

**Thermodynamics for Biological Systems:
Classical and Statistical Aspects
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**Lecture - 7
Chemical Potential**

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Comparison of the first two terms of Eq. 2.10 with Eq. 2.5a

$$dU^T = \left(\frac{\partial U^T}{\partial S^T}\right)_{V^T, n_i} dS^T + \left(\frac{\partial U^T}{\partial V^T}\right)_{S^T, n_i} dV^T + \sum_{i=1}^{i=k} \left(\frac{\partial U^T}{\partial n_i}\right)_{V^T, S^T, n_j} dn_i \quad \text{Eq. 2.10}$$

$$dU^T = T dS^T - P dV^T \quad \text{Eq. 2.5a}$$


 $\left(\frac{\partial U^T}{\partial S^T}\right)_{V^T, n_i} = T$
 and
 $\left(\frac{\partial U^T}{\partial V^T}\right)_{S^T, n_i} = -P$

Welcome!

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Now, let us define

$$\left(\frac{\partial U^T}{\partial n_i}\right)_{V^T, S^T, n_j} \equiv \mu_i \quad \text{Eq. 2.11}$$

μ_i is called the chemical potential of species i

Thus, we can write Eq. 2.10 as

$$dU^T = T dS^T - P dV^T + \sum_i \mu_i dn_i \quad \text{Eq. 2.12}$$



Let us look at this quantity a little more closely. $\left(\frac{\partial U^T}{\partial n_i}\right)_{V^T, S^T, n_j}$. Let us define that quantity $\left(\frac{\partial U^T}{\partial n_i}\right)_{V^T, S^T, n_j}$ as something called μ_i ; where μ_i is actually called the chemical potential of species i .

$$\left(\frac{\partial U^T}{\partial n_i}\right)_{V^T, S^T, n_j} \equiv \mu_i$$

Note this term chemical potential. This is going to be a very important intrinsic variable in this entire course. So, please take a note of this. We will call this equation as 2.11. Therefore, in terms of, whatever we have defined just now, the chemical potential; we can write our equation 2.10 that we saw earlier as

$$dU^T = T dS^T - P dV^T + \sum_i \mu_i dn_i$$

We will call this equation 2.12. So, that was for internal energy.

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Using a similar argument, but for the total Gibbs free energy, G , we can say, that since,

$$G^T = f(T, P, n_1, n_2, n_3, \dots, n_k) \quad \text{Eq. 2.13}$$

$$dG^T = \left(\frac{\partial G^T}{\partial T}\right)_{P, n_i} dT + \left(\frac{\partial G^T}{\partial P}\right)_{T, n_i} dP + \sum_{i=1}^{i=k} \left(\frac{\partial G^T}{\partial n_i}\right)_{T, P, n_j} dn_i \quad \text{Eq. 2.14}$$

On comparison with Eq. 2.8a, $dG^T = -S^T dT + V^T dP$

$$dG^T = -S^T dT + V^T dP + \sum_i \mu_i dn_i \quad \text{Eq. 2.15}$$

$$\left(\frac{\partial G^T}{\partial n_i}\right)_{T, P, n_j} \equiv \mu_i \quad \text{Eq. 2.16}$$

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Let us follow a similar argument for the total Gibbs free energy, G . G^T , we know ... Rather, let us take G^T as a function of temperature, pressure and all the number of moles of the components; so, n_1, n_2 till n_k ;

$$G^T = f(T, P, n_1, n_2, n_3, \dots, n_k)$$

equation 2.13. Therefore, the total differential G^T ; dG^T can be written as

$$dG^T = \left(\frac{\partial G^T}{\partial T}\right)_{P, n_i} dT + \left(\frac{\partial G^T}{\partial P}\right)_{T, n_i} dP + \sum_{i=1}^{i=k} \left(\frac{\partial G^T}{\partial n_i}\right)_{T, P, n_j} dn_i$$

We will call this equation 2.14.

If we compare this 2.14 with 2.8a, which for pure component system or a single component system

$$dG^T = -S^T dT + V^T dP$$

We can write dG^T equals minus $S^T dT$ plus $V^T dP$ plus the remaining term, which comes in for the multi-component case, since, we have for a multi-component case, G^T to be a function of all other components also ... the amounts of components also. So, we have this extra term here. Summation of $\mu_i dn_i$. We will call this equation 2.15. Let us define

$$\left(\frac{\partial G^T}{\partial n_i}\right)_{T, P, n_j} \equiv \mu_i$$

as μ_i itself. Somebody can ask me this question now. You had earlier said that, dG^T by dn_i at certain variables being constant S^T , V^T , and n_i s being constant, n_j s being constant, you had defined as μ_i , whereas now you are defining dG^T by dn_i at constant T , P and n_j is μ_i . How can you do that?


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Note that we have used the same term, μ_i , to denote two different quantities in Eqs. 2.11 and 2.16. Let us justify that

Let us begin with Eq. 2.12

$$dU^T = T dS^T - P dV^T + \sum_i \mu_i dn_i \quad \text{Eq. 2.12}$$

Let us add $d(PV^T - TS^T)$ to both sides of Eq. 2.12

$$dU^T + d(PV^T - TS^T) = T dS^T - P dV^T + d(PV^T - TS^T) + \sum_i \mu_i dn_i$$


Let us now go ahead and show that, these two quantities are actually the same. To justify, let us begin with equation 2.12 where we first defined our μ_i . We started with

$$dU^T = T dS^T - P dV^T + \sum_i \mu_i dn_i$$

Let us add


$$d(PV^T - TS^T)$$

to both sides of equation 2.12. If we do that,

$$dU^T + d(PV^T - TS^T) = T dS^T - P dV^T + d(PV^T - TS^T) + \sum_i \mu_i dn_i$$

Note that, we can combine these together. They are additive. The differentials are additive. You can visualize the differential in a graph or in two dimensional graph or a three dimensional graph. Then, you can understand why this is additive.

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$$d(U^T + PV^T - TS^T) = \cancel{T dS^T} - \cancel{P dV^T}$$

$$+ \cancel{P dV^T} + V^T dP - \cancel{T dS^T} - S^T dT + \sum_i \mu_i dn_i$$

$$dG^T = -S^T dT + V^T dP + \sum_i \mu_i dn_i$$

Same as earlier
 μ_i in Eqs. 2.11 and 2.16 are the same

So, if we recognize that this is additive and you can expand the term on this side using chain rule. Let us do that; therefore,

$$d(U^T + P V^T - T S^T) = T dS^T - P dV^T + P dV^T + V^T dP - T dS^T - S^T dT + \sum_i \mu_i dn_i$$

We can cancel $T d S T$ and minus $T d S T$. We can also cancel plus $P d V T$ and minus $P d V T$. And therefore,

$$dG^T = -S^T dT + V^T dP + \sum_i \mu_i dn_i$$

This was the same μ_i that we started with earlier, we did not do anything to it. Therefore, the μ_i 's in equation 2.11, which was for the internal energy $d U T$ and equation 2.16 for the Gibbs free energy $d G T$ are the same. Although we had used different ways to define the chemical potential, they all turn out to be the same quantity.


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$$\left(\frac{\partial G^T}{\partial n_i}\right)_{T,P,n_j} \equiv \mu_i \quad \text{Eq. 2.16}$$

The above is a better preferred definition of μ_i
 because of experimental relevance

T, P and n_j can be easily maintained constant

*Although the chemical potential, μ_i , was introduced in
 the context of a multi-component system, it is an
 intrinsic thermodynamic property. Thus, it is an
 equally valid concept for a pure component too.*


 NPTEL

Let us look at this equation 2.16 a little closer. Is a very useful equation.

$$\left(\frac{\partial G^T}{\partial n_i}\right)_{T, P, n_j} \equiv \mu_i$$

Note that, we had defined we have defined μ_i in the context of a multi-component system. That is purely didactic; that is purely to facilitate a certain flow in the development of the subject. The chemical potential itself is an intrinsic quantity for a particular system. Therefore, whether the system is multi-component or whether it is a single component or whatever it is, whatever other combinations you may have, the concept of chemical potential is still valid for any system that you take. To put a little more formally, although the chemical potential μ_i was introduced in the context of a multi-component system, it is an intrinsic thermodynamic property. Thus it is an equally valid concept for a pure component also.

Let us look at this a little closer now. Look at the quantities that are being held constant in this definition; ∂G^T by ∂n_i at constant temperature, pressure and all other n_j 's different from n_i , which is taken for the differentiation. These are very easy to maintain constant in an experimental situation. Therefore this definition of chemical potential has significant experimental relevance. ... Because of that, although you could define chemical potential as ∂U^T by ∂n_i at certain variables constant ∂A^T by ∂n_i at certain other variables constant and so on and so forth, this definition is most commonly used.


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By extension, we can define chemical potential in terms of H and A also. Thus, for a multi-component, single phase system:

$$dU^T = T dS^T - P dV^T + \sum_i \mu_i dn_i \quad \text{Eq. 2.12}$$

$$dG^T = -S^T dT + V^T dP + \sum_i \mu_i dn_i \quad \text{Eq. 2.15}$$

$$dH^T = T dS^T + V^T dP + \sum_i \mu_i dn_i \quad \text{Eq. 2.17}$$

$$dA^T = -S^T dT - P dV^T + \sum_i \mu_i dn_i \quad \text{Eq. 2.18}$$


I have already mentioned this, but let me complete the definitions here. So, by extension, as mentioned earlier, we can define chemical potential in terms of enthalpy, H, or the Helmholtz free energy, A. Therefore, for a multi-component, single phase system, dU^T equals $T dS^T$ minus $P dV^T$ plus summation of $\mu_i dn_i$, equation 2.12. We just compiling all the equations together, we will have it at one place for easier reference later.

$$dU^T = T dS^T - P dV^T + \sum_i \mu_i dn_i$$

$$dG^T = -S^T dT + V^T dP + \sum_i \mu_i dn_i$$

$$dH^T = T dS^T + V^T dP + \sum_i \mu_i dn_i$$

$$dA^T = -S^T dT - P dV^T + \sum_i \mu_i dn_i$$

dG^T equals minus $S^T dT$ plus $V^T dP$ plus summation of $\mu_i dn_i$ over all i , equation 2.15; dH^T equals $T dS^T$ plus $V^T dP$ plus summation of $\mu_i dn_i$, equation 2.17 and dA^T equals minus $S^T dT$ minus $P dV^T$ plus summation of $\mu_i dn_i$, we will call this equation 2.18.

What I would like you to do, is to actually sit down and verify the same way that we did, that these two μ_i 's were the same. It will be good if you verify and convince yourself that these two

mu i's are also the same. And, they are the same as the other mu i's that are here. That way, there will be no doubts in your mind.

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$$dU^T = T dS^T - P dV^T + \sum_i \mu_i dn_i$$

$$\left(\frac{\partial U^T}{\partial S^T}\right)_{V^T, n_i} \quad \left(\frac{\partial U^T}{\partial V^T}\right)_{S^T, n_i} \quad \left(\frac{\partial U^T}{\partial n_i}\right)_{V^T, S^T, n_j}$$

$$dH^T = T dS^T + V^T dP + \sum_i \mu_i dn_i$$

$$\left(\frac{\partial H^T}{\partial S^T}\right)_{P, n_i} \quad \left(\frac{\partial H^T}{\partial P}\right)_{S^T, n_i} \quad \left(\frac{\partial H^T}{\partial n_i}\right)_{S^T, P, n_j}$$

$$dA^T = -S^T dT - P dV^T + \sum_i \mu_i dn_i$$

$$\left(\frac{\partial A^T}{\partial T}\right)_{V^T, n_i} \quad \left(\frac{\partial A^T}{\partial V^T}\right)_{T, n_i} \quad \left(\frac{\partial A^T}{\partial n_i}\right)_{T, V^T, n_j}$$

$$dG^T = -S^T dT + V^T dP + \sum_i \mu_i dn_i$$

$$\left(\frac{\partial G^T}{\partial T}\right)_{P, n_i} \quad \left(\frac{\partial G^T}{\partial P}\right)_{T, n_i} \quad \left(\frac{\partial G^T}{\partial n_i}\right)_{T, P, n_j}$$

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These are the four equations that we have just derived. I would like to point out some interesting features here. These four equations dU^T equals $T dS^T$ minus $P dV^T$ plus $\mu_i dn_i$, dH^T equals $T dS^T$ plus $V^T dP$ plus $\mu_i dn_i$, dA^T equals minus $S^T dT$ minus $P dV^T$ plus $\mu_i dn_i$ and dG^T equals minus $S^T dT$ plus $V^T dP$ plus $\mu_i dn_i$... summation.

Let us look at it in terms of writing a total differential, in terms of its partial differentials.

$$dU^T = T dS^T - P dV^T + \sum_i \mu_i dn_i$$

$$\left(\frac{\partial U^T}{\partial S^T}\right)_{V^T, n_i} \quad \left(\frac{\partial U^T}{\partial V^T}\right)_{S^T, n_i} \quad \left(\frac{\partial U^T}{\partial n_i}\right)_{V^T, S^T, n_j}$$

This term, as you can recall can be written as $\text{d}U^T$ by $\text{d}S^T$ at constant V^T, n_i ; both these held constant. And, this term can be written as $\text{d}H^T$ by $\text{d}S^T$ at constant P and n_i . Similarly, this term P can be written as $\text{d}U^T$ by $\text{d}V^T$ at constant S^T and n_i .

The same P here, which appears in this other equation can be represented in terms of the Helmholtz free energy as $\text{d}A^T$ by $\text{d}V^T$ at constant T and n_i .

Now, you get the scheme with which we are going to do things. Minus S^T is dU^T by dS^T at constant V^T and n_i . Let us look for another S^T . It is here; $S^T \dots$ minus S^T equals dG^T by dS^T at constant P and n_i . Similarly, this V^T can be written as dH^T by dP at constant S^T and n_i . And, V^T which appears here can be written as dG^T by dP at constant T and n_i . And, as we have already seen, the chemical potentials can be written in terms of either the partial of U^T with n_i at constant S^T , V^T and all other n_j 's here. In terms of enthalpy, you could write μ_i as dH^T by dn_i at constant S^T , V^T and all other n_j 's. dA^T by dn_i at constant T , V and all other n_j 's. And, dG^T at constant ... dG^T by dn_i at constant T , P and all other n_j 's.

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$$T = \left(\frac{\partial U^T}{\partial S^T} \right)_{V^T, n_i} = \left(\frac{\partial H^T}{\partial S^T} \right)_{P, n_i} \quad \text{Eq. 2.19}$$

$$P = - \left(\frac{\partial U^T}{\partial V^T} \right)_{S^T, n_i} = - \left(\frac{\partial A^T}{\partial V^T} \right)_{T, n_i} \quad \text{Eq. 2.20}$$

$$S^T = - \left(\frac{\partial G^T}{\partial T} \right)_{P, n_i} = - \left(\frac{\partial A^T}{\partial T} \right)_{V^T, n_i} \quad \text{Eq. 2.21}$$

If we put this all together in one place,

$$T = \left(\frac{\partial U^T}{\partial S^T} \right)_{V^T, n_i} = \left(\frac{\partial H^T}{\partial S^T} \right)_{P, n_i}$$

equation 2.19

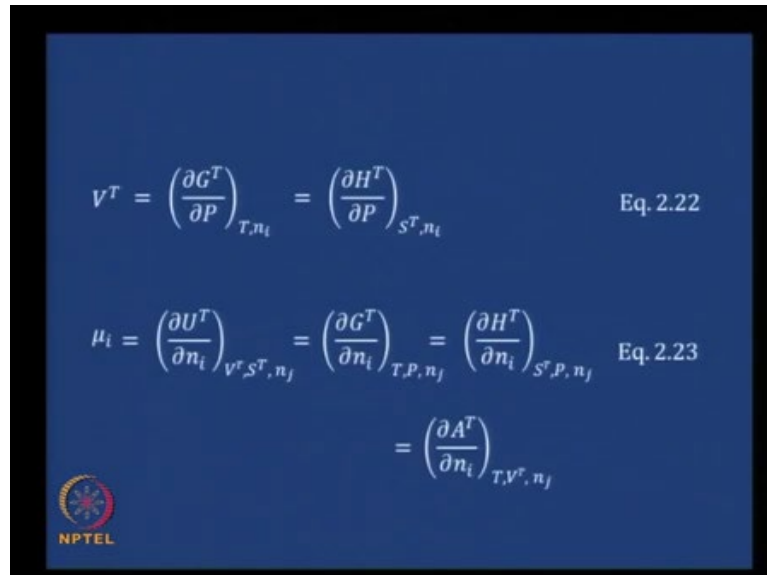
$$P = - \left(\frac{\partial U^T}{\partial V^T} \right)_{S^T, n_i} = - \left(\frac{\partial A^T}{\partial V^T} \right)_{T, n_i}$$

equation 2.20

$$S^T = -\left(\frac{\partial G^T}{\partial T}\right)_{P, n_i} = -\left(\frac{\partial A^T}{\partial T}\right)_{V^T, n_i}$$

equation 2.21

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The slide contains two equations on a blue background with a black border. Equation 2.22 is $V^T = \left(\frac{\partial G^T}{\partial P}\right)_{T, n_i} = \left(\frac{\partial H^T}{\partial P}\right)_{S^T, n_i}$. Equation 2.23 is $\mu_i = \left(\frac{\partial U^T}{\partial n_i}\right)_{V^T, S^T, n_j} = \left(\frac{\partial G^T}{\partial n_i}\right)_{T, P, n_j} = \left(\frac{\partial H^T}{\partial n_i}\right)_{S^T, P, n_j} = \left(\frac{\partial A^T}{\partial n_i}\right)_{T, V^T, n_j}$. The NPTEL logo is in the bottom left corner.

$$V^T = \left(\frac{\partial G^T}{\partial P}\right)_{T, n_i} = \left(\frac{\partial H^T}{\partial P}\right)_{S^T, n_i}$$

equation 2. 22.

And, all the four ways of representing the chemical potential as

$$\mu_i = \left(\frac{\partial U^T}{\partial n_i}\right)_{V^T, S^T, n_j} = \left(\frac{\partial G^T}{\partial n_i}\right)_{T, P, n_j} = \left(\frac{\partial H^T}{\partial n_i}\right)_{S^T, P, n_j} = \left(\frac{\partial A^T}{\partial n_i}\right)_{T, V^T, n_j}$$

Let me call this equation 2.23.

These are very useful equations as we will see.

See you in the next class.