

**Satellite Attitude Dynamics And Control**  
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**Lecture - 23**  
**Rigid Body Dynamics (Contd)**

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Lecture - 23

Rigid Body Dynamics

Torque free rigid body dynamics

$\dot{\omega}_3 = 0 \Rightarrow \omega_3 = \text{a constant}$

$\omega_3 = \dot{\phi} + \dot{\psi} \cos \theta = \text{a constant} = n \Rightarrow \omega_3 = n$

$I_1 = I_2 = I$   
 $I_3 = I_0$  } Assumption

$I_1 \dot{\omega}_1 - (I_2 - I_3) \omega_2 \omega_3 = M_1 = 0$

$\dot{\omega}_1 = \frac{(I_0 - I)}{I} \omega_2 \omega_3$

$\dot{\omega}_2 = \frac{(I_0 - I)}{I} \omega_3 \omega_1$

$\dot{\omega}_3 = \frac{(I - I)}{I_0} \omega_1 \omega_2 = 0 \Rightarrow \omega_3 = \text{a constant} = n$

$\ddot{\omega}_1 = \frac{(I - I_0)}{I} \dot{\omega}_2 \omega_3 = \frac{(I - I_0)}{I} \frac{(I_0 - I)}{I} \omega_2^2 \omega_1$

$\ddot{\omega}_1 = -\left(\frac{I_0 - I}{I}\right)^2 n^2 \omega_1 = -\lambda^2 \omega_1$

where  $\lambda = \left(\frac{I_0 - I}{I}\right) n$

Welcome to the 23rd lecture. We will continue with Rigid Body Dynamics, we have been discussing in the last lecture ok. So, recalling that for the rigid body dynamics, we have used the Euler's equation and especially for this case where the it is torque free. So, what we have got we had three equations.

So, from the last equation, we got that the omega 3 dot, this will turn out to be 0 ok. And this implies omega 3, this is a constant. Moreover, if you remember that omega 3 already using the kinematics equation, we have written this as phi dot plus psi dot cos theta ok. So, this implies this quantity is going to be a constant.

Beside this we have derived the equation of motion in this format, where we have assumed that. So, whatever we have derived this was done under the assumption that I 1 equal to I 2 equal to I and I 3 equal to I 0, this is the assumption ok. So, under this condition we have got I times omega 1 double dot plus and let us write this omega 3 constant this equal to n. This is a means this omega 3 this equal to n, implying this is omega 3 equal to n.

So, under that condition if we look back into the previous equation, so it can be written as  $n$  times  $n$  square times let me write it in another way, we will start from the scratch and then write it. So, we had  $I_1$  times  $\omega_1$  dot, this is the angular acceleration plus and will follow the sequence. So,  $I_1$  then we have  $I_2$  minus  $I_3$  times  $\omega_2$  times  $\omega_3$  and on this side the  $M_1$  which we have set it to 0, because it was torque free (Refer Time: 03:15) just I am recalling the previous thing.

And from there we have written this as  $\omega_1$  dot minus or this equal to  $I_2$ , we are writing as  $I$  and this as  $I_0$  and  $I_1$  then becomes  $I \omega_2 \omega_3$ . And similarly,  $\omega_2$  dot we have written as following the same kind of rule. So, here these are starting with 2, so it was here  $I_0$ , this is  $I_3$  basically  $I_3, 2_3$  and then 1, so this is  $I$  and  $\omega_2$ , after this is 3, and then the 1 will come. And obviously here the  $I_2$  is there, so  $I_2$  will come here in the denominator which we are writing as  $I$ , so this is what we have written last time.

And  $\omega_3$  dot, this was taken from the Euler's equation. So, in that case  $\omega_3$  there after the  $I_1$  and  $I_2$  comes. So,  $I_1$  equal to  $I$  and out  $I_2$  equal to also  $I$ , so this will cancel out. And then we have 3 after this 1 and this is 2 divided by  $I_3$ , which is  $I_0$  so, this is set to 0. And from there this is what we have recovered ok. So, this implies  $\omega_3$ , this is a constant and which we are writing as  $n$  ok.

And we double different we differentiated once this quantity and  $\omega_3$  is already constant. So, once we difference here, this remains constant. And from there we have got this as  $\omega_1$  double dot, this equal to  $I$  minus  $I_0$  divided by  $I$  times  $\omega_2$  dot times  $\omega_3$ . And  $\omega_2$  dot, we are inserting from this place, so this becomes  $I$  minus  $I_0$  divided by  $I$  and  $I_0$  minus  $I$  divided by  $I$  times  $\omega_3 \omega_1$ . And  $\omega_3$  from this place, because we have replaced to, so this becomes  $\omega_3$  square and then  $\omega_1$ .

So, this way we are getting  $\omega_1$  double dot, this equal to now we will write in terms of  $I_0$  minus  $I$ . So, this is  $I_0$  minus  $I$  divided by  $I$  whole square, then  $\omega_3$  this is written as  $n$  square  $\omega_1$  and this quantity we have written as and here there will be a minus sign ok. So, this is minus  $\lambda$  square  $\omega_1$ , where  $\lambda$  we have written as  $I_0$  minus  $I$  divided by  $I$  times  $n$  ok. So, the proceeding in the same way and we got  $\omega_2$  double dot equal to minus  $\lambda$  square  $\omega_2$ .

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$$\begin{cases} \ddot{w}_2 + \lambda^2 w_2 = 0 \\ \ddot{w}_1 + \lambda^2 w_1 = 0 \end{cases} \Rightarrow \begin{cases} w_1 = A \cos \lambda t + B \sin \lambda t \\ w_2 = B \cos \lambda t - A \sin \lambda t \end{cases}$$

$$\Rightarrow w_1^2 + w_2^2 = A^2 \cos^2 \lambda t + B^2 \sin^2 \lambda t + 2AB \cos \lambda t \sin \lambda t + B^2 \cos^2 \lambda t + A^2 \sin^2 \lambda t - 2AB \cos \lambda t \sin \lambda t$$

$$w_1^2 + w_2^2 = A^2 + B^2 = \text{constant}$$

$$\begin{array}{l|l} I_1 \dot{w}_1 - w_2 w_3 (I_2 - I_3) = 0 & I_2 \dot{w}_2 - w_3 w_1 (I_3 - I_1) = 0 \quad \left( \begin{array}{l} I_1 = I_2 = I \\ I_3 = I_0 \end{array} \right) \\ \text{multiply this by } w_1 & \text{multiply this by } w_2 \end{array}$$

and add.

$$I_1 w_1 \dot{w}_1 + I_2 w_2 \dot{w}_2 - w_1 w_2 w_3 (I_2 - I_3) - w_1 w_2 w_3 (I_3 - I_1) = 0$$

$$I (w_1 \dot{w}_1 + w_2 \dot{w}_2) - w_1 w_2 w_3 \left[ \frac{I_2 - I_3}{3} + \frac{I_3 - I_1}{3} \right] = 0$$

$$w_1 \dot{w}_1 + w_2 \dot{w}_2 = 0 \Rightarrow w_1^2 + w_2^2 = \text{constant}$$

So, the second equation in the same way, we get as  $\omega_2$  double dot plus  $\lambda^2 \omega_2$  equal to 0. And the first equation we have  $\omega_1$  this equal to 0, and as I told you the solution to because this is a simple harmonic motion format the standard format.

So, we can write  $\omega_1$  this equal to  $A \cos \lambda t + B \sin \lambda t$ . And  $\omega_2$  we have written as because it is a  $\omega_2$  will differ from  $\omega_1$  by 90 degree, but it is a perpendicular to that. And therefore, in this case if we put here  $\phi$  by  $\pi/2$ , so at the phase difference of  $\phi$  by  $\pi/2$ , so this becomes  $B \cos \lambda t - A \sin \lambda t$  ok.

And if you a square and add, so this will be  $A^2 \cos^2 \lambda t + B^2 \sin^2 \lambda t$  from the previous one and the first one  $2AB \cos \lambda t \sin \lambda t$  and from this one we get  $B^2 \cos^2 \lambda t + A^2 \sin^2 \lambda t - 2AB \cos \lambda t \sin \lambda t$  ok. So, adding these two will cancel out plus minus signs are there and these two will add up.

So,  $A^2$  we can take it as common, so this becomes  $A^2 + B^2$  ok. So, this is a constant, because  $A$  and  $B$  they are constant. So,  $\omega_1^2 + \omega_2^2$  square, it turns out to be a constant. So, this is not the only way of doing this, this can be grouped in other way also. So, for this what we need to do that once we are writing this equation  $I_1 \dot{\omega}_1 - I_2 \omega_2 \omega_3$  and on the right hand side it is 0, so multiply this  $\omega_1$  ok.

Similarly, we have  $\mathbf{i} \cdot 2\boldsymbol{\omega}_2 - \boldsymbol{\omega}_3 \times \boldsymbol{\omega}_1$   $\mathbf{i} \cdot 3 - \mathbf{i} \cdot 1$  ok, on the right hand side there is no torque, so this is 0. So, multiply this by  $\boldsymbol{\omega}_2$  and add. So, if we do this, so this will be  $\mathbf{i} \cdot 1$  times  $\boldsymbol{\omega}_1$  times  $\boldsymbol{\omega}_1 \cdot$  and plus  $\mathbf{i} \cdot 2$  times  $\boldsymbol{\omega}_2$  times  $\boldsymbol{\omega}_2 \cdot$ . And here this term will be  $\boldsymbol{\omega}_1$  times  $\boldsymbol{\omega}_2$   $\boldsymbol{\omega}_3$  and now  $\mathbf{i} \cdot 2$   $\mathbf{i} \cdot 3$  are there, so we need to replace them, this is  $\mathbf{i} \cdot 2 - \mathbf{i} \cdot 3$ . Because, we have already assumed  $\mathbf{i} \cdot 1$  equal to  $\mathbf{i} \cdot 2$  this is the case of symmetry, we are assuming. If we do not assume case of symmetry, the result will be different. So,  $\mathbf{i} \cdot 2 - \mathbf{i} \cdot 3$  and then from this place again  $\boldsymbol{\omega}_1 \boldsymbol{\omega}_2 \boldsymbol{\omega}_3$   $\mathbf{i} \cdot 3 - \mathbf{i} \cdot 1$  ok, and this quantity is going to be 0.

And therefore, if we replace this by  $\mathbf{i}$ , so  $\mathbf{i}$  can be taken as common  $\mathbf{i} \cdot 1$  time  $\boldsymbol{\omega}_1$  times  $\boldsymbol{\omega}_1 \cdot$   $\boldsymbol{\omega}_2$  times  $\boldsymbol{\omega}_2 \cdot$  minus here  $\boldsymbol{\omega}_1 \boldsymbol{\omega}_2 \boldsymbol{\omega}_3$  this we can take as common and here this is  $\mathbf{i} \cdot 2$ . So,  $\mathbf{i} \cdot 2 - \mathbf{i} \cdot 3$  and plus  $\mathbf{i} \cdot 3 - \mathbf{i} \cdot 1$  ok, this equal to 0. So, you can see that these and this they cancel out and  $\mathbf{i} \cdot 2$  and  $\mathbf{i} \cdot 2$  and  $\mathbf{i} \cdot 1$  they being equal so, this gets reduced to 0.

And therefore, we can write  $\boldsymbol{\omega}_1$  times  $\boldsymbol{\omega}_1 \cdot$  plus  $\boldsymbol{\omega}_2$  times  $\boldsymbol{\omega}_2 \cdot$ , this equal to 0 ok. And this implies that if we integrate it with respect to  $t$ , so this will be  $\boldsymbol{\omega}^2$  square plus  $\boldsymbol{\omega}_2$  square, this is a constant. And this is what exactly we have got here also in this place ok. So, both way you can look into so the results are consistent in both the ways.

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The image shows a handwritten derivation and diagrams illustrating the relationship between angular velocities and their magnitudes. The derivation starts with the vector equation  $\vec{\omega} = \omega_1 \hat{e}_1 + \omega_2 \hat{e}_2 + \omega_3 \hat{e}_3$  and the magnitude equation  $\omega^2 = \omega_1^2 + \omega_2^2 + \omega_3^2 = a \text{const.}$ . It then shows that  $\omega_3 = a \text{const.}$  and  $\omega = a \text{constant}$ , where  $\omega$  is the magnitude. The diagrams show vectors  $\vec{\omega}_1, \vec{\omega}_2, \vec{\omega}_3$  and their resultant  $\vec{\omega}$  in a 3D coordinate system with unit vectors  $\hat{e}_1, \hat{e}_2, \hat{e}_3$ . The resultant vector  $\vec{\omega}$  is shown to be coplanar with the vectors  $\vec{\omega}_1$  and  $\vec{\omega}_2$ . The final result is  $\vec{\omega} = \omega_1 \hat{e}_1 + \omega_2 \hat{e}_2 + \omega_3 \hat{e}_3$  and  $\vec{h} = I_1 \omega_1 \hat{e}_1 + I_2 \omega_2 \hat{e}_2 + I_3 \omega_3 \hat{e}_3 = I(\omega_1 \hat{e}_1 + \omega_2 \hat{e}_2) + I_3 \omega_3 \hat{e}_3$ . The text states that  $\vec{h}$  and  $\vec{\omega}$  are coplanar.

So, what we get that  $\omega_1^2 + \omega_2^2$ , this is a constant. And already we have observed that  $\omega_3$ , this is also a constant. And therefore, this implies that  $\omega^2 = \omega_2^2 + \omega_3^2$ , this will be a constant plus a constant. And this implies the  $\omega$  is also a constant.

So, therefore  $\omega$  turns out to be a constant  $\omega_1^2 + \omega_2^2$ , this is a constant ok. Now, this is constant in magnitude, but not in direction. As already we have seen that  $\omega_1$ , this will oscillate at the frequency of  $\lambda$  and  $\omega_2$  also this oscillates at the frequency of  $\lambda$ . Therefore, we can draw the conclusion that  $\omega$  will also oscillate with the frequency of  $\lambda$ .

Now, we will look into this picture  $e_1, e_2$ , and  $e_3$ , this is the body frame. And let us say this is the vector  $\omega$ , if we take the projection of this. So, this component this is  $\omega_2$ , and this component is  $\omega_1$ . So,  $\omega_1^2 + \omega_2^2$ , this equal to we can write this as  $\omega^2$  which is a constant.

So, this implies that instead of writing  $\omega_1$  and  $\omega_2$ , we can replace in terms of  $\omega^2$  so,  $\omega^2$  this is a constant. So,  $\omega$  it lies in the plane  $e_1$  and  $e_2$  to remember that  $\omega$  we have described in terms of the components in the body frame, which are  $\omega_1$  and  $\omega_2$ . So, your  $\omega$ , this is  $\omega$  vector. And if you want to show it in the vector form, so we need to put a arrow of this, but I am not showing like that only the components, I have shown here in this place.

Now, already we have written this like  $\omega_1 e_1 + \omega_2 e_2 + \omega_3 e_3$ . And if we remember this  $h$  vector we are writing as  $I_1 \omega_1 e_1 + I_2 \omega_2 e_2 + I_3 \omega_3 e_3$ . So,  $I_1$  and  $I_2$  though are they are equal, so are we can take it outside. And this remains as  $\omega_2 e_2 + \omega_3 e_3$  plus and this  $I_3$  we are writing as  $I_0$  so, this  $\omega_3 e_3$ ,  $\omega_3$  is nothing but  $n$ .

So,  $h$  this can be written as if you look here in this place, this part if we write this as  $\omega$  like this ok. It was a magnitude of this if we write here,  $\omega^2 = \omega_1^2 + \omega_2^2$ , you can see that the magnitude of this vector will be  $\sqrt{\omega_1^2 + \omega_2^2}$  ok. This is a constant as we have written earlier, this is a constant.

So, from this place what we can try it that this is  $I \omega t$  ok, this is a magnitude. And let us consider that unit vector here in this direction is  $\hat{e}_t$  cap. So, this can be written as  $\hat{e}_t$  cap and plus  $I \omega$  times  $\omega_3 \hat{e}_3$  cap ok. And this part  $\omega_3$  also can be rewritten as  $\omega_1 \hat{e}_1$  cap  $\omega_2 \hat{e}_2$  cap is equal to  $\omega t$  times  $\hat{e}_t$  cap and this part  $\omega_3$  times  $\hat{e}_3$  cap.

So, now let us look into the figure here. Like if we write here in this direction say, this is  $\omega t$  ok, so  $\omega t$  is here. And this is  $I \omega t$  means, you are magnifying it ok. So, one is this is  $\omega t$  and further magnitude let us say this is  $I \omega t$ , so from here to here. And then multiplied by this vector  $\hat{e}_t$ , so it will get converted into will multiply here by  $\hat{e}_t$  cap and the here also as  $\hat{e}_t$  cap to indicate these are the vectors ok.

Similarly, we have here in this place this is  $\omega_3$  and these two are perpendicular to each other, if you see the  $\omega_3$  is along this direction ok. So, we need to take perpendicular one, so for the perpendicular one, we have let us say this is  $\omega_3$ . So,  $\omega_3$  times  $\hat{e}_3$  cap and the another vector which is  $I \omega$  times  $\omega_3$ . So,  $I \omega$  times  $\omega_3$  let us make it here in this place. So, from this place it runs from this place and this also runs from this place from here to here and from this place to this place. So,  $I \omega$  times  $\omega_3$ .

So, if we make a linear combination of this, so how it will appear? Vector addition it will give you for this blue one, this is the vector here ok. And this vector depending on this magnitude how much this is and how much this one is so you will get another vector, which may look like this from here to here and this place to this place. So, then this is your this vector and this vector is this ok.

So, what it implies that  $\mathbf{h}$  and  $\omega$  are co-planar ok, this is what it implies, because  $\hat{e}_3$  and it depends on  $\hat{e}_3$  and  $\hat{e}_t$  and it is a just different linearly ok, you can see this component  $\omega t$  here this  $\omega t$ . So, here it is  $I \omega t$  so, just referring in magnitude, while they are in the same plane ok. And therefore, they are bound to be co-planar, so this implies that  $\mathbf{h}$  and  $\omega$  they are co-planar ok.

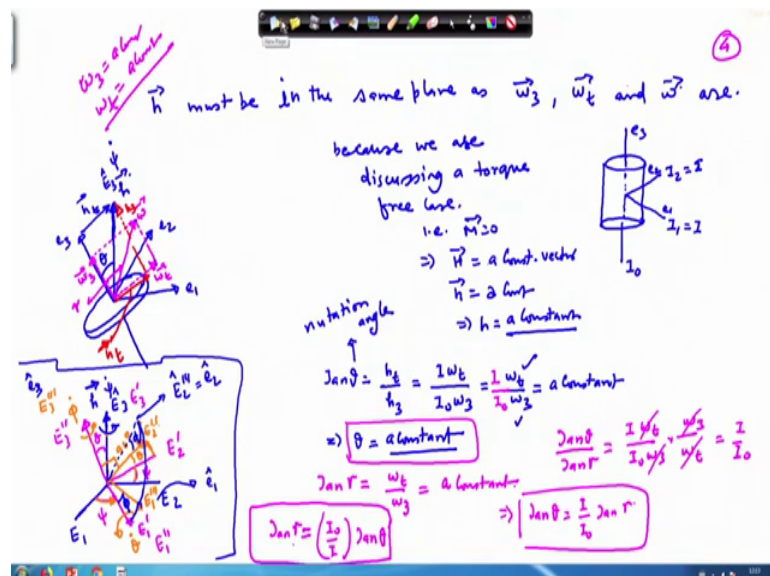
So, what we can conclude here that  $\mathbf{h}$   $\omega$ , then  $\omega t$  and  $\omega_3$ , they are co-planar. Now, already we have observed that vector  $\omega_1$  and  $\omega_2$ , it is a oscillating at the rate of it is a rotating at the rate of this frequency  $\lambda$  read in per second ok.

And therefore,  $\omega_1^2 + \omega_2^2$ , this we are writing as  $\omega^2$  ok. So, therefore this will also rotate at the frequency of this  $\lambda$  ok, but from these two equations, these vectors need to be all the time in the same plane. So, what does it imply that if  $\omega$  is rotating ok,  $\omega$  is rotating, because  $\omega_1$  and  $\omega_2$ , they are rotating as a consequence of this  $\omega$  is rotating and  $\omega_3$  is always perpendicular to this.

And therefore, the combination of this which is the  $\omega$  vector;  $\omega$  vector will also rotate at the this angle of frequency of  $\lambda$ . So, this implies that this is also rotating, this is also rotating ok, this is also rotating and therefore  $h$  must also be a rotating, because  $h$  has to be in the same plane ok. Here you can see that this is the  $h$  defined here ok, this is the  $h$  vector and they are bound to lie in the same plane.

So, if this vector is rotating which is the  $\omega$  vector, here I should rather show it like this. And this one rather I should show like this ok. So, this is the  $\omega$  vector and this is the  $h$  vector. So, if your  $h$  vector is rotating ok, so the if your  $\omega$  is rotating, so that also implies that it  $h$  vector also must rotate means, the  $h$  vector lying in the same plane. We now go to the next page.

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So, this implies  $h$  must be in the same plane as  $\omega_3$   $\omega_3$  vector,  $\omega$  vector, and  $\omega$  ok. So, let us consider the case here we have a disc or it may be a cylinder.

So, the case we have been discussing the torque free case, it can be represented either by cylinder or either by a disc.

As for the cylinder we know that along this axis, we have written this as  $I_0$  and the other two axis we have written as  $I_1$  equal to  $I$  and  $I_2$  equal to  $I$ , it is the moment of inertia along the other two axis  $e_3$ ,  $e_1$  and this is  $e_2$  ok. So, the same thing it applies in the case of the disc also.

So, what we assumed that let us assume that  $h$  is along this direction, because we are discussing a torque free case that is  $M$  equal to 0. So, this implies  $H$  will be a constant  $H$  equal to a constant vector or equally instead of capital notation, we have used here a small notation so,  $h$  is a constant and this implies there is a magnitude also remains here constant ok.

So, we have here  $e_1$  axis and this is the  $e_2$  axis, which is lying in the plane of this disc, and here we have this  $e_3$  axis. Some where your  $\omega$  vector is lying along this direction, so  $\omega t$  will lie here in this plane and it can be shown like this. So, this is your  $\omega t$  that implies that this component is  $\omega_3$  is along this direction as shown by this plane.

Similarly,  $h$  vector can be broken along these directions, so we will write the because they are all co-planar  $\omega t$ ,  $\omega h$  and  $\omega_3$ , they are co-planar. And therefore, the component of this  $h$  vector, it must lie in this plane itself. So, if they drop it, the somewhere shown by this red line, this is your  $h t$  ok and this component this will be your  $h_3$ . Obviously,  $h t$  is a combination of its component along the  $e_1$  direction and the  $e_2$  direction.

So, what exactly we have done here that given this initial reference frame  $E_1$ ,  $E_2$ ,  $E_3$ . And a body is there which is torque free and we are assuming that it is a angular momentum vector, it is along the  $E_3$  direction. So, this is also your happens to be your  $E_3$  direction or either we can show it like the  $h t$  cap. And we are assuming that the angular velocity vector which along this direction, so this simplifies the case you can assume it to be in other directions also no problem, but assuming this it definitely it simplifies the whole thing and as earlier we have done.



Let first we are giving rotation about this by angle  $\psi$ . So, this is  $\psi$  dot along this direction, so here  $\psi$  dot will be lying along this direction, then once it comes here in this place somewhere here. This angle being  $\psi$ , then you are rotating it by angle  $\theta$  here. So, as a consequence of this, this will rotate from this place to this place.

So, here you have  $E_1'$  as we have shown earlier,  $E_2'$  we have shown here, and so once this is rotated this is  $\psi$  ok and  $E$  lies along the same direction. And once we rotate by  $\theta$ , so this will also rotate and here we will have  $\theta$ . So,  $E_3'$  lies along this direction,  $E_1'$ ,  $E_2'$ ,  $E_3'$  and then  $E_3''$  along this direction,  $E_1''$  along this direction. So, once you are rotating about this, so this will rotate by this amount and here we have shown it like this.

So, once this rotates from this place to this place the corresponding thing, this angle is 90 degree, this is your  $\theta$ . First we have rotated by  $\psi$ , this rotates by  $\psi$  ok. Then we have rotated by  $\theta$ , so this rotates from this place to this place by  $\theta$ . And this rotates from this place to this place by  $\theta$  and then the final rotation you are giving about this line which is by  $\phi$ . So, as a consequence of this  $\phi$  dot is appearing here, and  $\theta$  dot is appearing here in this place ok.

So, in the final configuration your  $E_3'''$  is here,  $E_2'''$  is here and ok. And this will rotate once we rotate about this  $\phi$ , so this will go away from here ok. So, figure I am not showing here this will go away from this place, this will also go away from this place.

So, somewhere it will come near this point and this is your  $E_1'''$ ,  $E_1'$ , this is  $E_2$ , this is  $E_2'$ , next rotation we have given. So,  $E_2''$ ,  $E_2'''$ ,  $E_2''''$  and  $E_3''''$ , it will go out like this. Let me try we can add just here, it will go like this. So, this and this they lie in the same plane. Say here this part, this part they lie here in the same plane, figure is getting complicated, this is  $\phi$ . This will rotate from this position to this position, this is  $\phi$ .

So, instead of dealing with this picture, which is getting very complicated. What we showed that this is my initial configuration  $E_3$  vector is  $E_3$  cap here is in this direction already and  $E_1$  and  $E_2$ , this is the final orientation which we are showing by this part  $E_2'''$  by showing it by  $E_2'''$ , this equal to  $e_2$  cap. And  $E_1'''$ , we have shown this as  $e_1$  cap similarly this one as a  $e_3$  cap.

So, this configuration once we show it here ok, so you can forget about this figure. If you have confusion in understanding this for the time being just forget about this and just look at this figure. So, finally as we have discussed last time that once we rotate the body axis from the initial axis, so it will come here some what is the sort of orientation. And in this sort of orientation this is how your  $h$  vector is located. So,  $h$  vector is fixed, it is a non-rotating vector, but to your  $\omega$  vector it rotates ok, because  $\omega_1$  and  $\omega_2$  it is a rotating and  $h$  is fixed vector. So, it is so happens that once this plane is rotating, the plane containing this  $\omega$ ,  $\omega$  and  $e_3$  this is rotating ok.

So, we go back here on the previous one, they are co-planar. We have written them as a co-planar, but this one is fixed this is non-rotating, because it is a torque free case. So, it is a fixed in inertial place ok. So, here this is fixed in initial case place. And about this vector all these vectors are rotating and it, so happens that all of them remain co-planar ok.

So, finally what we can write that for this particular case, we indicate this angle as  $\theta$  ok. So,  $\tan \theta$  we can write this as see the component here. This  $\tan \theta$  this is nothing but  $\tan \theta$  we are measuring from this vertical axis to this is called a nutation angle,  $\theta$  is your nutation angle. So, how the  $\tan \theta$  will describe?  $\tan \theta$  will be the component of  $h$  here in this direction, so which is  $h \sin \theta$  along this direction, this is your  $h \sin \theta$  divided by this quantity here, which is  $I \omega_3$  ok. And  $h \sin \theta$  is nothing but going back on the previous page  $h \sin \theta$  is nothing but  $I \omega_3 \sin \theta$  ok. So,  $h \sin \theta$  is nothing but  $I \omega_3 \sin \theta$  and  $I \omega_3$  is nothing but  $I \omega_3$ .

So, this implies  $\sin \theta$  divided by  $\omega_3$ .  $\omega_3$  is a constant this is a constant and therefore this is turns out to be a constant and this implies  $\theta$  is a constant. So, again and again I am telling you refer to the previous lecture, where this picture was made in a much better way here the in the sort of space, it has got much complicated. And repeating the same figure and again we wasted a lot of time over that. So,  $\theta$  turns out to be a constant, this is one part.

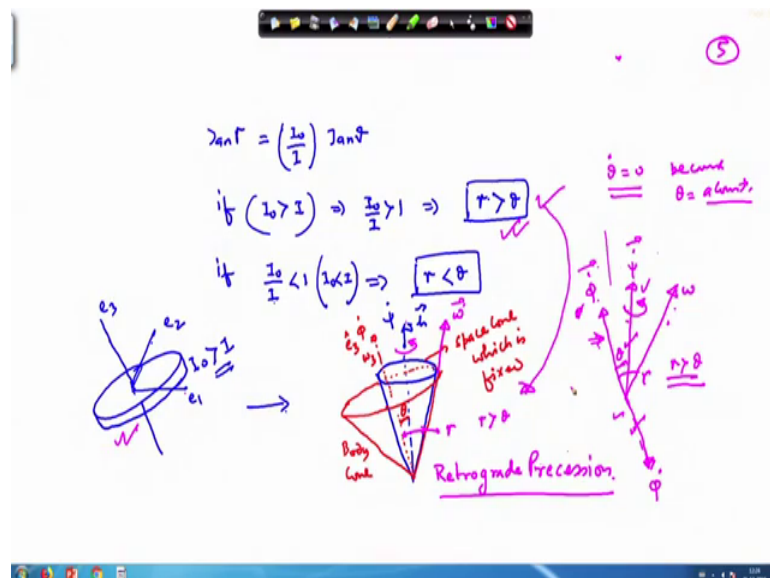
Now, there is another angle involved, which is of interest to us which is the angle  $\gamma$ . The  $\gamma$  angle is between this and this, this is your  $\gamma$  angle means between the  $\omega_3$  and  $\omega$  vector. So,  $\tan \gamma$  we can write as  $\tan \gamma$  you can see from this place, this is nothing but  $\omega_3 \sin \gamma$ ;  $\omega_3 \sin \gamma$  divided by  $\omega_3$ . So,

omega t divided by this is here we have missed out on those particular term, this is I times I 0, this is constant. So, this is a this also turns out to be a constant. But, here they are not independent, they are related to each other.

So, how we can get this just divide this tan theta by tan gamma. And once we divided, so this is I times omega t I 0 times omega 3 and here this becomes omega t omega 3 will cancel out and we get I by I 0. So, this implies tan theta equal to I by I 0 tan gamma or in the other way tan gamma equal to I 0 divided by I times tan theta. And this relationship is of great importance, this where getting from geometry ok. So, just by showing the geometry, we have come to this conclusion. And obviously, we have used the information's like omega 3, this is a constant. Then omega t, this is a constant as per our theoretical work ok.

Now, if we look here in this equation, so if I 0 is greater than I, which happens in the case of disk ok. So, for the case of disc I 0 is greater than I and therefore this quantities going to be greater than 1. So, this equality will hold, if and only if gamma is greater than theta. So, we write it on the next page tan gamma equal to tan gamma this equal to I 0 by I times tan theta.

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So, if I 0 is greater than I, this implying that I 0 by I is greater than 1, so this will imply that this quantity is greater than 1, so this implies gamma is greater than theta. On the

other hand, if  $I_0$  by  $I$  this is less than 1 ok, so this will imply that  $\gamma$  will be less than  $\theta$ .

So, we have two cases where  $I_0$  is. So, equally we can write in terms of  $I_0$  is less than  $I$ . So,  $I_0$  greater than  $I$  this implies  $\gamma$  is greater than  $\theta$ ,  $I_0$  less than  $I$  this implies  $\gamma$  is less than  $\theta$ . So, what does this mean, what does this imply. Here take this case take case of the disc. In this case  $I_0$  is greater than  $I$ , because the moment of inertia along this axis will be more than the other two axis this is  $e_1$ ,  $e_2$  and this is  $e_3$ .

So, moment of inertia along this direction is more. So, for this case it implies that  $\gamma$  will be greater than  $\theta$ , the situation is something like this as shown here;  $\gamma$  is greater than  $\theta$ . So, this is your  $\gamma$  angle from this place see here from this place to this place, this is your  $\gamma$  angle ok. So, this angle is greater than  $\theta$  ok. The other one if  $I_0$  is less than  $I$ , in that case  $\gamma$  turns out to be less than  $\theta$  means  $\gamma$  will this angle will turn out to be less than  $\theta$ .

So, here we have to be careful in describing the whole issue, the picture looks like something like this. For this disc, the motion it will appear as we will follow further, it will appear something like this. We have a cone here, fixed cone this is called, the space cone this is a fixed cone in the in inertial frame. And we show our  $h$  along this direction and  $\dot{\psi}$  also along this direction.

And this motion can be the motion of this particular disc this can be described as remember that this is your torque free case and the disc case initially inclined with the vertical by angle  $\theta$  ok, which we are calling as a nutation angle ok. So, the motion of the whole disc can be perceived not only in terms of see the perception in terms of  $\omega_1$ ,  $\omega_2$ ,  $\omega_3$ , it is a difficult how it is a rotating along the 3 axis it is a difficult to perceive. And therefore, we try to perceive this in terms of Euler angles which has physical meaning along which axis how it has rotated.

So, see here in this case first it is rotating like this, then it is rotating like and then it is rotating along this axis. Firstly, it has gone from this place to this place, then it is a rotating about this axis, then rotating about this axis. So, this is physically visualize able, but  $\omega_1$ ,  $\omega_2$ ,  $\omega_3$  these are not physically visualize able you would not be able to get anything from that.

And therefore, for this particular case the around this space cone we describe a body cone. This two ends are touching, this is called a body cone. And this blue one, it is called the space cone which is fixed ok. So, body cone you can consider that your body is towards the center of this cone ok. It is a the third axis, it is a lying along the center of this cone so, you have  $e_3$  along this direction ok. So,  $\dot{\phi}$  is along this direction so, you can consider it is a spinning ok.

So,  $\dot{\phi}$  is here then  $e_3$  is also here and  $\omega_3$  is also here. Then the angle between  $\omega_3$  direction and this one this is a  $\theta$  angle ok. So, here we can show this as  $\theta$  and  $\omega$  vector here in this case it so happens, it we can show it by some other color  $\omega$  vector it lies along this direction ok. And the angle from  $\omega_3$  to this place, this is your  $\gamma$ . So, here in this case  $\gamma$  is greater than  $\theta$  as we have shown in this place.

And obviously, you see that  $\dot{\theta}$  will be equal to 0 because  $\theta$  this is a constant. So,  $\theta$  is a constant, therefore,  $\dot{\theta}$  it turns out to be 0 means now we have the situation that. The motion which we have shown about this line with the motion which we have soon about this particular line here this way this motion is not present, when  $\theta$  is remaining constant want it has tilted and it is remaining like that. The only motion that appears there is constituted by  $\dot{\phi}$  and  $\dot{\psi}$  in terms of Euler angles which we can see here, it is a like this is your  $\dot{\phi}$  and this is  $\dot{\psi}$ .

But as you see that the  $\omega$  is lying here on this side ok, so in this kind of situation, it will so happen that the sense of rotation means if we are taking this to be positive means it is a going like this anticlockwise ok, we have taken this direction to be positive, here it is like this ok. So, this direction we have taken to be positive, and  $\omega$  it appears here in this direction.  $\omega$  is appearing somewhere here. You can see that if we look in the terms of Euler angles ok so, this is one vector and this is another vector. So, the combination of these two vectors should lie here in this place, but the  $\omega$  is lying here as per our conclusion from this place ok.

So, therefore, this implies that these vector cannot be in this direction rather this vector has to lie here in this direction  $\dot{\phi}$ .  $\dot{\phi}$  will be lying opposite to this. And as a consequence of this the summation of this vector and vector some of this and this vector this will be lying here in this direction. So, this kind of rotation where the  $\dot{\phi}$  is not

having the same spins as the  $\dot{\psi}$  ok, the spin is not having the same says as the precession is called a retrograde rotation. So, this is retro grade or retrograde precession, it is called the retrograde precession.

So, from this geometry, we have been able to find it out this particular particularly you remember this angle here this is  $\theta$ , this angle is  $\gamma$  and where  $\gamma$  is greater than  $\theta$ . So,  $\dot{\phi}$  is bound to lie along this direction. It cannot be here in this direction ok. We initially we have shown that it is a lying here in this direction, but for this kind of disc it is a bound to go downward and vector some of this and this will give you this  $\omega$ . And this  $\omega$ , obviously, you can break along three free body axis  $e_1$ ,  $e_2$  and  $e_3$ . But here in this case especially this  $\dot{\theta}$  is absent and because of this the whole thing gets simplified. So, we continue in the next lecture.

Thank you.