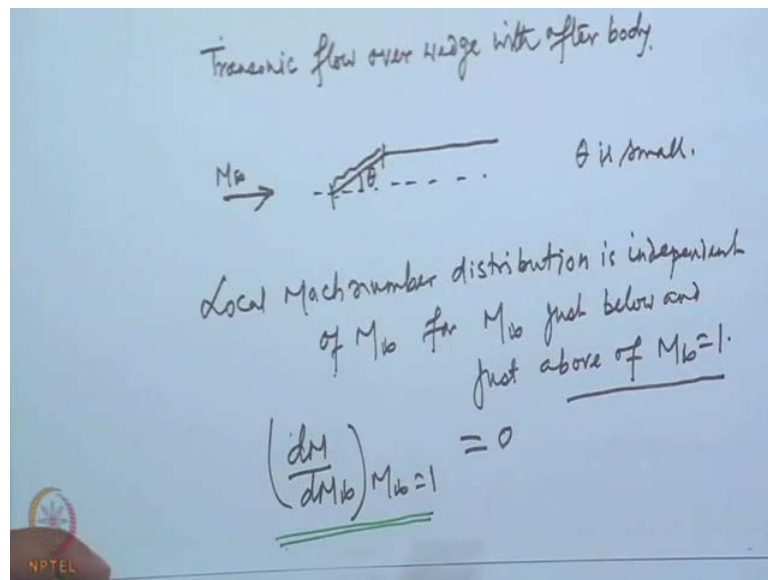


High Speed Aerodynamics
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Lecture No. # 39
Transonic Flow (Contd.)

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So, we will continue our discussion, on transonic flow over a wedge with upper body, where we are considering a wedge, which is extended and our approximation; so that, we are within the framework of small perturbation theory, that M in θ is small; as θ increases, the accuracy of the results decreases. However in this case, mostly our discussion, the qualitative nature of the flow, so they of course, remain more or less the same. We have seen that, the local Mach number distribution on the wedge; that is, on this part of the wedge, local **Mach number distribution** Mach number distribution is independent of M_∞ .

For M_∞ , **just below and above** just below and just above of M_∞ **just above**

of $M \rightarrow \infty$ equal to 1. That is, when the system Mach number is very close to 1? The Mach number distribution on the surface remains the same, which we expressed mathematically as $\frac{dC_p}{dM} \bigg|_{M \rightarrow \infty} = 0$.

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The image shows a handwritten derivation on a blue background. At the top right, there is a small box with the text "© CET I.I.T. KGP". The derivation starts with the equation for the pressure coefficient C_p in terms of the Mach number M_∞ and the specific heat ratio γ :

$$C_p = \frac{2}{\gamma M_\infty^2} \left[\frac{2 + (\gamma - 1) M_\infty^2}{2 + (\gamma - 1) M_\infty^2} \right]^{\frac{\gamma}{\gamma - 1}} - 1$$

Next, the equation is rearranged to take the logarithm of both sides:

$$\Rightarrow \log\left(\frac{\gamma}{2} M_\infty^2 C_p + 1\right) = \frac{\gamma}{\gamma - 1} \left\{ \log(2 + (\gamma - 1) M_\infty^2) - \log[2 + (\gamma - 1) M_\infty^2] \right\}$$

Then, the instruction "Differentiate with M_∞ and let $M_\infty \rightarrow 1$ " is written. This leads to the final result for the derivative of C_p with respect to M_∞ at $M_\infty = 1$:

$$\Rightarrow \frac{\gamma C_p^* + \frac{\gamma}{2} \left(\frac{dC_p}{dM_\infty} \right)_{M_\infty=1}}{1 + \frac{\gamma}{2} C_p^*} = \frac{2\gamma}{\gamma + 1}$$

In the bottom left corner, there is a logo for NPTEL.

Now let us see, how this affix the pressure coefficient near Mach number unity? The pressure coefficient is given by $2 + \gamma - 1 M_\infty^2$ by $2 + \gamma - 1 M_\infty^2$ to the power $\frac{\gamma}{\gamma - 1} - 1$. So, this is the pressure coefficient at a point, where the local Mach number is M . Taking logarithm, this gives $\log\left(\frac{\gamma}{2} M_\infty^2 C_p + 1\right) = \frac{\gamma}{\gamma - 1} \left\{ \log(2 + (\gamma - 1) M_\infty^2) - \log[2 + (\gamma - 1) M_\infty^2] \right\}$. Differentiate this with respect to M_∞ , **differentiate with M_∞** and let M_∞ approaches 1; this gives then, $\gamma C_p^* + \frac{\gamma}{2} \frac{dC_p}{dM_\infty} \bigg|_{M_\infty=1} = \frac{2\gamma}{\gamma + 1} \left(1 + \frac{\gamma}{2} C_p^* \right)$, that is $\frac{2\gamma}{\gamma + 1}$.

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$$C_p = \frac{1}{2} M_\infty^2 \left[\left(\frac{2 + (\gamma - 1) M_\infty^2}{2 + (\gamma - 1) M_\infty^2} \right)^{-1} - 1 \right]$$

$$\Rightarrow \log\left(\frac{1}{2} M_\infty^2 C_p + 1\right) = \frac{\gamma}{\gamma - 1} \left\{ \log(2 + (\gamma - 1) M_\infty^2) - \log[2 + (\gamma - 1) M_\infty^2] \right\}$$

Differentiate with M_∞ and let $M_\infty \rightarrow 1$

$$\Rightarrow \frac{\gamma C_p^* + \frac{\gamma}{2} \left(\frac{dC_p}{dM_\infty} \right)_{M_\infty=1}}{1 + \frac{\gamma}{2} C_p^*} = \frac{2\gamma}{\gamma + 1}$$

$$C_p^* = (C_p)_{M_\infty=1}$$

In this, the C_p star is C_p at M equal to 1. In which, this C_p star is C_p at M infinity is equal to 1 with this, then dC_p/dM infinity; at M infinity equal to $1/4$ by $\gamma + 1$ to 1 minus half.

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$$\left(\frac{dC_p}{dM_\infty} \right)_{M_\infty=1} = \frac{4}{\gamma + 1} \left(1 - \frac{1}{2} C_p^* \right)$$

Contribution of C_p to C_D , is $C_p \sin^2 \theta \approx C_p \theta$.

$$\Rightarrow \left(\frac{dC_D}{dM_\infty} \right)_{M_\infty=1} = \frac{4\theta}{\gamma + 1} - \frac{2C_D^*}{\gamma + 1}, \quad C_D^* = (C_D)_{M_\infty=1}$$

Within the framework of small perturbation theory

$$\frac{1}{M_\infty} \ll 1, \Rightarrow \frac{1}{2} C_p^* \ll 1.$$

Now at any point, where the pressure coefficient is C_p , that contributes to the drag as C

$p \sin \theta$. So, contribution of C_p to C_D is $C_p \sin \theta$; where θ is the semi vertex angle of the wedge, which for small θ is C_p into θ . So, we can now get that, $dC_D/dM_\infty = 4\theta/(\gamma+1)$ and again C_{D^*} is C_D at $M_\infty = 1$. So, this very simple analysis gives us, some indication about the drag curve flow, near M_∞ .

Now, within the small perturbation theory, where; within the framework of small perturbation theory, that is for u/v_∞ is much less than 1 and in this case, this half C_p^* is much less than 1. Of course, this result holds for wedge with very small semi vertex angle; as the vertex angle increases, the accuracy of this approximation decreases.

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$$\left(\frac{dC_p}{dM_\infty}\right)_{M_\infty=1} = \frac{4}{\gamma+1}$$

and $\left(\frac{dC_p}{dM_\infty}\right)_{M_\infty=1} = \frac{4\theta}{\gamma+1}$

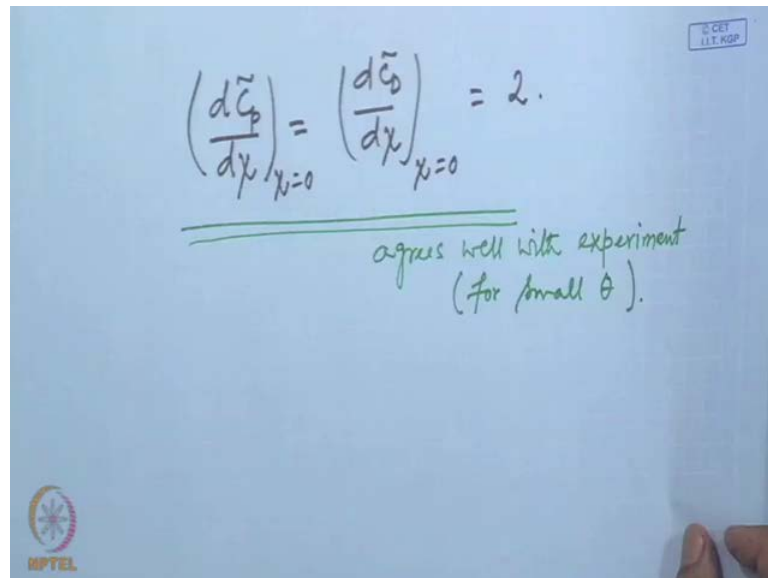
introducing $\tilde{C}_p = C_p \frac{[(\gamma+1)M_\infty^2]^{1/3}}{\theta^{2/3}}$, $\tilde{C}_D = C_D \frac{[(\gamma+1)M_\infty^2]^{1/3}}{\theta^{5/3}}$

$$X_tilde = \frac{M_\infty^2 - 1}{[(\gamma+1)M_\infty^2\theta]^{2/3}}$$

So with this approximation, **with this approximation** we have $dC_D/dM_\infty = 4\theta/(\gamma+1)$. So, this is where is the slope of the drag versus Mach number curve near $M_\infty = 1$? Now, if you introduce the transonic similarity parameter, that modifier pressure coefficient corresponding to transonic flow $C_p^* = C_p (\gamma+1) M_\infty^2 / \theta$, C_{D^*} is $C_D (\gamma+1) M_\infty^2 / \theta$ to the power $1/3$ by θ , C_{D^*} is C_D into $(\gamma+1) M_\infty^2$ to the power $1/3$ by θ to be power $2/3$ and this is θ to the power $5/3$ and the transonic similarity parameter that is $M_\infty^2 - 1$ divided by $(\gamma+1) M_\infty^2 \theta$ to the power $2/3$.

1 into M infinity squared into theta to the power 2 by 3. So, if these definitions are introduced here; if these are introduced in these two relations.

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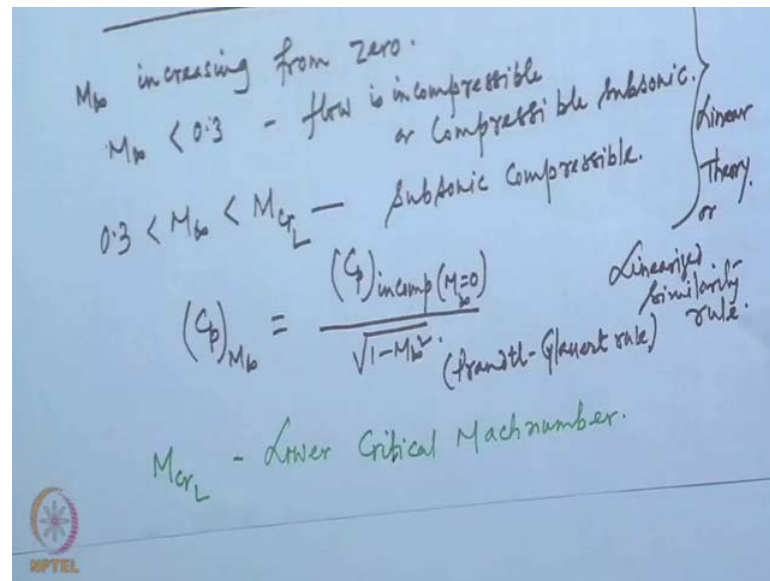

$$\left(\frac{d\tilde{C}_p}{dx}\right)_{x=0} = \left(\frac{d\tilde{C}_p}{dx}\right)_{x=0} = 2.$$

agrees well with experiment
(for small θ).

We get 2 so that, the modified drag coefficient in terms of the transonic similarity parameter, has a slope of 2 at the sonic point; free stream sonic point of course. And it found that, this result agrees well with experiment for small theta. Say of, theta is of the order of 10 to 15 degree, this result agrees quite well.

Meaning of for a wedge of 20 to 30 degree vertex angle, this result compares very well with experimental observation. And of course, that validates the transonic similarity theory and as well as, transonic small perturbation theory. Now once, after this discussion; we will now, have a brief discussion of transonic flow, about airfoils which is of course, a practical interest and perhaps, the most important as far as the aerodynamics are concerned.

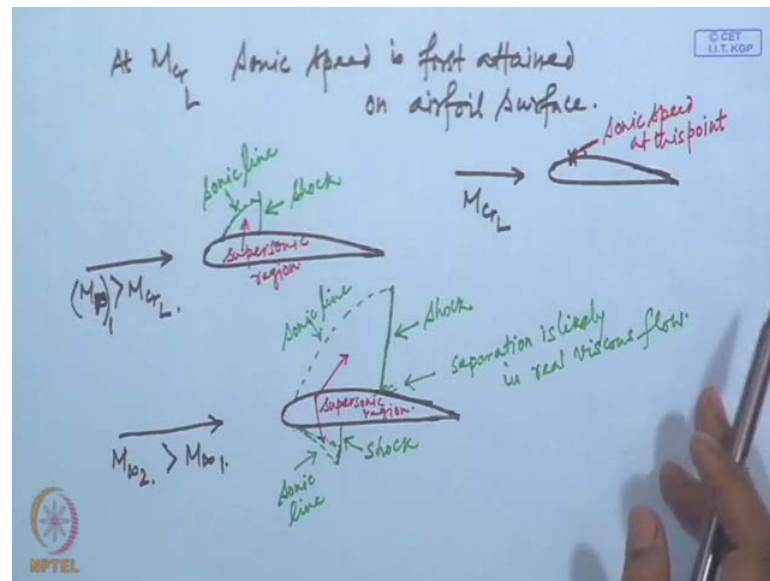
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So, let us now, consider transonic flow about airfoils. Again, let us consider that, M_∞ increasing from zero, initially up to certain Mach number; let say, as a thumb rule, which we have taken as that M_∞ less than point three, this flow is incompressible and we know, what is the flow? Of course, they can be treated as compressible subsonic, then say; we will call this M_∞ critical subsonic compressible and both can be used by both can analyzed by linearized theory linear theory or even, linear similarity rule; meaning if, we think in terms of prandtl-glauert rule.

Then once we know, the pressure distribution on the airfoil at an incompressible case, which correspondence really to Mach number 0; free stream Mach number 0. Then of course, at any other Mach number within this range, we can get the pressure coefficient simply by the prandtl-glauert rule, which is the C_p incompressible divided by square root of, so C_p at M_∞ is C_p incompressible divided by root over $1 - M_\infty^2$. When the free stream Mach number reaches M_{crL} , which is lower critical Mach number; which we have already defined in our last lecture, that M_{crL} is lower critical Mach number.

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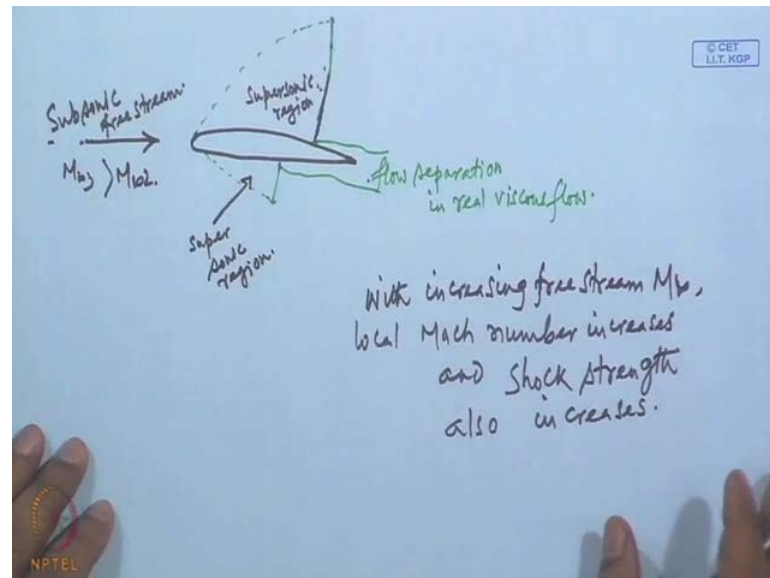
Where, sonic speed is first attained on the airfoil surface at one point. At $M_{cr,L}$, sonic speed is first attained on airfoil surface. This of course, inherently means that, we are considering inviscid flow; where the flow slips over the body and as we have seen earlier, the maximum velocity is reached on the surface. So if now, free stream Mach number increases beyond this. Then of course, there will be a subsonic region on the surface of the airfoil over certain region and as this Mach number increases, the size of the supersonic region increases and this supersonic region terminates with a shock

So let us consider, one such a free stream; remember in this discussion, we are also considering that the angle of attack is remaining constant, because this lower category Mach number depends also on the angle of attack. As this directly, effects the configuration. Let say, now this is $M_{\infty 1}$ and once again, let us denote it $M_{\infty 1}$. So, there will be; so this is shock; this is a sonic line, just to show for that $M_{\infty 1}$ sonic speed at this point. If the Mach number increase still slightly, now will have a much larger sonic region and in a viscous flow, they are might be even a separation is likely in real viscous flow.

This $M_{\infty 2}$ is of course, larger than $M_{\infty 1}$ and it is also possible to have, ((no audio 27:21 to 27:51)) this is; **super** these are all supersonic region. ((no audio 27:56

to 28:35)) Of course, the actual location where this shock will occur and at with Mach number, the shock will be relatively weak and sonic region will be little smaller. They of course, change from airfoil to airfoil at and also, changes with angle of attacks, but qualitatively that as Mach number increases, the flow behavior becomes like this.

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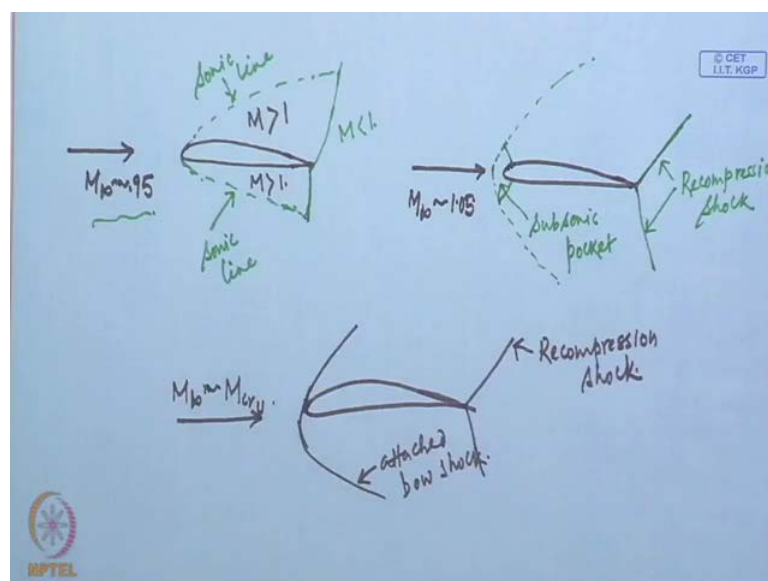
Let us now consider, free stream Mach number which is little less than one; say of the order of 0.9 or little higher then of course, the and flow separation in real viscous flow. So, in real viscous flow, there will be a strong interaction between the shock and the boundary layer and which will effects the location and strength of the shock, as well as the growth of the boundary layer. So here we see that, there is an interdependence between the shock wave and the boundary layer of course, in real viscous flow. In inviscid flow of course, that flow separation and shock boundary layer interaction will not be present.

However, as we can see that, as Mach number increases; the supersonic region on the upper surface, as well as on the lower surface increases. In all these case, subsonic free stream as we mentioned that, this corresponds to a case where Mach number may be of the order of 0.9 or so, and so you see that, as the Mach number goes on increasing; free stream Mach number goes on increasing, the supersonic region also increases and

consequently, the local Mach number also increases and the shocks become stronger. So, with increasing free stream Mach number and shock strength also increases, because if, we recall that shock strength for normal shock is $M^2 - 1$. So since, the upstream Mach ahead of the shock, the Mach number is increase.

Now, compared to the earlier cases, the shock is much strong here. So as this subsonic free stream Mach number increases towards unity, the shock strength goes on increasing. Now, let us consider two cases, which are just marginally less than so if, we go on increasing this Mach number, when the Mach number comes very close to unity? Then, the shock moves downstream, to say near about the trailing edge. Let us say two cases, if, we consider two such cases, very close to unity one subsonic and other supersonic.

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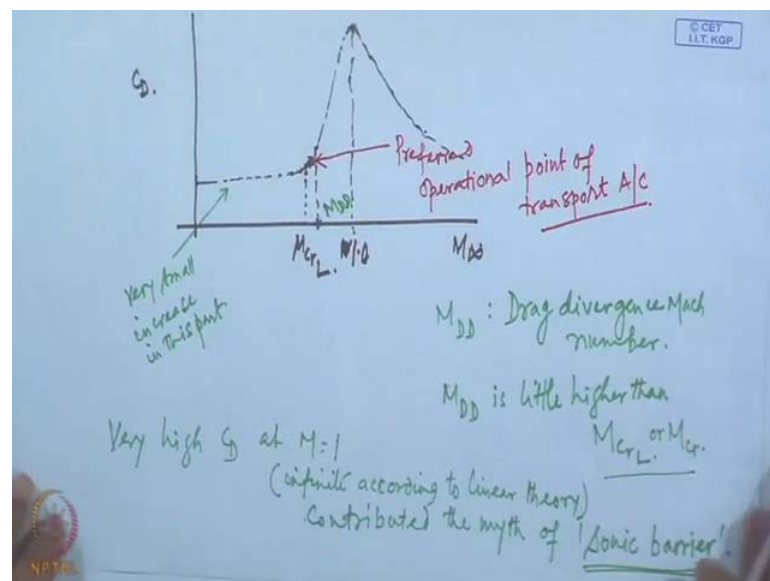


Let us say, this is M infinity; say near about 0.95 then the flow, where the airfoil is all most entirely supersonic? Of course, this number has no real significant; it simply implies that, Mach number is very close to unity. It of course, depends on configuration to configuration that is, the particular geometry of the airfoil and the angle of attack

Similarly, let us say another case, where again Mach number is very close to unity? In this case, as we know that, there will be a bow shock wave with a subsonic pocket. These

are the sonic line ((no audio 37:49 to 38:26)) and these are recompression shock. Further, higher Mach number see if, we have further higher Mach number then and of course, a recompression shock. This is M infinity, still M critical upper. So, we have briefly describe, the flow behavior or flow pattern in the transonic regime and talked about different type of shock appearance at different Mach numbers and particularly, this have a tremendous effect on the drag coefficient.

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Let us say, if we plot drag coefficient that is; it is Mach number say, nearly up to critical Mach number, very hardly any effect of increasing Mach number the drag coefficient increases marginally. In this part however, after the Mach number has cross the lower critical Mach number then, shock wave appears on the airfoil and drag started increasing very rapidly and ((no audio 43:45 to 44:16)) this is called M D D. M D D is drag divergence Mach number, which of course is, little higher than critical Mach number; little higher than.

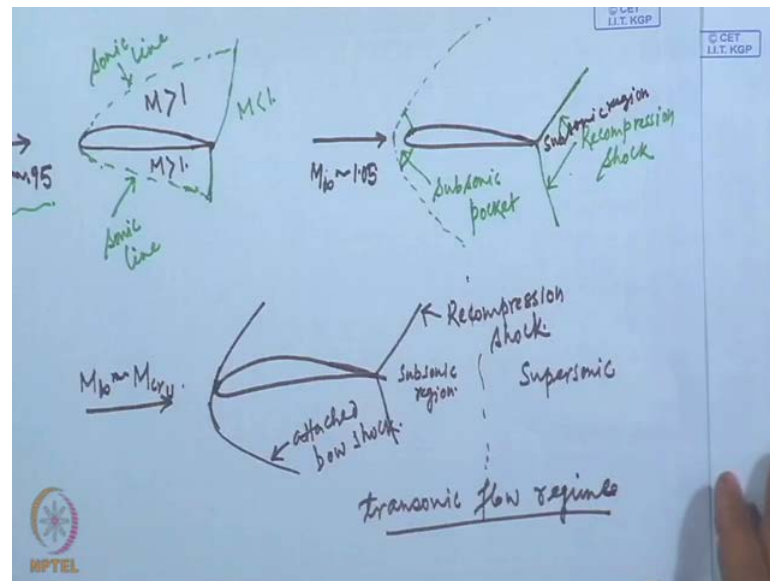
As this quite obvious, that the amongst the two critical Mach number, the lower critical Mach number is perhaps, more important as in terms of design or for aviation and hence, usually the critical Mach number is refer to lower critical Mach number. As we mentioned that, at M lower critical Mach number, the sonic point is first reach and after

that as Mach number increases, there is a supersonic sonic pocket on the airfoil, terminated by a normal shock. So, when the Mach number is increased a little bit from lower critical Mach number or critical Mach number then, the shock appears and once, the shock appears that contributes the rapid increase in drag, this and also, we have seen that, as now Mach number increases the strength of the shock increases considerably and of course, the drag rise is quite first and this high value of drag so of course, which we earlier mentioned that according to linearized theory, this would have been infinite.

So, very high C_D at infinite according to linear theory. This contributed to the myth of sonic barrier **this contributed the myth of sonic barrier**. Now of course, aircrafts which wants to operate at high speed, but without going supersonic; that is without paying the heavy price of drag rise and fuel price; that is for the transport aircraft. The most preferable point of operation is here; this is the preferred operational point of transport aircrafts that is getting the benefit of high speed, but without paying the penalty of heavy drag or high drag.

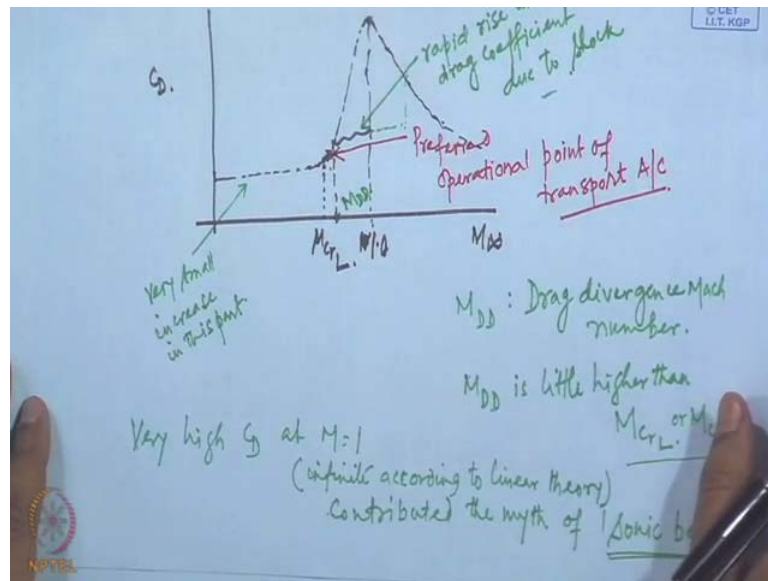
So; this critical; so we have discussed about the flow over airfoil in the transonic range and we have seen, how the supersonic region appears first in the upper surface? **Then consequently on the** Then subsequently on the lower surface, as well and then, this shock moves downstream towards the trailing edge point and the Mach number reaches very close to unity and downstream of these trailing edge shock, the flow again become subsonic. Even in this case also, this region here is a subsonic region in these cases as well.

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A subsonic region is present here and here also and then of course, a little downstream; the flow become as the free streaming supersonic, that little downstream in these cases also, the flow becomes supersonic somewhere here and so these also belongs to that transonic flow regime. We also have indicated a possible shock boundary layer interaction and separation and rapid growth in boundary layer with thickness, for a real viscous flow and with this, we have discussed how the rapid drag rise.

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So in this case, rapid rise in drag coefficient in to shock and we also have defined, what is the divergence drag? Divergence Mach number, where the shock wave becomes started contributing largely to the drag coefficient, that becomes the drag divergence Mach number. Now, it a very important to particularly, estimate this lower critical Mach number, because it is that more or less fixes the preferred operational point or cruising speed of a particular aircraft.

Of course, as we have mentioned earlier, that the lower critical Mach number changes with the airfoil shape and airfoil angle of attack. So, we will briefly describe, how we can have a preliminary approximate estimate of lower critical Mach number? And anyway we will consider that, in our next lecture the estimation of lower critical Mach number. So with this, we conclude; so in our next lecture, we will discuss about lower estimation of lower critical Mach number approximately.