 had a complementary function part which was $B e^{i\omega_0 T_0}$ plus a particular integral which was $A^3/8 \omega_0^2 e^{3i\omega_0 T_0}$.

And here also like the previous term, I am going to write down its complex conjugate because that will help us do the multiplication and keep only the ones that we want to keep and the rest will go under complex conjugation. So, I will put B^* , so that is the complex conjugate of this term $e^{-i\omega_0 T_0}$ plus $A^3 e^{-3i\omega_0 T_0}$ ok.

And now, let us see what do we obtain from here ok. So, for this, let us square the first term. The first term is just $A^2 e^{2i\omega_0 T_0}$ plus $A^2 A^*$

e to the power minus twice $i\omega t$ plus twice the first term into the second term.

The exponentials will cancel and then, we will just have A, A^* . Multiplied by the second term exactly remains the same as what I have written earlier. Now, let us see how do we multiply these things term by term ok. So, we can immediately see that if I multiply. So, I am just going to indicate which terms I am multiplying by which. So, I will compare the multiplication of this term with the multiplication of that term because what I have put here is just the term on the top.

So, let us multiply the first term of this with the first term of that. That will give me minus 3 $A^2 B e^{3i\omega t}$, then I will multiply the first term of this with the third term of that ok. That will give me minus 3 $A^2 B^*$ into $e^{i\omega t}$ ok.

Notice the pattern that I am following. I only want e to the power of positive exponent; we do not want any multiplications, where it leads to e to the power a negative exponent. So, for example, if I multiply this term with this term, then you can immediately see that it will give me exponential minus $i\omega t$.

We do not want that term because we will get the complex conjugate of that term and that will be plus $i\omega t$. So, I am going to follow this procedure, where I multiply every term and make sure that I write only those terms whose exponent for the exponential is positive. I do not write any term, where the exponent is negative. I will put them all under c.c. ok.

So, I leave it to you to understand this. I am writing down the final expression. You can multiply each term in these two brackets. You will get whole lot of terms and you will see that I have not missed out anything which has a positive exponent. I am only not writing those terms, where the exponent is negative and I will put them under c.c. ok. So, you will see that

there is a, it always occurs in pairs. I only write one of the pair and I put the other one in c.c.
ok.

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$$\begin{aligned}
 4: & -3u_0^2 u_1 \\
 = & -3A^2 B e^{3i\omega_0 T_0} - 3A^2 \bar{B} e^{i\omega_0 T_0} - \frac{3A^5}{8\omega_0^2} e^{5i\omega_0 T_0} \\
 & - \frac{3A^3 (\bar{A})^2}{8\omega_0^2} e^{i\omega_0 T_0} - 6A\bar{A}B e^{i\omega_0 T_0} - \frac{6A^4 \bar{A}}{8\omega_0^2} e^{3i\omega_0 T_0}
 \end{aligned}$$

So, with that structure, the final answer here looks like. So, again, we are looking at the fourth term minus 3 u 0 square u 1 and this gives us quite a few terms which are minus 3 A square B e to the power thrice i omega naught T naught minus thrice A square B bar power i omega naught T naught.

So, of course, I have not combined all the terms. You can see that there are some terms which can be combined here; but I am just writing what we obtain after multiplication. We will continue this.