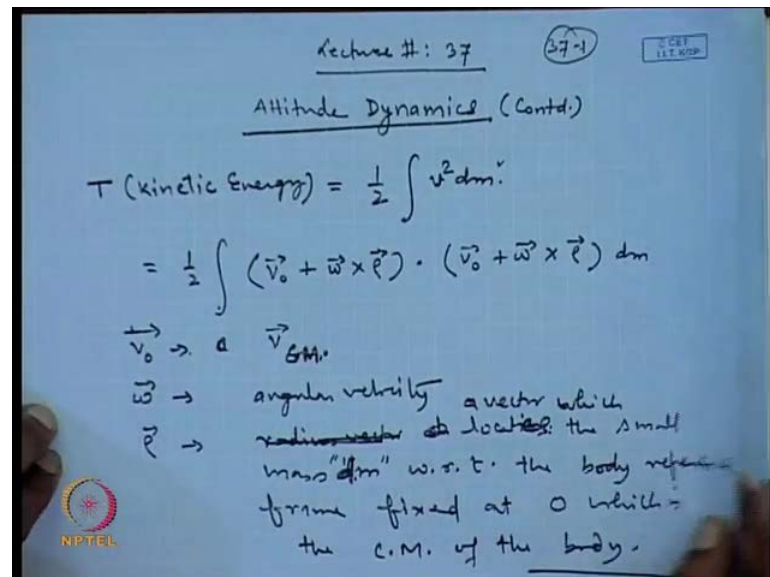


Space Flight Mechanics
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Model No. 01
Lecture No. 37
Attitude Dynamics (Contd.)

We have been continuing with attitude dynamics so, in that context last time we derived the angular momentum equation. So, we continue further so, today will find out the equation for the kinetic energy of a rotating body and rotating rigid body and also the Euler's equation, Euler's dynamical equation. So, which describes if the torque is applied to a rigid body so how, the angular momentum of the rigid body will change. So, the equation derived the Euler equation it relates that.

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Lecture #: 37 (37-1)

Attitude Dynamics (Contd.)

$$T \text{ (Kinetic Energy)} = \frac{1}{2} \int v^2 dm$$

$$= \frac{1}{2} \int (\vec{v}_0 + \vec{\omega} \times \vec{r}) \cdot (\vec{v}_0 + \vec{\omega} \times \vec{r}) dm$$

$\vec{v}_0 \rightarrow$ \vec{v}_{GM}
 $\vec{\omega} \rightarrow$ angular velocity a vector which
 $\vec{r} \rightarrow$ radius vector at location of the small mass "dm" w.r.t. the body reference frame fixed at O which is the C.M. of the body.

So, let us continue with the kinetic energy first. So, let us say this is the kinetic energy. So, this can be written as $\frac{1}{2} \int V^2 dm$ where, V is the velocity of the mass dm . We can write this as $\frac{1}{2} \int V_0^2 + 2 \vec{v}_0 \cdot (\vec{\omega} \times \vec{r}) + \omega^2 r^2 dm$ once we are considering the O to be located at the center of mass, this is V_{CM} , ω this is the angular velocity and r is the radius vector or say it is a vector, r is a

vector which locates the small mass dm with respect to the body reference frame fixed at O which is the center of mass of the body.

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$$\begin{aligned}
 T &= \frac{1}{2} \int \left[\vec{v}_O \cdot \vec{v}_O + 2 \vec{v}_O \cdot (\vec{\omega} \times \vec{\rho}) + (\vec{\omega} \times \vec{\rho}) \cdot (\vec{\omega} \times \vec{\rho}) \right] dm \\
 &= \frac{1}{2} \int v_O^2 + 2 \vec{v}_O \cdot (\vec{\omega} \times \vec{\rho}) + (\vec{\omega} \times \vec{\rho}) \cdot (\vec{\omega} \times \vec{\rho}) dm \\
 &= \frac{1}{2} \int v_O^2 dm + \vec{v}_O \cdot \int (\vec{\omega} \times \vec{\rho}) dm + \frac{1}{2} \int (\vec{\omega} \times \vec{\rho}) \cdot (\vec{\omega} \times \vec{\rho}) dm \\
 &= \frac{1}{2} M v_O^2 + \vec{v}_O \cdot \left[\vec{\omega} \times \int \vec{\rho} dm \right] + \frac{1}{2} \int (\vec{\omega} \times \vec{\rho}) \cdot (\vec{\omega} \times \vec{\rho}) dm = 0
 \end{aligned}$$

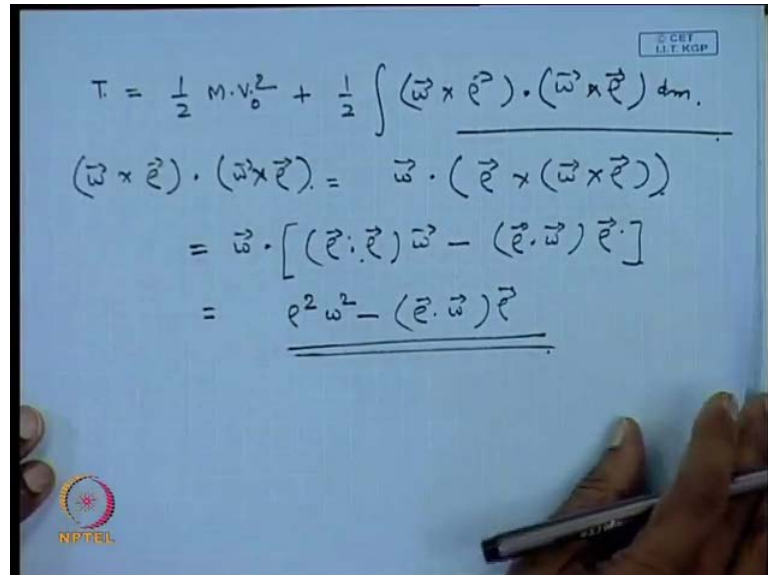
Now, we can expand the equation number. This is the equation that we have got, let us name this as equation number 1 so, we can expand it plus we can write here 2 times ω cross ρ ω cross ρ dot. So, from here we get V_O^2 and this is $2 V_O$ times $2 V_O$ dot ω cross ρ . So, V_O is the velocity of the center of the mass and integration is being done over the mass of the body. So, we can take it outside the integration sign. So, let us write here also in this place, we will write $\frac{1}{2} V_O^2 dm$ plus V_O dot.

Find this V_O^2 this can be taken outside and directly we can write $\frac{1}{2}$, the mass of the body times V_O^2 plus V_O this is the V_O dot. So, V_O dot we have already taken out now we can look into this moreover, we can put the whole thing in this fashion ω cross this is ρdm . So, this can be expanded and you can look from this place the quantity, ρdm this is nothing but because we are taking the point O to be the center of mass and ρ is being measured from the center of mass.

We have the body here and this is the point O which is the center of mass and ρ is the reduce vector of any point P that we choose earlier if you remember from our last lecture. So, this quantity will turn out to be equal to 0 because at the center of mass if you integrate or the moment of mass about the center of mass it will be 0. So, integration

taken over the whole body. So, this term drops out. So, this term will drop out leaving us only these two terms.

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Handwritten derivation on a blue background:

$$T = \frac{1}{2} M v_0^2 + \frac{1}{2} \int (\vec{\omega} \times \vec{r}) \cdot (\vec{\omega} \times \vec{r}) dm$$

$$(\vec{\omega} \times \vec{r}) \cdot (\vec{\omega} \times \vec{r}) = \vec{\omega} \cdot (\vec{r} \times (\vec{\omega} \times \vec{r}))$$

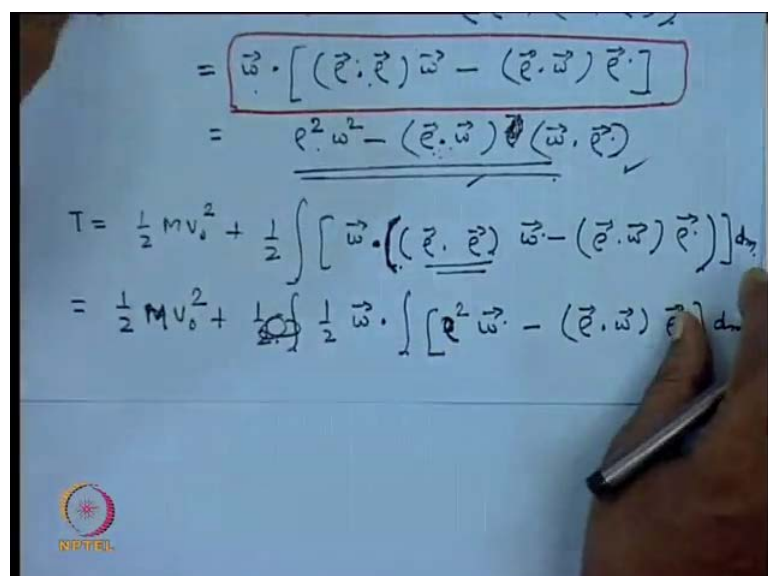
$$= \vec{\omega} \cdot [(\vec{r} \cdot \vec{r}) \vec{\omega} - (\vec{r} \cdot \vec{\omega}) \vec{r}]$$

$$= \underline{\underline{r^2 \omega^2 - (\vec{r} \cdot \vec{\omega})^2}}$$

NPTEL logo is visible in the bottom left corner.

Therefore, the kinetic energy we can write as $\frac{1}{2} M v_0^2$ plus now we need to find out what this quantity $\omega \times r$, what $\omega \times r$ will be? So, we can write this as $\omega \cdot r$ and then we can expand it further, this is the simple relationship here using for the dot and the cross product. So, this becomes $r \cdot r \omega$ minus $r \cdot \omega$ and now we can take again this quantity becomes $r^2 \omega^2$ this becomes ω^2 and minus $r \cdot \omega$.

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Continuation of the handwritten derivation:

$$= \vec{\omega} \cdot [(\vec{r} \cdot \vec{r}) \vec{\omega} - (\vec{r} \cdot \vec{\omega}) \vec{r}]$$

$$= \underline{\underline{r^2 \omega^2 - (\vec{r} \cdot \vec{\omega})^2}}$$

$$T = \frac{1}{2} M v_0^2 + \frac{1}{2} \int [\vec{\omega} \cdot ((\vec{r} \cdot \vec{r}) \vec{\omega} - (\vec{r} \cdot \vec{\omega}) \vec{r})] dm$$

$$= \frac{1}{2} M v_0^2 + \frac{1}{2} \int \vec{\omega} \cdot [r^2 \vec{\omega} - (\vec{r} \cdot \vec{\omega}) \vec{r}] dm$$

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So, this relationship we are going to utilize later on where we are measuring the term $\omega \cdot \rho \rho \omega$ and this is $\rho \cdot \omega$ times $\omega \cdot \rho$ or $\rho \cdot \omega$ it is the same. Now, the kinetic energy finally, we can write as $\frac{1}{2} M v_0^2$ plus $\frac{1}{2}$ inside the integration sign we have $\omega \cdot \rho \rho \omega$ we can choose from here times $\rho \cdot \rho$.

This $\omega \cdot \rho \rho \omega$ we can take it outside. So, let us take it outside $\omega \cdot \rho \rho \omega$ and then inside the integration sign this quantity will represent as $\rho \rho$ minus $\omega \cdot \omega$ and this becomes $\rho \cdot \rho$.

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$$T = \frac{1}{2} M v_0^2 + \frac{1}{2} \vec{\omega} \cdot \left(\int [\vec{r}^2 \vec{E} - \vec{r} \vec{r}] dm \right) \vec{\omega}$$

$$= \frac{1}{2} M v_0^2 + \frac{1}{2} \vec{\omega} \cdot \left(\int [\vec{r}^2 \vec{E} - \vec{r} \vec{r}] dm \right) \vec{\omega}$$

$$T = \frac{1}{2} M v_0^2 + \frac{1}{2} \vec{\omega} \cdot \vec{I} \cdot \vec{\omega}$$

Equation for kinetic energy of a rotating rigid body

So, further working on this equation so, we can write T equal to $\frac{1}{2} M v_0^2$ plus $\frac{1}{2} \omega \cdot \rho \rho \omega$ and inside the integration now if you remember that any vector can be written in terms of a unit dyadic. So, our unit dyadic we wrote as \vec{E} double bar. So, $\vec{E} \cdot \omega$ this is nothing but equal to ω this is identically equal to ω . So, therefore, this equation can be written in this fashion and here we can write this as $\rho \rho$, this ρ we can take from this place from the this side to this side and we can write in front of this bracket. And then we can write here dm and this whole thing we can put in a bracket and dot ω so, this ω we have taken it.

Now, this ω the ω , which was present here let us keep it here itself, right now we have $\rho \cdot \omega$. So, this dot product is between ρ and ω , in the next step $\omega \cdot \rho \rho \omega$ \vec{E} double bar minus now if from here we can write here $\rho \rho$

$\mathbf{m} \cdot \boldsymbol{\omega}$. We are taking the dot product this $\boldsymbol{\omega}$ outside this whole bracket from here and here both the places we are taken it outside. Now, if you remember this quantity is nothing but our inertia dyadic that we have defined earlier so, we can write this as \mathbf{I} double bar and so $\boldsymbol{\omega}$. So, this gives the equation for the kinetic energy and the same thing if you try to write in the form of the matrix notation.

So, this will change to $\boldsymbol{\omega}^T \mathbf{I} \boldsymbol{\omega}$ and \mathbf{I} is indicating the inertia matrix or either we can put \mathbf{g} to indicate that this is inertia matrix and so, $\boldsymbol{\omega}^T \mathbf{g} \boldsymbol{\omega}$. So, this gives you the equation for kinetic energy. Kinetic energy of a rotating rigid body and it is easy to identify. This is the kinetic energy of the body rigid body, considering this as a particle or a of mass m and this term represent the energy due to the rotational motion. So, total energy becomes the translational energy plus translational kinetic energy plus rotational kinetic energy.

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(37-5)

Euler's Dynamical Equation.

$$\vec{T} = \frac{d\vec{H}}{dt}$$

where \vec{H} is the angular momentum at time "t" in inertial reference frame.

$$\vec{H} = h_1 \vec{e}_1 + h_2 \vec{e}_2 + h_3 \vec{e}_3$$

$\vec{e}_1, \vec{e}_2, \vec{e}_3$ are unit vectors along the body axes.

$$\vec{T} = \frac{d\vec{H}}{dt} = \dot{h}_1 \vec{e}_1 + h_1 \frac{d\vec{e}_1}{dt} + \dot{h}_2 \vec{e}_2 + h_2 \frac{d\vec{e}_2}{dt} + \dot{h}_3 \vec{e}_3 + h_3 \frac{d\vec{e}_3}{dt}$$

So, we have completed this. Now, we go to the next step of deriving the Euler's dynamical equation. So, we know from our elementary physics that the if torque T is applied on a rigid body, then this can be written as torque is equal to dH by dt where, h is the angular momentum at time t in inertial reference frame. Now, this h this can be either capital H or a small h , last time we have used the small h notation. So, any one of them you can use. So, this can be written as $h_1 \mathbf{e}_1 + h_2 \mathbf{e}_2 + h_3 \mathbf{e}_3$.

What we have doing here that the angular momentum vector of the rigid body, this is in the inertial reference frame, but we can take the components of this angular momentum along the body axis. So, this is broken along the body $\hat{e}_1, \hat{e}_2, \hat{e}_3$, these are the unit vectors along the body. So, if we differentiate this so, we can write this as $\dot{h}_1 \hat{e}_1 + h_1 \dot{\hat{e}}_1 + \dot{h}_2 \hat{e}_2 + h_2 \dot{\hat{e}}_2 + \dot{h}_3 \hat{e}_3 + h_3 \dot{\hat{e}}_3$.

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The image shows a handwritten derivation on a blue background. At the top right, there is a circled number '37-5' and a small logo for 'CET IIT KGP'. The main derivation starts with the equation:

$$\Rightarrow \vec{T} = \frac{d}{dt} (\vec{H}) = \frac{d}{dt} (h_1 \hat{e}_1 + h_2 \hat{e}_2 + h_3 \hat{e}_3) + (h_1 \dot{\hat{e}}_1 + h_2 \dot{\hat{e}}_2 + h_3 \dot{\hat{e}}_3)$$

Below this, the first term is equated to the time derivative of the total angular momentum in the body frame:

$$h_1 \dot{\hat{e}}_1 + h_2 \dot{\hat{e}}_2 + h_3 \dot{\hat{e}}_3 = \left. \frac{d\vec{H}}{dt} \right|_B$$

Then, the derivatives of the unit vectors are given as:

$$\begin{aligned} \dot{\hat{e}}_1 &= \vec{\omega} \times \hat{e}_1 \\ \dot{\hat{e}}_2 &= \vec{\omega} \times \hat{e}_2 \\ \dot{\hat{e}}_3 &= \vec{\omega} \times \hat{e}_3 \end{aligned}$$

At the bottom left, there is a small circular logo with 'NPTEL' text.

So, this implies torque is equal to $\dot{h}_1 \hat{e}_1 + h_1 \dot{\hat{e}}_1 + \dot{h}_2 \hat{e}_2 + h_2 \dot{\hat{e}}_2 + \dot{h}_3 \hat{e}_3 + h_3 \dot{\hat{e}}_3$. Now, in this two this is one part and this is the second part. This we can identify as this can be written as $d\vec{H}/dt$ in the body or we are using capital notation for H . So, we can write in this way $d\vec{H}/dt$ with respect to the body. And this part we need to explore. So, now, we have to first evaluate $\dot{\hat{e}}_1$. So, $\dot{\hat{e}}_1$ will be nothing but the angular velocity of the rigid body times \hat{e}_1 and it is a very easy to prove we have earlier proved also this things $\dot{\hat{e}}_2$ similarly, this can be written as $\dot{\hat{e}}_2 = \vec{\omega} \times \hat{e}_2$ and $\dot{\hat{e}}_3 = \vec{\omega} \times \hat{e}_3$.

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The image shows a handwritten derivation on a whiteboard. At the top right, there is a circled number '37-7' and a small logo for 'SCET U.T. KGP'. The derivation starts with the equation:

$$\left. \frac{d\vec{H}}{dt} \right|_I = \vec{T} = \left. \frac{d\vec{H}}{dt} \right|_b + (h_1 \vec{\omega} \times \hat{e}_1 + h_2 \vec{\omega} \times \hat{e}_2 + h_3 \vec{\omega} \times \hat{e}_3)$$

This is followed by an equation where the terms in parentheses are grouped into a single vector \vec{h} :

$$= \left. \frac{d\vec{H}}{dt} \right|_b + \vec{\omega} \times (h_1 \hat{e}_1 + h_2 \hat{e}_2 + h_3 \hat{e}_3)$$

The final result, underlined, is:

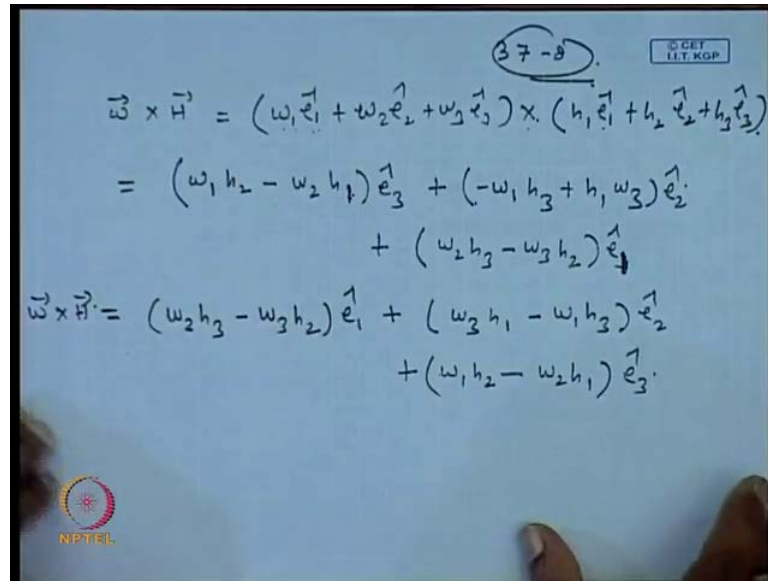
$$\left. \frac{d\vec{H}}{dt} \right|_I = \left. \frac{d\vec{H}}{dt} \right|_b + \vec{\omega} \times \vec{H}$$

In the bottom left corner, there is a logo for 'NPTEL'.

So, therefore, let us say this is equation number 3. So, equation number 3 gets reduced to $d\vec{H}/dt$ is in the inertial reference frame, we can put here I to indicate this is with respect to the inertial reference frame. This is equal to torque applied and this is with respect to the body (()) and the rest of the things you can write as h_1 times $\vec{\omega}$ cross \hat{e}_1 then h_2 times $\vec{\omega}$ cross \hat{e}_2 plus h_3 times $\vec{\omega}$ cross \hat{e}_3 , this is \vec{h} . So, now, you can recognize that this thing can be written as $h_1 \hat{e}_1$ plus $h_2 \hat{e}_2$ plus $h_3 \hat{e}_3$ and this quantity is nothing but \vec{h} .

So, $d\vec{H}/dt$ in the inertial reference frame, this can be written as $d\vec{H}/dt$ in the body reference frame. So, rate of change of the angular momentum in the body reference frame plus $\vec{\omega}$ cross the angular momentum vector.

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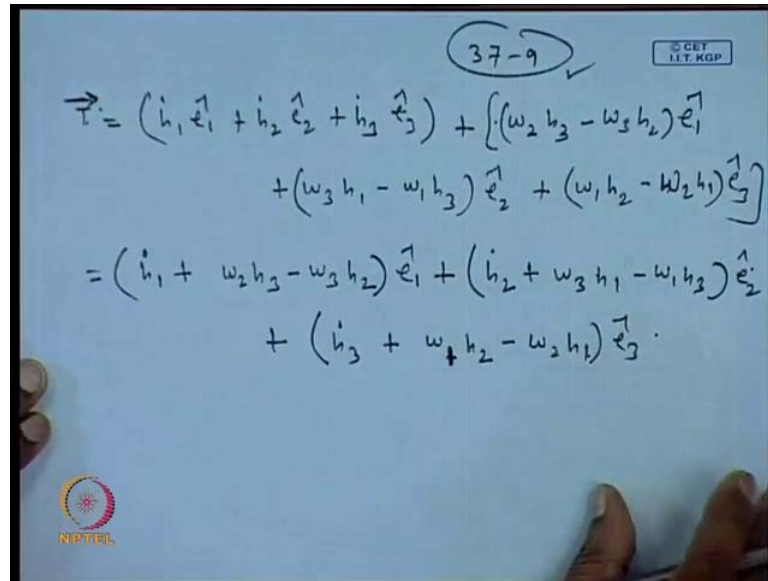


$$\begin{aligned}\vec{\omega} \times \vec{H} &= (\omega_1 \hat{e}_1 + \omega_2 \hat{e}_2 + \omega_3 \hat{e}_3) \times (h_1 \hat{e}_1 + h_2 \hat{e}_2 + h_3 \hat{e}_3) \\ &= (\omega_1 h_2 - \omega_2 h_1) \hat{e}_3 + (-\omega_1 h_3 + h_1 \omega_3) \hat{e}_2 \\ &\quad + (\omega_2 h_3 - \omega_3 h_2) \hat{e}_1 \\ \vec{\omega} \times \vec{H} &= (\omega_2 h_3 - \omega_3 h_2) \hat{e}_1 + (\omega_3 h_1 - \omega_1 h_3) \hat{e}_2 \\ &\quad + (\omega_1 h_2 - \omega_2 h_1) \hat{e}_3.\end{aligned}$$

So, we know that $\vec{\omega} \times \vec{H}$ can be written as $\omega_1 \hat{e}_1 + \omega_2 \hat{e}_2 + \omega_3 \hat{e}_3$. And we have to take the cross product of this, with \vec{h} and \vec{h} we can write as $h_1 \hat{e}_1 + h_2 \hat{e}_2 + h_3 \hat{e}_3$. Expanding this we can re write them as $\omega_1 h_2$. So, what we have to do once we take the cross product the self $\hat{e}_1 \times \hat{e}_1$ this will vanish only thing it will adjust with the other 2. So, this can be written as $\omega_2 h_1$.

Times \hat{e}_1 cap. Re writing the whole thing we bring in the third term make it the first term, $\omega_2 h_3$ minus $\omega_3 h_2$ \hat{e}_1 cap from here $\omega_3 h_1$ minus $\omega_1 h_3$ \hat{e}_2 cap plus now taking this term \hat{e}_3 . So, this is $\vec{\omega} \times \vec{H}$.

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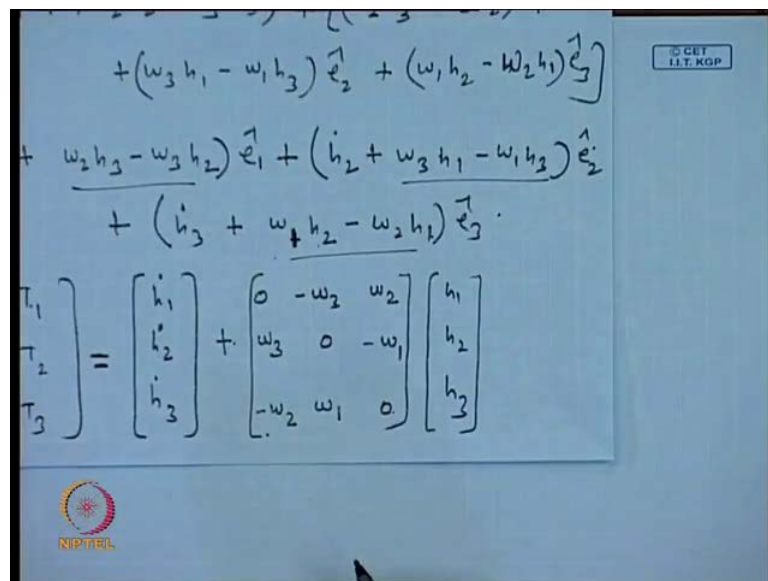
Handwritten derivation of \vec{T} in terms of \dot{h} and $\omega \times h$. The derivation is as follows:

$$\begin{aligned}\vec{T} &= (\dot{h}_1 \hat{e}_1 + \dot{h}_2 \hat{e}_2 + \dot{h}_3 \hat{e}_3) + [(\omega_2 h_3 - \omega_3 h_2) \hat{e}_1 \\ &\quad + (\omega_3 h_1 - \omega_1 h_3) \hat{e}_2 + (\omega_1 h_2 - \omega_2 h_1) \hat{e}_3] \\ &= (\dot{h}_1 + \omega_2 h_3 - \omega_3 h_2) \hat{e}_1 + (\dot{h}_2 + \omega_3 h_1 - \omega_1 h_3) \hat{e}_2 \\ &\quad + (\dot{h}_3 + \omega_1 h_2 - \omega_2 h_1) \hat{e}_3.\end{aligned}$$

Therefore, this T what we have written as $\dot{h}_1 \hat{e}_1 + \dot{h}_2 \hat{e}_2 + \dot{h}_3 \hat{e}_3$ and the $\omega \times h$ now we put here in this place $\omega \times h$.

We can combine the terms with \hat{e}_1 , \hat{e}_2 and \hat{e}_3 separately. So, this will become $\dot{h}_1 + \omega_2 h_3 - \omega_3 h_2$ times \hat{e}_1 plus $\dot{h}_2 + \omega_3 h_1 - \omega_1 h_3$ times \hat{e}_2 plus $\dot{h}_3 + \omega_1 h_2 - \omega_2 h_1$ times \hat{e}_3 .

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Handwritten matrix representation of the vector equation:

$$\begin{bmatrix} T_1 \\ T_2 \\ T_3 \end{bmatrix} = \begin{bmatrix} \dot{h}_1 \\ \dot{h}_2 \\ \dot{h}_3 \end{bmatrix} + \begin{bmatrix} 0 & -\omega_2 & \omega_3 \\ \omega_3 & 0 & -\omega_1 \\ -\omega_2 & \omega_1 & 0 \end{bmatrix} \begin{bmatrix} h_1 \\ h_2 \\ h_3 \end{bmatrix}$$

If now, if you look into this equation so, this T if you try to write in the matrix notation. So, you can write this as T_1, T_2, T_3 and here you will have $\dot{h}_1, \dot{h}_2, \dot{h}_3$ and

plus the other terms which are present. So, similarly, you can represent them. So, for representing them in a very simple format, there are compact notations for representing them. So, let us say that there are two ways of doing this $\omega_2 \omega_3$ like this terms are present here. So, first we take the ω outside this and write here h_1 , like $h_1 h_2 h_3$ and ω we write here in this place and later on the h_1 can be expanded in terms of the moment of inertia matrix and the angular velocity vector components along the body $(())$.

So, h_1 here if you take the first term here. So, you can see from here this the ω_2 is related to h_3 and positive sign while h_2 is related with ω_3 and this is having a negative sign. And ω_1 is not present in this so, we can write here 0. Similarly, we can take the next term here so; next term is h_1 times ω_3 and minus ω_1 times h_3 . So, ω_1 will come with a minus sign and here we will have 0, next we take this term, but this is related to h_2 . So, with h_2 we have the ω_1 term present and while with h_1 we have the minus ω_2 term present. So, we represent it in this format.

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The image shows a handwritten derivation on a blue background. At the top right, there is a small box with the text "© CBT I.I.T. KGP". Below it, the number "37-10" is circled. The main equation is:

$$\begin{bmatrix} h_1 \\ h_2 \\ h_3 \end{bmatrix} = \begin{bmatrix} I_{11} & -I_{12} & -I_{13} \\ -I_{12} & I_{22} & -I_{23} \\ -I_{13} & -I_{23} & I_{33} \end{bmatrix} \begin{bmatrix} \dot{\omega}_1 \\ \dot{\omega}_2 \\ \dot{\omega}_3 \end{bmatrix} + \begin{bmatrix} 0 & -\omega_3 & \omega_2 \\ \omega_3 & 0 & -\omega_1 \\ -\omega_2 & \omega_1 & 0 \end{bmatrix} \begin{bmatrix} I_{11} & -I_{12} & -I_{13} \\ -I_{12} & I_{22} & -I_{23} \\ -I_{13} & -I_{23} & I_{33} \end{bmatrix} \begin{bmatrix} \omega_1 \\ \omega_2 \\ \omega_3 \end{bmatrix}$$

Below this, the equation is simplified to:

$$\begin{bmatrix} h_1 \\ h_2 \\ h_3 \end{bmatrix} = \begin{bmatrix} I_{11} & -I_{12} & -I_{13} \\ -I_{12} & I_{22} & -I_{23} \\ -I_{13} & -I_{23} & I_{33} \end{bmatrix} \begin{bmatrix} \dot{\omega}_1 \\ \dot{\omega}_2 \\ \dot{\omega}_3 \end{bmatrix}$$

The NPTTEL logo is visible in the bottom left corner of the slide.

So, this equation so, this you can write further as $T_1 T_2 T_3$.

Now, look into this equation, once we wrote the $h_1 h_2 h_3$ dot so, what we assume that our body axis fixed in the body, rigid body itself. And therefore, the moment of inertia will not change. If you do not fix the body $(())$ in the rigid body so, moment of inertia will keep changing and that is what happens with respect to the inertial $(())$. In the

inertial $(())$ the moment of inertia of the body 1 because the moment of inertia we are defining with respect to the 3 $(())$ and if the body orientation is changing so, with respect to this $(())$ moment of inertia will change.

So, whichever point you choose the moment of inertia of the body changes. So, we need to fix the body $(())$ in the rigid body. So, once we fix it then \dot{h}_1 can be simply written as \dot{h}_1 we have \dot{h} we have written as let us write below first $\dot{h}_1 \dot{h}_2 \dot{h}_3$ this we have written as the I matrix $I_{11} \text{ minus } I_{12} \text{ so, here minus } I_{13} \text{ minus } I_{23} \text{ } I_{33}$ and here we have I_{12} with minus sign $I_{22} \text{ minus } I_{23} \omega_1 \omega_2 \omega_3$. So, \dot{h}_1 is the product of this vector and this one. So, this is basically we are taking you multiply each of the term here. So, if you differentiate you take write \dot{h}_1 if you are writing \dot{h}_1 so, these terms are not changing because the body axis is fixed in the body itself and therefore, only thing will happen that this $\omega_1 \dot{\omega}_2 \dot{\omega}_3$ will keep changing with time. So, here we can put the dot for that.

So, therefore, $T_1 T_2 T_3$ this will become equal to now this $\dot{h}_1 \dot{h}_2 \dot{h}_3$ in this equation need to be replaced by this equation. So, we insert here in this phase this is $I_{11} \omega_1 \dot{\omega}_2 \dot{\omega}_3$ and plus the matrix just now we have written here. So, $0 \text{ minus } \omega_3 \omega_2$ and then this is $\omega_3 0 \text{ minus } \omega_1$ you can see this is anti-symmetric matrix or asymmetric, this is basically asymmetric matrix. And then $\dot{h}_1 \dot{h}_2 \dot{h}_3$ we can similarly, write here \dot{h}_1 will be $I_{11} \text{ minus } I_{12} \text{ minus } I_{13}$ minus $I_{12} I_{22} \text{ minus } I_{23}$ with minus sign and I_{33} and this 2 be multiplied by $\omega_1 \omega_2$ and ω_3 .

So, whatever we have got it represents that if we apply a torque to a rigid body. So, the torque is directly taken in the inertial reference frame. Now, $T_1 T_2$ it's a taken the inertial reference frame, but you can consider that we have the components $T_1 T_2 T_3$ along the body axis. So, these are we are writing we have broken along the body $(())$ because all these $\omega_2 \omega_3$ they are all about the along the body $(())$. So, if you apply the torque whose components in the body $(())$ are $T_1 T_2$ and T_3 , then you get this dynamical equation which relates the rotational motion of the body. So, how the angular velocity will change it will be obtained from this equation.

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Now if the off-diagonal terms of inertia matrix are zero.

$$\begin{bmatrix} T_1 \\ T_2 \\ T_3 \end{bmatrix} = \begin{bmatrix} I_1 & 0 & 0 \\ 0 & I_2 & 0 \\ 0 & 0 & I_3 \end{bmatrix} \begin{bmatrix} \dot{\omega}_1 \\ \dot{\omega}_2 \\ \dot{\omega}_3 \end{bmatrix} + \begin{bmatrix} 0 & -\omega_3 \omega_2 & \omega_2 \\ \omega_3 & 0 & -\omega_1 \\ -\omega_2 & \omega_1 & 0 \end{bmatrix} \begin{bmatrix} I_1 & 0 & 0 \\ 0 & I_2 & 0 \\ 0 & 0 & I_3 \end{bmatrix} \begin{bmatrix} \omega_1 \\ \omega_2 \\ \omega_3 \end{bmatrix}$$

Euler's Dynamical Eq.

NPTEL

Now, if the off diagonal terms of inertia matrix are 0 then what we get T_1, T_2, T_3 now we write as $I_1 \dot{\omega}_1, I_2 \dot{\omega}_2, I_3 \dot{\omega}_3$ to indicate that these are serving as the principal moment of inertia along the 3 axes. And similarly, here we have $0 \pm \omega_3 \omega_2$ and this is nothing but your Euler's dynamical equation. This is what we intended to derive you can break it down and write each of the component separately.

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$$\left. \begin{aligned} T_1 &= I_1 \dot{\omega}_1 + \omega_2 \omega_3 (I_2 - I_3) \\ T_2 &= I_2 \dot{\omega}_2 + \omega_1 \omega_3 (I_1 - I_3) \\ T_3 &= I_3 \dot{\omega}_3 + \omega_1 \omega_2 (I_2 - I_1) \end{aligned} \right\} \begin{matrix} \hat{e}_1, \hat{e}_2, \hat{e}_3 \\ \downarrow \\ \text{principal} \\ \text{axis direction} \end{matrix}$$

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So, the first equation if you break it up first equation can be written as $I_1 \dot{\omega}_1$ dot multiplying this row by the whole column. So, we get $I_1 \dot{\omega}_1$ and then

we have to multiply this similarly, and write here. So, what you get? $I_2 \dot{\omega}_2$ plus $I_1 \dot{\omega}_1$ minus $I_3 \dot{\omega}_3$.

So, here your directions \hat{e}_1, \hat{e}_2 and \hat{e}_3 they are basically serving as the principal directions because the off diagonal terms are 0 here.

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$$\begin{aligned} T_1 &= I_1 \dot{\omega}_1 + \omega_2 \omega_3 (I_2 - I_3) \\ T_2 &= I_2 \dot{\omega}_2 + \omega_1 \omega_3 (I_1 - I_3) \\ T_3 &= I_3 \dot{\omega}_3 + \omega_1 \omega_2 (I_2 - I_1) \end{aligned} \quad \left\{ \begin{array}{l} \hat{e}_1, \hat{e}_2, \hat{e}_3 \\ \downarrow \\ \text{principal} \\ \text{axes directions} \end{array} \right.$$

Body is rotating about one of the principal axes.
 Say $\omega_1 \neq 0 = \omega_0$ $\omega_2 = \omega_3 = 0$

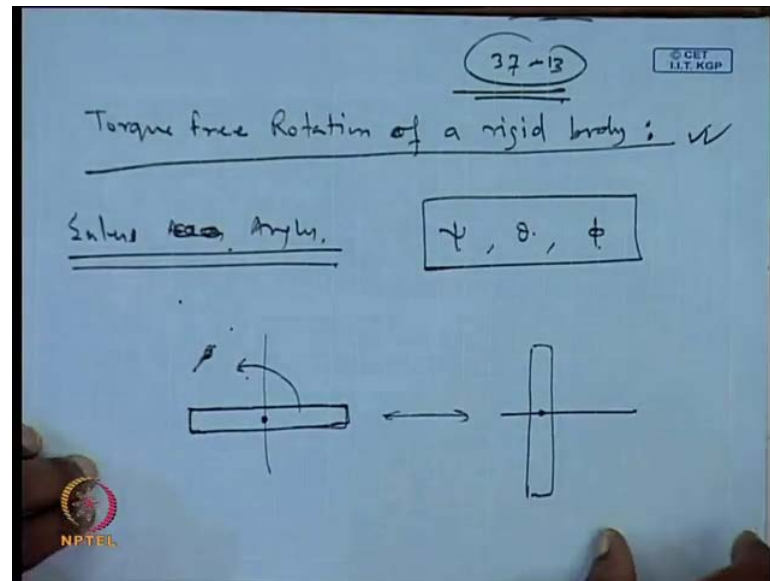
$$\begin{aligned} T_1 &= I_1 \dot{\omega}_1 = I \dot{\omega}_0 \\ T_2 &= 0 \\ T_3 &= 0 \end{aligned}$$

 $T_1 = I \dot{\omega}_0$

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Now, we in this case, if you considered that body is rotating about one of the principal directions, then let us say this is ω_1 which is not non 0 and this we are putting as ω_0 and ω_2 and ω_3 equal to 0. So, the above equation then will get reduce to T_1 equal to $I_1 \dot{\omega}_1$ and $\omega_2 \omega_3$ they are 0. So, therefore, those quantities vanish. Similarly, T_2 equal to now ω_2 itself is 0 is not present. So, this term also vanishes ω_3 is 0 so, this will vanish. And all similarly, other terms will become 0. This is the very simple equation that we have got T_1 equal to I times ω_0 dot. And earlier I have told you that the angular velocity vector and the angular momentum vector, they will coincide if and only if rotation is about 1 of the principal directions.

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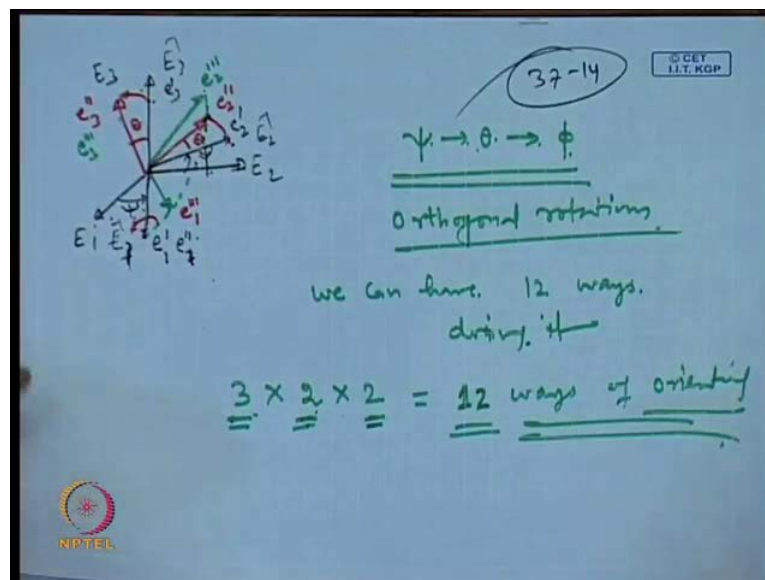


So, we go to the next step. Now, we consider the torque free rotation of a rigid body. So, why this is important for our study that say if we have in this space some satellite is left alone and no torque is applied. The torque's whatever the torque's magnetic torque's or the rockets which are on board are the reaction will nothing is working neither the suppose the solar pressure is also not there. So, torque due to solar pressure is also not present. So, if it is the satellite is free from the all sort of torques, then how the satellite is going to behave? So, if we orient the satellite say this is my satellite and this is moving in the orbit. So, I can have the orbit like this in which the satellite has to move in this way, always pointing in one direction, always pointing towards the earth.

So, is it possible that once we are putting in this configuration that it is going to remain in this configuration. So, in the case of because it is a long body and we have seen that because of the gravity gradient it tends to stabilize in this configuration only. But we will do a general derivation right now and we will see that if the bodies now, if the satellite is left in the orbit and depending on it is a moment of inertia along the different axis how much it is. So, whether the satellite will remain in the same configuration or means the same orientation or the orientation will change. So, this is related to stability of the rotational motion. If it is rotating in this direction like this, if it is rotating like this so, whether it will maintain that direction that becomes important to know or satellite orientation will or either it will change over a period of time.

So, this we can derive if we work on the torque free rotation of a for a rigid body. So, first of all we need to know about the few the details like the Euler angles. So, Euler angles are represented by psi theta and phi. So, in general if you see anybody we are having so, it can be oriented in the space, say if I have any object which is lying like this. So, you can bring it to this configuration by rotating it about the axis which is passing through this page. If you rotate anti clock wise by 90 degree so, it will come in this configuration. So, the next rotation you give so, the configuration will change. So, for any orientation of the body this three rotations are three consecutive rotations can be given and it will get oriented in certain direction.

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So, let us say that E_1 , E_2 and E_3 , these are the 3 axes with respect to which I have to orient the body. So, let us say that these are the vectors $E_3 \cap E_2 \cap E_1$ along this direction. And if I give a first rotation by angle psi about this z axis so, this is my new position and new position let us write this as e_1' and from here we write this as e_2' and e_3' remains along the same direction and this angle is basically your psi this angle is psi. The next rotation so, we have the new configuration in which this axis came to this position this axis moved to this position. So, we can give a rotation here about this by angle theta.

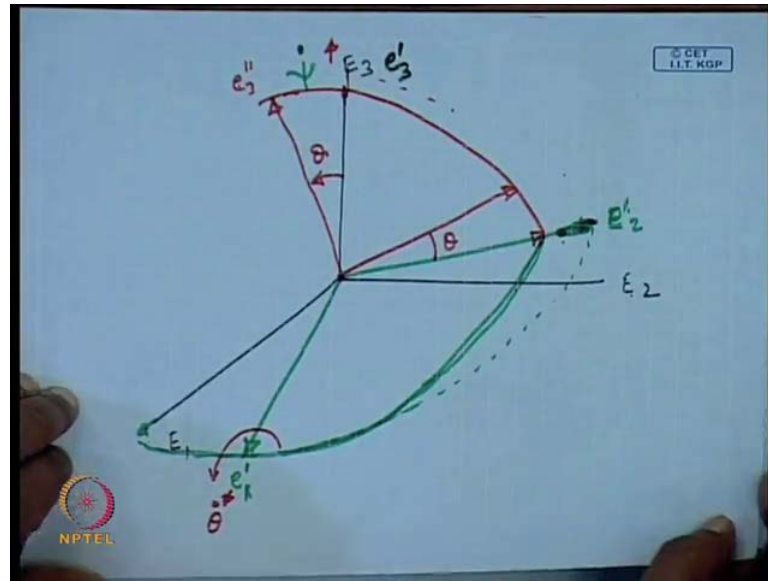
So, if we rotate it, this will move from this place to this place and this vector will move from this place to this place. So, our new vector will be e_3'' and here we

can write as e_2' and this is our new vector. And by how much angle we are rotating here? This by angle θ and you see that they are not in the same plane. First we rotated about the z axis then we are rotating about this is e_3 axis is serving as the z axis here and even is serving as the x axis. So, next we are rotating about this axis. The third rotation we can give in the plane about e_3' axis.

So, e_1 will move in this plane from here to here, the e_1' and e_2' now is the same. e_1' and e_2' it would be e_1'' they are the same. So, e_1'' will move from this place and it will come to e_1''' here in this place and this will move in this plane from here and it will go into this place and here, we write as e_2''' . While our e_3''' will remain in same direction because we have rotated about this axis so, we are basically giving three consecutive rotation which are indicated by ψ and then we are getting rotating by θ and then we are rotating it by ϕ . And these are called orthogonal rotations. So, for orienting a body in the inertial reference frame, we can have 12 ways of doing it. So, the first rotation I can choose among any of the 3 axis.

So, the first rotation can be given in three ways, once we have given the first rotation. So, the next rotation can be given only about the remaining 2 axis. We cannot give about the first at $(())$ like here we choose the third axis to rotate it so, there after we are left with e_1 and e_2 $(())$ to rotate it. We cannot give the next rotation about the e_3 $(())$ again if you do that so, it is the same rotation nothing it is a rotating in the same plane. So, only thing you are adding to this ψ is $\Delta\psi$ more. Some more angle you are adding so, the first rotation can be given in 3 ways, the next rotation can be chosen in two ways once you are chosen the next rotation about any remaining of the 2 $(())$. So, suppose that we choose the e_1 x. So, next rotation can be chosen among the rest of the two x so, the total of 3 into 2 into 2 so, total twelve ways of orienting. So, giving the different magnitude of the ψ , θ and ϕ and reshuffling their order, you can rise till the same orientation in 12 different ways.

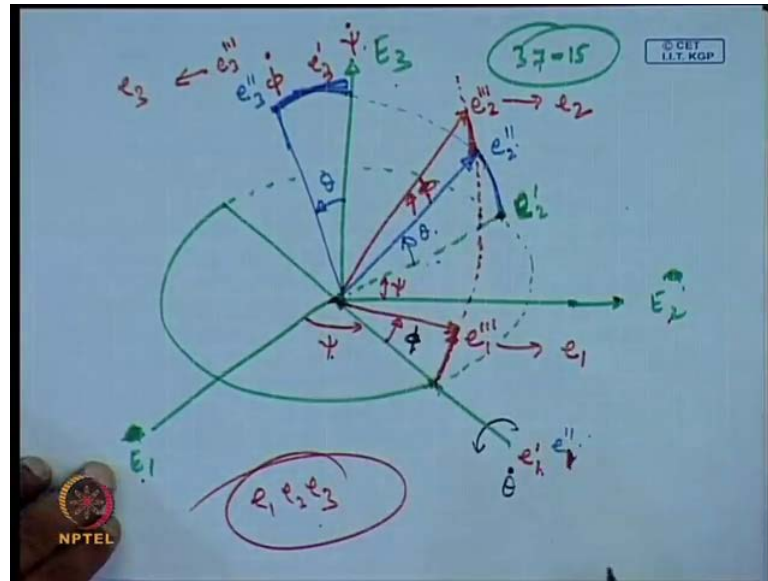
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So, what we have done just now, we can explode it little more. This is our E_1, E_2 and E_3 and first we are rotating about E_3 axis, and rotating about by angle ψ . So, the rate of change of this angle $\dot{\psi}$ we indicate here so, the next vector is from here to here. So, this we write as e_1' . And similarly, in the same plane e_2 will move to this place and this we have indicated as e_2' . Next, the rotation we gave about this x by θ . So, $\dot{\theta}$ can be shown along this direction this is the direction for $\dot{\psi}$. So, if we are rotating about this. So, this x will rotate, it will move from this place to this place. So, we take it up and this will go from here to here.

So, this vector will move from here to here, this vector is moving from this place to this place. So, here you have e_3, e_3' and in this place you have e_3'' . And this vector will also move from this place to this place so, this is your let us put this vector here. So, first of rotation we show by this line so, this has rotated from here to here by θ . You see these are the different planes, they are not in the same plane rather I will make another figure to make with more clear, this figure is not very good.

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And this is E_1 and this is the direction E_2 , this is the direction E_3 . So, the first rotation we give by angle ψ . So, ψ dot is along this direction and your new orientation is e_1 prime and E_2 prime similarly, it will rotate in the $E_1 E_2$ plane, E_2 will move in the $E_1 e_2$ plane and let us say this gets oriented from this place to this place. Obviously, goes into this place, this is E_2 prime and here we have the e_1 prime. The next rotation we are giving is about this by angle θ so, θ dot will point in this direction. So, if you rotate it so, this vector and this vector now it is forming a triode e_1 prime e_2 prime and here we have the e_3 prime. So, it is forming a triode. So, this e_2 prime and e_3 prime will rotate about the e_1 prime. So, we can make it rotate. This angle is turned to this ψ and let us say, that we rotate it so, it rotates like this.

And a move from this has been shown by dotted line to show that this is in the $e_1 e_2$ plane. It moves from here to here, this is here e_2 double prime, e_2 double prime is along this direction itself. And this will move from this place to this place. This is the angle θ and here we will write as e_3 double prime. Next rotation we can give which we will be contained in e_2 double prime, this is e_1 double prime here and e_2 double prime plane. And we will rotate about this xso, e_3 triple prime will remain here and we are rotating about and this angle will rotate by ϕ . And we can show it that this will move from this place to this place. So, this is angle ϕ .

And this will move from this place to this place, this is angle ϕ . So, they are not in the same plane you can see. So, these are the three orthogonal rotations. So, any set of vector is once need to be converted from the inertial reference frame to the body reference frame so, say the final orientation which we have got as e triple prime we write as e_3 and here this is moving from this place to this place. So, we write this as e_1 triple prime which we will show as e_1 and e_2 moves from this place to this place. So, we write as e_2 triple prime which will write as e_2 . So, e_1 e_2 e_3 they constitutes the body reference frame and capital E_1 capital E_2 and capital E_3 it constitutes the inertial reference frame. So, changing from the inertial reference frame to the body reference frame it requires this three rotations and they can be represented in terms of the matrix. So, we will continue with this next time. Thank you very much.