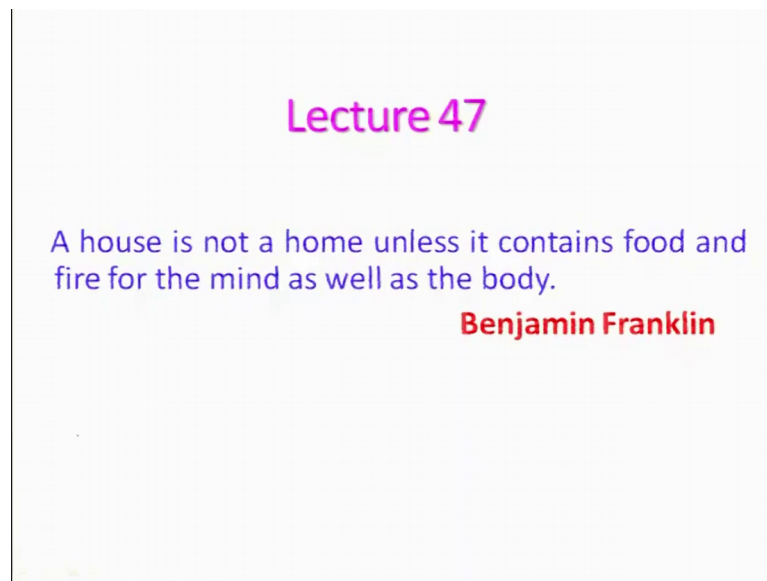


Fundamentals of Combustion (Part 2)
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Lecture – 47
Laminar Flame Theory for Premixed Flames (Contd..)

Let us start this lecture with a thought process from Benjamin Franklin.

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Who says, a house is not a home, unless it contains food and fire for the mind as well as the body; which is quite true and commensurate with the thought process of Indian scriptures. And let us recall what we learnt, in the last lecture we are trying to derive a relationship for burning velocity and in that we invoke basically all the equations of conservation of mass, momentum and energy and spaces. But we looked at the basically the energy equation and mass conservation.

And we have also found that momentum conservation equation need not to taken for 1 dimensional steady flow. Because it is a change in pressure is approximately equal to 0. And if you look at I had derived expression for energy.

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The image shows handwritten mathematical derivations and a graph. At the top left, equation (6) is written: $\rho V x \frac{dG_T}{dx} = k_g \frac{dT}{dx} + \dot{m}''_F \Delta H_c$. Below it, for the preheat zone, equation (11) is derived: $\left(\frac{dT}{dx}\right)_{i,s} = \frac{\dot{m}''_F C_p}{k_g} (T_{i,g} - T_u)$. For the reaction zone, it is noted that $T_{i,g} = T_f$ and $\frac{dT}{dx} \approx 0$, so $\rho V x \frac{dG_T}{dx} = 0$. A graph on the right shows a temperature profile T versus distance x . The preheat zone is from $x=0$ to $x=x_0$, and the reaction zone is from $x=x_0$ to $x=x_b$. The temperature starts at T_u at $x=0$, rises to $T_{i,g}$ at $x=x_0$, and then drops to T_f at $x=x_b$. The reaction rate \dot{m}''_F is shown as a curve that peaks in the reaction zone. Equation (12) is derived from equation (6) by neglecting the reaction rate term: $\int_{x=0}^x \frac{dT}{dx} k_g \frac{dT}{dx} dx = \dot{m}''_F \Delta H_c \frac{dT}{dx}$. This leads to equation (13): $\int_{x=0}^x \left(\frac{dT}{dx}\right)^2 dx = \frac{1}{2} \left(\frac{dT}{dx}\right)^2_{x=0} + \left(\frac{dT}{dx}\right)^2_{x=x_0} = -\frac{\Delta H_c}{k_g} \int_{T_u}^{T_f} \dot{m}''_F dT = \frac{\Delta H_c}{k_g} \int_{T_u}^{T_f} \dot{m}''_F dT$. Equation (14) is derived from equation (13) by using the boundary condition $\left(\frac{dT}{dx}\right)_{i,s} = \frac{2 \Delta H_c}{k_g} \int_{T_u}^{T_f} \dot{m}''_F dT$. Finally, equation (15) is derived by using equation (14) in equation (13): $\frac{\dot{m}''_F C_p}{k_g} (T_{i,g} - T_u) = \sqrt{\frac{2 \Delta H_c}{k_g} \int_{T_u}^{T_f} \dot{m}''_F dT}$.

And that is basically $\rho V x \frac{dG_T}{dx} = k_g \frac{dT}{dx} + \dot{m}''_F \Delta H_c$. This I had mentioned as equation number 6. So, later on what we did? We basically looked at the for the preheat zone. In the case of preheat zone, the reaction rate is 0.

So, therefore, we neglect the reaction rate or the heat release rate. And then we applied the boundary condition arrive at this expression; that is $\frac{dT}{dx}$ by dx ignition is equal to $\frac{\dot{m}''_F C_p}{k_g} (T_{i,g} - T_u)$. I had given equation 11; if I have not given you can take this as a equation number 11, right to be consistent. Now for the reaction zone we will have to consider, for reaction zone as the name indicates reaction zone means you know reaction will be very important.

And conduction also will be important then which term will be not important, what can be approximately any idea? Let us draw a temperature profile and this is my what we call let us say this is my $T_{i,g}$. This is T_f , this is T_u . And this is basically x is equal to 0; x is equal to infinity, minus infinity.

And this is my reaction rate you can say \dot{m}''_F right. And we know that $\frac{dT}{dx}$ is equal to 0 here. And similarly $\frac{dT}{dx}$ is equal to 0. These are the boundary condition what we have seen. In the reaction zone you can see that T and which is your reaction zone? This portion is your reaction zone right. And this portion is your preheat zone. We have

already discussed this, I am just repeating it. So, in the reaction zone, T_{ig} is approximately a closer I can say to T_F right.

So, therefore, I can write down dT/dx is equal to 0, approximately right this is not exactly. And this is again an approximation ok. As a result, we can say that $C_p \rho V x$ is approximately 0 right. Then equation 6 becomes $k_g dT^2/dx$ is equal to $-\Delta H_c$ and we will have to basically integrate this equation right. This I can say as equation 12. What I will do? I will just multiply by dT/dx here, dT/dx here right.

And I will have to integrate this equation. So, I can say this is dx/dx . So, therefore, I can write down here as dT/dx whole square right $d/dx, dx$ and this is equal to basically what? x is equal to 0 to x is equal to infinity, right. And similarly here and that itself I can write down basically half dT/dx x is equal to 0 minus dT/dx x is equal to infinity, right, is if you look at sorry, this will be infinity and this will be 0.

And by using the boundary condition, if you look at the boundary condition is basically, what is that? x is equal to 0; dT/dx at is equal to dT/dx at ignition right. And x is equal to infinity or I can write down here x is equal to infinity; dT/dx is equal to 0, right. So, if I will use this one, x is equal to infinity this will be 0, right. Yes or no? And we will I will take now the right hand side, right hand side is basically I can say that is ΔH_c by k_g ok. I have taken k_g from left hand side to right hand side in the equation 12.

And that is triple dash F integration dT ; this is basically T_{ig} right at x is equal to 0 T is equal to T_{ig} and x is equal to infinity T is equal to what? T_F , right yes or no? Make sense? And because the and if you look at the whole reaction rate from this region this is the portion of things, reaction rate which is negligible. Instead of integrating from $T_{ignition}$ to T_F , I can write down as approximately equal to ΔH_c by $k_g T_u$ by T_F m dot triple dash f dT . Can I write down this? Because I am saying this is very small quantities as compared to this reaction rate. This is the major one right.

So, therefore, instead of taking from here I can take from there. This is very important this thing you should keep in mind, ok. Are you getting? These are negligible. So, instead of say I can take from here to there from the T_u to the T_F or the x minus infinity to x is equal to infinity I can say.

Student: By F, how can you say negligible?

Because if you look at the area right because this area is very, very small, and it is why it is so? I will come to later on, but let us also we can discuss now; that is, if you look at activation energy, right E by $R U T$, we will be discussing that. Unless the temperature is having that higher, you cannot have any term from that. So, therefore, this will be very small term. But however, it will be finite, it is not that exactly 0, it will be approximately 0 or it is as compared to the heat release in the reaction zone is very, very what you call higher in the heat release rate at the reaction zone is much higher as compared to the heat release in the preheat zone that is the thing we have already discussed right.

So, therefore, I am saying instead of taking this T_{ignition} to T_{ignition} to T_F integrating I can say that. Making sense? Ok. So now, this if you look at and keep in mind that I can cancel also this minus sign, right then I can say let me tell you this is 13. The by equation 13 becomes right, I am just rewriting that equation that is all, dT by dx_{ignition} is equal to right $\frac{\Delta H_c}{2 \Delta H_c}$, there is a little problem like 2 square this will be square, ΔH_c by $k g T_u$ by T_F m dot triple dash $F dT$.

So, I can write down this as dT by dx_{ignition} right; is equal to root over 2 ΔH_c by $k g T_u$ by T_F m dot triple dash $F dT$ right. This is I am saying equation 14. Now, we have already derived this you know expression for temperature gradient at x is equal to 0 or ignition point right. You can see the equation 11 right. By using.

Student: Equation number.

Equation number 11 right, in equation 14, we can get C_p by $k g T_{\text{ignition}}$ minus T_u is equal to root over 2 ΔH_c by $k g$ integrated T_u by T_F triple dash $F dT$ right.

So, I can say this is basically equation 15 kind of things. What we will do now? We will have to now basically invoke the continuity equation right, because here the mass flux is there, that mass flux is equal to basically $\rho u S L$. Why we are doing? Because we want to find out a relationship for burning velocity that is $S L$ right. $S L$ means basically burning velocity for laminar flame, laminar premix plane.

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From Eq (7), $\ddot{m} = \rho u S L$
 we can also find out, $\Delta H_c = (\nu+1) C_p (T_F - T_u)$
 By substituting the above expressions in Eq. (15), we can get,

$$S_L = \frac{k_g}{\rho u C_p (T_F - T_u)} \sqrt{\frac{2(\nu+1) C_p (T_F - T_u)}{k_g} \int_{T_u}^{T_F} \ddot{m}_F dT} \quad \text{--- (16)}$$

Let us consider, $T_{ig} = 0.75 T_F + 0.25 T_u$; $\Rightarrow T_{ig} - T_u = 0.75 T_F + 0.25 T_u - T_u = 0.75 (T_F - T_u)$

$$\Rightarrow \frac{T_{ig} - T_u}{T_F - T_u} = 0.75 = \frac{3}{4} \quad \text{--- (17)}$$

By using Eq 17 in Eq 16, we can have

$$S_L = \frac{k_g}{\rho u C_p} \frac{1}{\sqrt{9}} \frac{1}{(T_F - T_u)^{1/2}} \frac{2(\nu+1) C_p (T_F - T_u)}{k_g} \int_{T_u}^{T_F} \ddot{m}_F dT$$

$$\Rightarrow S_L = \frac{k_g}{\rho u C_p} \frac{2(\nu+1) C_p}{9 \sqrt{9} (T_F - T_u)^{1/2}} \int_{T_u}^{T_F} \ddot{m}_F dT = \frac{2(\nu+1) C_p}{9 \sqrt{9} (T_F - T_u)^{1/2}} \int_{T_u}^{T_F} \ddot{m}_F dT = \sqrt{\frac{2(\nu+1) C_p}{9 \sqrt{9} (T_F - T_u)^{1/2}} \int_{T_u}^{T_F} \ddot{m}_F dT}$$

Where $\ddot{m}_F = \frac{1}{T_F - T_u} \int_{T_u}^{T_F} \ddot{m}_F dT = MW_F A_f C_p^{n1} C_{ox}^{n2} e^{(-E/R_u T)}$
 1 Pre exponential factor; $E =$ Activation Energy

$S_L = f(\phi, T_u, P_u, \text{ inert addition})$

So, from equation 7 right that is the continuity equation, we know that \ddot{m} is equal to $\rho u S L$. And we can also derive, we can also find out ΔH_c is approximately equal to $\nu + 1 C_p T_F - T_u$. That I think I have given as an assignment you can do that it is a very simple thing right.

So, by using by substituting the above expressions in equation 15, we can get S_L is basically is equal to k_g by $\rho u C_p T_{ig} - T_u$ right root over $2 \nu + 1 C_p T_F - T_u$, divided by $k_g T_u T F dT$, is it fine?

Now, keep in mind let me tell this is 16. Keep in mind that this we will have to simplify little further also right. I can say that I can assume that T_F , let us consider the T_F is equal to 0.75. Let us consider T_{ig} as $0.75 T_F + 0.25 T_u$. I can write down this as $T_{ig} - T_u$ is equal to $0.75 T_F + 0.25 T_u - T_u$; which is nothing but your $0.75 T_F - T_u$ right. So, by that I can write down that $T_{ig} - T_u$ divided by $T_F - T_u$ is equal to 0.75 is equal to $3/4$. This is an again an approximation because I want to relate the T_{ig} with the T_F and T_u , right.

So, if you substitute these values in equation 6, 16, right. So, I can say this is 17. Using equation 17 and equation 16, we can have S_L . What I am doing in this expression? Basically, I am taking this term right into the root. What I will get is basically, k_g square into ρu square C_p square right. In place of T_{ig} and T_u , I will be putting

basically $\frac{16}{9} \frac{1}{T_u}$ that is $\frac{1}{T_u} \left(\frac{16}{9} \right)$ minus T_u whole square right and $2 \nu + 1 C_p T_u$ minus $T_u k g T_u T F m \dot{F} dT$.

Now, these will cancel it out, kg will cancel it out, C_p will cancel it out right. So, I can write down that as $S L$ is equal to basically kg by $\rho u C_p$ I can write down $\frac{32}{9}$.

Student: Why sir?

Just hold on, I am coming, $\nu + 1$ right and $\rho u T_u T F m \dot{F} dT$. Keep in mind that these I will be considering as a α . What is α ? Thermal diffusivity, I must also tell that kg will be evaluated at the average temperature right, kind of things which will take an example and do, but T_u will be evaluated at the ambient inlet temperature, right and C_p will be evaluated at the average temperature. I will take an example to illustrate that point right. So, I can write down this as α .

Student: $T F$ minus T_u .

Yes, yes you are right. I will consider this as basically $\frac{32}{9} \alpha$ by $\rho \nu + 1$, right $\frac{1}{T F \text{ minus } T_u T_u}$ by $T F F dT$. Keep in mind that this I can write down as $\frac{32}{9} \alpha \rho \nu + 1 m \text{ triple dash } F \text{ average}$. So, where the triple dash F average is equal to what is that? $\frac{1}{T F \text{ minus } T_u T_u T F m \dot{F} \text{ triple dash } F dT}$ is equal to basically $M W F A f C F n 1 C \text{ oxidizer } n 2 e \text{ power to the } R u T$, right.

And this is from the Arrhenius relationship. I can say this portion the average mass conjunction of fuel per unit volume is nothing but that right. And if you look at we can get this you know A_f is your pre exponential factor. And C_F is the concentration of fuel, $C_{oxidizer}$ is the constant oxidizer, E is the activation energy, R_u is the universal gas constant, T the temperature, right.

So, from these what you can get is basically $S L$ is a function of what? What is its function? It is function of basically ϕ or the fuel air ratio or the right. Because ν is coming, ν is the ν you know, like 1 mole of fuel is reacting with ν moles of oxidizer going to the product of $\nu + 1$ of product ok.

So, therefore, the equivalence ratio comes over there. And it is also a function of T_u , because ρu is there, this is you can say ρu , right ρu and it is also a function of P_u ; unburnt pressure right. Of course, in this case pressure is remaining almost constant

right. And it is also you know depending on the inert addition. Like suppose, you are replacing nitrogen with argon or helium or other things.

So, these are the things which we will get from the simple you know expression and which will be looking at particularly looking at the considering the experimental data later on. And in the next lecture, I will be taking an example how to calculate the burning velocity.

Thank you very much.